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**A FE MODEL FOR STATIC LOAD
IN TILTING PAD JOURNAL BEARINGS WITH
PAD FLEXIBILITY**

Research Progress Report to the TAMU Turbomachinery Research
Consortium

TRC-B&C-02-2014

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November 2014

COMPUTATIONAL MODEL FOR TILTING PAD JOURNAL BEARINGS
TEES Project 32513/15196B

EXECUTIVE SUMMARY

A FE Model for Static Load in Tilting Pad Journal Bearings with Pad Flexibility
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Tilting pad journal bearings (TPJBs) supporting high performance turbomachinery rotors have undergone steady design improvements to satisfy ever stringent operating conditions that include large specific loads, due to smaller footprints, and high surface speeds that promote flow turbulence and hence larger drag power losses. Simultaneously, predictive models continuously evolve to include minute details on bearing geometry, pads and pivots' configurations, oil delivery systems, etc. In general, predicted TPJB rotordynamic force coefficients correlate well with experimental data for operation with small to moderately large unit loads (1.7 MPa). Experiments also demonstrate bearing dynamic stiffnesses are frequency dependent, best fitted with a stiffness-mass like model whereas damping coefficients are adequately represented as of viscous type. However, for operation with large specific loads (> 1.7 MPa), poor correlation of predictions to measured force coefficients is common. Recently, an experimental effort [1] produced test data for three TPJB sets, each having three pads of unequal thickness, to quantify the effect of pad flexibility on the bearings' force coefficients, in particular damping, over a range of load and rotational speed conditions. This paper introduces a fluid film flow model accounting for both pivot and pad flexibility to predict the bearing journal eccentricity, drag power loss, lubricant temperature rise and force coefficients of typical TPJBs. A FE pad structural model including the Babbitt layer is coupled to the thin film flow model to determine the mechanical deformation of the pad surface. Predictions correlate favorably with test data, also demonstrating that pad flexibility produces a reduction of up to 50% in damping for the bearing with the thinnest pads relative to that with the thickest pads. A parametric study follows to quantify the influence of pad thickness on the rotordynamic force coefficients of a sample TPJB with three pads of increasing preload, $\bar{F}_p = 0, 0.25$ (baseline) and 0.5. The bearing pads are either rigid or flexible by varying their thickness. For design considerations, dimensionless static and dynamic characteristics of the bearings

are presented versus the Sommerfeld number (S). Pad flexibility shows a more pronounced effect on the journal eccentricity and the force coefficients of a TPJB with null pad preload than for the bearings with larger pad preloads (0.25 and 0.5), in particular for operation with a small load or at a high surface speed ($S > 0.8$).

Turbulent flow may occur in (large size) hydrodynamic bearings operating at high rotational speeds and with low viscosity lubricants. Appendix A describes a turbulent bulk-flow model for thin films and presents predictions from the model for a large size four-pad water lubricated TPJB. The agreement with test data is remarkable as the bearing operating with turbulent flow has a much larger power loss and temperature rise.

Hydrodynamic bearings commonly use Babbitt layer as the liner to protect the surfaces of rotors while ensuring low friction. As a substitute for Babbitt, PTFE (polytetrafluoroethylene) shows better performance at high temperatures while offering a low breakaway friction. However, the high flexibility of PTFE does affect the static and dynamic forced performance of a hydrodynamic bearing. Appendix B describes a simple model for the liner elastic deformation due to hydrodynamic pressure and temperature differential. This deformation changes the operating clearance thus affecting the bearing performance. The model predictions are in good agreement with other model predictions in the literature. Liner flexibility reduces the peak film pressure and journal eccentricity.

Lastly, an addendum provides a (practical) graphical overview on the creation of the pad structural stiffness matrix using a commercial finite element software.

NOMENCLATURE

A	Cross-sectional area of a pad [m ²]
C_B	Bearing radial clearance [m]
C_p	Pad radial clearance [m]
C_{XX}, C_{YY}	Bearing damping force coefficients [N·s/m], $c=C\Omega C_p/W$
D	Bearing diameter [m]
E	Material elastic modulus [N/m ²]
E_{eq}	Equivalent elastic modulus of composite material [N/m ²]
e	Journal eccentricity [m]
f	Drag friction coefficient
h	Fluid film thickness [m]
$h_x, h_y, h_\delta, h_\xi, h_\eta$	Perturbed film thickness components due to pad motions
I	$Lt^3/12$. Pad area moment of inertia [m ⁴]
K_{XX}, K_{YY}	Bearing stiffness coefficients [N/m] $k=KC_p/W$
K_{piv}	Pivot stiffness [N/m], $k_{piv}=K_{piv}C_p/W_{max}$
K_{pad}	Pad bending stiffness [N/m] $k_{pad}=K_{pad}C_p(LD)/W$
L	Bearing length [m]
M_b	Bending moment [N·m]
M_{XX}, M_{YY}	Bearing virtual mass force coefficients [kg] $m=M\Omega^2C_p/W$
N	Journal rotational speed [rev/s], $N=\Omega/2\pi$
P	Pressure field on pad surface [Pa]
$P_X, P_Y, P_\delta, P_\xi, P_\eta$	Perturbed (first-order) pressure fields due to journal and pad motions [Pa/m]
\bar{p}	Uniform pressure applied on a pad [Pa]
R, R_B	Journal and bearing radii [m]
R_d	Distance from pivot to bearing center [m]
R_{eq}, R_n	Neutral axis of arc pad, single and two materials [m]

R_p	Pad radius [m]
R	$(R+i)/R$. Ratio of pad inner & outer radii
r_p	Pad dimensional preload [m], $r_p = C_p - C_B$ $\bar{r}_p = r_p / C_p$
S	Sommerfeld number, $S = \mu NDL (R/C_p)^2 / W$
T	Fluid film temperature [$^{\circ}$ C]
<i>Torque</i>	Bearing drag torque [N.m]
T	Pad thickness [m]
U	Pad surface nodal displacement [m]
u_t	Deformation at pad edges [m]
W	Static load applied on the bearing [N]
Z	$K + i \omega C$. Complex dynamic stiffness coefficients [N/m]
$\tilde{Z}_{\alpha\beta}$	Fluid film complex dynamic stiffness coefficients [N/m, Nm/rad] $\alpha, \beta = X, Y, \delta, \zeta, \eta$
δ_p	Pad tilt angle [rad]
$\Delta\sigma$	Perturbation in parameter σ
ΔW	External dynamic force [N]
η_{piv}	Pad transvers displacement [m]
Θ_p	Pad arc length [rad]
θ_p	Pivot angular position [rad]
μ	Oil viscosity [Pa.s]
ζ_{piv}	Pad radial displacement [m]
ρ	Oil density [kg/m ³]
Ω	Journal rotational speed [rad/s]
ω	Excitation frequency [rad/s]

Coordinate Systems

(X, Y)	Cartesian coordinates for journal center displacements
(ζ, η)	Pad pivot local coordinates
(r, θ, z)	Cylindrical coordinates for pad finite element structural model

Vectors and Matrices

A	Matrix of element surfaces
$\bar{\mathbf{F}}$	Reduced external force vector
K	Pad stiffness matrix
$\bar{\mathbf{K}}$	Reduced stiffness matrix
L	Lower triangular form
P	Pressure filed
S	Surface traction vector
U	Pad displacement vector
$\bar{\mathbf{u}}$	Reduced pad displacement vector

Superscripts

G	Global matrix
K	k^{th} pad

Subscripts

0	Static equilibrium position
α, β, σ	X, Y, δ, ξ, η
P	Pad upper surface

Acronyms

DOF	Degree of freedom
FE	Finite element
LOP	Load on pad
LBP	Load between pad
TPJB	Tilting pad journal bearing

LIST OF CONTENTS

EXECUTIVE SUMMARY.	2
NOMENCLATURE.....	4
LIST OF FIGURES.....	8
LIST OF TABLES.....	11
INTRODUCTION.....	12
REYNOLDS EQUATION FOR THIN FILM BEARINGS.....	16
ANALYSIS FOR PAD MECHANICAL DEFORMATIONS.....	18
PERTURBATION ANALYSIS OF THE THIN FILM FLOW EQUATION.....	21
COMPARISON OF PREDICTIONS AGAINST TEST DATA.....	24
PARAMETRIC STUDY ON THE EFFECT OF PAD FLEXIBILITY ON TPJB PERFORMANCE.....	41
CONCLUSION.....	53
REFERENCES.....	55
APPENDIX A HIRS TURBULENT FLOW MODEL.....	58
APPENDIX B A MODEL FOR BEARING LINER DEFORMATION.....	65

ADDENDUM. A Guide to get a pad structural stiffness matrix in ANSYS

LIST OF FIGURES

Figure 1	Schematic view of an idealized TPJB. Film thickness (h), pad deflection (u_p), pad rotation (δ_p) and pivot deflections (ζ_{piv}, η_{piv}) greatly exaggerated.17
Figure 2	(a) Typical FE model and mesh for a bearing pad, and (b) boundary conditions on pad as modeled in Ref. [14]. u_r , u_θ , u_z are the nodal displacements along the radial, angular, axial directions, respectively.19
Figure 3	Load configuration for three pad TPJB and photograph of one pad.....25
Figure 4	Illustration of three pad bearing and set up for measurement of pivot stiffness.....26
Figure 5	Journal eccentricity (e_j) along the load direction versus unit load $W/(LD)$. Journal speed $\Omega=6$ krpm and 12 krpm. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.27
Figure 6	Maximum pad-surface temperature versus unit load $W/(LD)$. Journal speed $\Omega=6$ krpm and 12 krpm. Predictions (without and with pad flexibility) and test data from Gaines [1]. Inlet oil temperature, $T_{in} = 49^\circ\text{C}$. Results shown for thin, medium and thick pads.....28
Figure 7	Real part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=6$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).....30
Figure 8	Real part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=12$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibi.....32
Figure 9	Imaginary part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=6$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).....34
Figure 10	Imaginary part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=12$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).....35
Figure 11	Direct stiffnesses (k_{xx} and k_{yy}) versus unit load and two shaft speeds. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.37
Figure 12	Predicted film pressure and film thickness at bearing mid plane. Operation with unit load $W/(LD)=172$ kPa and shaft speed $\Omega=6$ krpm. Location of the maximum film pressure for each pad: $\theta_1=33^\circ$ (pad 1), $\theta_2=153^\circ$ (pad 2) and $\theta_3=273^\circ$ (pad 3). Location of the minimum film thickness for each pad:

	$\theta_1=53^\circ$ (pad 1), $\theta_2=173^\circ$ (pad 2) and $\theta_3=301^\circ$ (pad 3).....	38
Figure 13	Direct damping coefficients (c_{XX} and c_{YY}) versus unit load and two shaft speeds. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.....	39
Figure 14	Direct virtual mass coefficients (m_{XX} and m_{YY}) versus unit load and shaft speed= 6krpm. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.....	40
Figure 15	Three-pad TPJB journal eccentricity (e/C_p) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload varies: LBP and LOP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	44
Figure 16	Three-pad TPJB drag friction coefficient (f) vs. Sommerfeld number. Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload varies: LBP and LOP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	45
Figure 17	Three-pad TPJB stiffness coefficients (k_{XX}, k_{YY}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0, 0.5$: LBP configuration. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	47
Figure 18	Three-pad TPJB stiffness coefficients (k_{XX}, k_{YY}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0.25$: LOP and LBP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	48
Figure 19	Three-pad TPJB damping coefficients (c_{XX}, c_{YY}) vs. Sommerfeld number. Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0, 0.5$: LBP configuration. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	49
Figure 20	Three-pad TPJB damping coefficients (c_{XX}, c_{YY}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0.25$: LOP and LBP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	50
Figure 21	Three-pad TPJB virtual mass coefficients (m_{XX}, m_{YY}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0, 0.5$: LBP configuration. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.....	51
Figure 22	Three-pad TPJB damping coefficients (m_{XX}, m_{YY}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0.25$: LOP and LBP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.	52
Figure A.1	Schematic view of large four-pad TPJB in Ref [A.5]	61
Figure A.2	Film thickness at mid-plane of pads #1 and #2. Current predictions and results in Ref. [A.5]. $W/(LD)=1,252$ kPa, $\Omega=3$ krpm, average $Re=2,043$.	

Figure A.3	Pressure profile at the mid-plane of pads #1 and #2. Current predictions and results in Ref. [A.5]. $W/(LD)=1,252$ kPa, $\Omega=3$ krpm, average $Re=2,043$.	61
Figure A.4	Temperature at the mid-plane of pads #1 and #2. Current predictions and results in Ref. [A.5]. $W/(LD)=1,252$ kPa, $\Omega=3$ krpm, average $Re=2,043$.	62
Figure A.5	Bearing friction loss at $W/(LD)=1,252$ kPa. Speed varying 1200 rpm to 4200 rpm. Current predictions and results in Ref. [A.5]. average Re from 1201 to 4774.	63
Figure B.1	Schematic view of plain journal bearing with a liner, as per Ref. [B.1].	66
Figure B.2	Fluid film thickness at the bearing mid-plane. Static load $W/(LD)=3,073$ kPa, rotor speed $\Omega=900$ rpm. Current predictions and data from Ref. [B.1].	68
Figure B.3	Fluid film pressure at the bearing mid-plane. Static load $W/(LD)=3,073$ kPa, rotor speed $\Omega=900$ rpm. Current predictions and data from Ref. [B.1].	68
Figure B.4	(a) Maximum pressure, (b) minimum fluid film thickness and (c) eccentricity ratio and (d) maximum temperature at different load. 504 kPa $< W/(LD) < 6,297$ kPa. $\Omega=900$ rpm. Current predictions and data from Ref. [B.1].	69

LIST OF TABLES

Table 1	Geometry, lubrication properties, and operating conditions of three test TPJBs [1].....	24
Table 2	Estimated average pivot stiffness for each pad set configuration in Ref. [1].	26
Table 3	Cases to assess effect of pad flexibility on the performance of a TPJB.....	41
Table A.1	Geometry and lubricant properties of the test TPJB in Ref.[A.5].....	60
Table B.1	Parameters of a plain journal bearing in Ref. [B.1].....	66
Table B.2	PTFE liner properties in Ref. [B.1].....	67
Table B.3	Cases to assess effect of liner flexibility on the performance of a plain journal bearing.....	67

INTRODUCTION

Tilting pad journal bearings (TPJBs) offer rotor dynamically stable performance, a distinct advantage over fixed geometry fluid film bearings [2]. This feature is attractive for high speed rotating machinery applications where high levels of damping are desirable with absence of (destabilizing) cross-coupling hydrodynamic effects. During the last decade there have been various test programs aiming to quantify the rotordynamic force coefficients of TPJBs while also elucidating the importance of excitation frequency dependence. Amongst these concerted efforts, Refs. [3-7] report that the actual damping offered by TPJBs is lower than predicted; possibly degrading with increasing rotor speed and applied load. Discrepancies between predictions and experimentally identified force coefficients, in particular under a heavy static load, $W/(LD) > 2.0$ MPa, are attributed to the physical model(s) not accounting for either pad and/or pivot flexibility [6,7].

Refer to San Andrés et al. [8,9] for a review on the effects of pivot stiffness on the force coefficients of TPJBs, the description of a predictive model and its validation against test force coefficients available in the literature, and a discussion on the impact of pivot flexibility on the performance of TPJBs.

Analyses including pad surface flexibility on the modeling of the forced response of TPJBs are well documented. In a few analyses, pad flexibility is modeled with a moment-bending stiffness derived from beam theory and a change in pad curvature that reduces the pad mechanical preload as the applied load increases [6,7,9-12]. In other analyses, the whole pad structure is modeled with finite elements (FE), for example, and active degrees of freedom synthesized to produce a system of algebraic equations relating the hydrodynamic pressure to a pad surface deformation field that affects the film thickness [13-16]. In 1978, Nilsson [10] predicts single pad force coefficients for arc lengths varying from 60° to 120° and realizes that the pad with the longest arc length bends more and affects more the damping coefficients. At a specific journal eccentricity ($e/C_B=0.5$), the direct damping coefficient along the load direction decreases by 6% for a 60° arc pad and by 21% for a 120° arc pad. Ettles [11] in 1980 produces a computational

one-dimensional beam model for pad deflections due to both a hydrodynamic film pressure and thermally induced stresses. Ettles contrasts the force coefficients of TPJBs with rigid pads against those with flexible pads, and finds an apparent reduction in damping (up to 40%) due to pad flexibility, further aggravated as the applied static load increases. Ettles notes that pad flexibility has a negligible effect on the journal eccentricity (and minimum film thickness), the maximum fluid film temperature; and for actual thick pads configurations, has no discernible effect on the bearing stiffnesses. In 1987, Lund and Pederson [12] detail a perturbation analysis to derive the bearing dynamic force coefficients including both pad and pivot flexibility. The authors find that the more flexible a pad is, the more is the reduction in the bearing damping coefficients.

To include for pad axial deformations, in 1989 Brugier and Pascal [13] build a three-dimensional (3D) finite element (FE) structural pad model to predict the elastic deformation of the pads. Predictions for a large scale three-pad TPJB with a diameter of 0.75 m and a slender ratio (L/D) of 0.75 show that both pad flexibility and pivot flexibility reduce significantly the bearing force coefficients, yet only produce a small decrease in the maximum film temperature and the minimum film thickness of the most loaded pad. The authors also predict an increase in direct stiffness coefficients at low loads due to both pad and pivot flexibility. Later in 1995, Desbordes *et al.* [14] evaluate predictions using two-dimensional (2D) and three-dimensional (3D) FE structural pad models and realize that pad deflections along the axial direction are not negligible in a heavily loaded TPJB.

Neglecting the variation of the pad mechanical deformation along the axial direction, in 1990 Earles *et al.* [15] use a two-dimensional (2D), plane strain FE model to evaluate pad flexibility effects on the forced performance of TPJBs. By assuming a pad keeps its original curvature, the complex multiple degree of freedom (DOF) system is reduced to one single parameter which reflects as change in pad radius. Frequency reduced dynamic force coefficients for a single pad correlate well, within 5% difference, with those obtained by Lund and Pedersen [12]. Earles *et al.* [16] also implement a “Pad Assembly Method” to obtain the dynamic force coefficients of a TPJB. With pivot flexibility

modeled using Hertzian contact theory, the TPJB stiffness and damping matrices contain each $(3N_{pad}+2)^2$ coefficients. Using synchronous speed reduced TPJB force coefficients in a particular rotor-bearing system, the authors report that pad flexibility produces a reduction of 6% in the predicted onset speed of instability.

Childs and students [1,3,5-7,17] utilize a test rig for measurement of the static and dynamic load performance of hydrodynamic bearings. This test rig comprises of a rigid rotor supported on rolling element bearings and with a softly supported (floating) bearing cartridge holding the test bearing element at the rotor mid span. Operators perform sets of tests spanning a range of applied static loads for one or more shaft speed conditions. Importantly enough, most tests are conducted with a constant lubricant supply rate, irrespective of the shaft speed or load condition. Orthogonally placed external shakers deliver multiple frequency dynamic loads exciting the test element over a prescribed frequency range. Bearing force coefficients – stiffness, mass and damping (K , C , M), are extracted from curve fits to experimentally derived complex stiffness functions. Recently, Wilkes [7] reports both measurements and predictions for a load on pad (LOP), five-pad TPJB with 50% pad pivot offset. Wilkes' predictions point to pad flexibility having a more pronounced effect under large loads. Do notice that Wilkes' meticulous work produces a method to measure pad flexibility by inducing moments on the leading and trailing edges of a pad and recording the pad strain. Wilkes plots the pad bending stiffness versus the applied bending moment curves and obtains the bending stiffness for the pad in his test bearing. Wilkes uses this bending stiffness to predict pad flexibility and regards the pad deformation as a change in pad clearance.

Based on the urge of Wilkes [7] to assess systematically the effect of pad flexibility on the dynamic force coefficients of TPJBs, Gaines [1] conducts a concerted effort to quantify the influence of pad thickness, e.g. its flexibility, on the static and dynamic forced performance of a three pad TPJB installed in a load between pads (LBP) configuration. Gaines used three sets of pads, each with an arc length of 90° , and varying in thickness $t=8.5$ mm, 10 mm and 11.5 mm. The unit load $W/(LD)$ applied to the bearings increases from 1,72 kPa to 1,724 kPa while operating at three rotor speeds: 6, 9

and 12 krpm (32, 48, and 64 m/s). Test data show that the direct damping coefficients decrease up to 20% as pad flexibility increases. The reduction in direct damping coefficients due to pad flexibility prevails as the rotor speed increases but keeps nearly constant as the unit load varies. However, the effect of pad flexibility on the bearing direct stiffnesses is minor. Similar to Brugier and Pascal's findings [13], Gaines also notices up to a 12% increase in direct stiffnesses as pad flexibility increases for operation with a low unit load, $W/(LD) < 689$ kPa. Even at the largest unit load of $W/(LD) = 1,726$ kPa, the direct stiffnesses of the TPJB with thin pad sets ($t = 8.5$ mm) is only 3% smaller than those of the TPJB with thick pad sets ($t = 11.5$ mm). Refer to Gaines [1] for a lucid review of the archival literature on experiments related to pad flexibility and its effects on TPJB force coefficients. See also Ref. [17] for an abridged version of the original reference and a discussion on the effects of pad flexibility within the context of practical TPJB configurations.

The current analysis extends the physical model in Refs. [8,9] to include pad flexibility to predict a TPJB static performance (e.g., journal eccentricity and minimum film thickness, fluid film temperature, etc.) and its dynamic force coefficients. A FE pad structural model including the Babbitt layer is coupled to the thin film flow model to determine the mechanical deformation of the pad surface. Exhaustive predictions are compared to the test data reported by Gaines [1]. The discussion produces relevant correlations while still stressing differences. Lastly, a parametric study follows to investigate the effects of pad and pivot stiffnesses, pad preload, and load orientation on TPJB force performance, static and dynamic. The graphs constructed using dimensionless parameters provide a bearing designer with a guide to estimate the effect of pad flexibility on journal eccentricity, drag friction coefficient, and the stiffness, damping and virtual mass force coefficients.

REYNOLDS EQUATION FOR THIN FILM BEARINGS

Figure 1 depicts a schematic view of an idealized TPJB with its major components and reference coordinate systems. The external static load (W) forces the journal spinning with rotational speed (Ω) to displace to eccentricity (e_x, e_y) away from the bearing center. The pads in the bearing generate a fluid film hydrodynamic pressure (P) that acts on the journal to oppose the applied load. The pressure also acts on each pad, tilting the pad about its pivot with rotation δ_p and displaces the pad pivot to ξ_{piv} and η_{piv} . The pressure field also deforms elastically the pads. Of importance is u_p the deformation field on the pad surface facing the fluid film.

An extended Reynolds equation¹ [8,9] with temporal fluid inertia effects governs the generation of hydrodynamic pressure (P^k) in the k^{th} pad with film thickness h^k ,

$$\vec{\nabla} \cdot \left\{ \frac{(h^k)^3}{12\mu} \vec{\nabla} P^k \right\} = \frac{\partial h^k}{\partial t} + \frac{\Omega}{2} \frac{\partial h^k}{\partial \theta} + \frac{\rho (h^k)^2}{12\mu} \frac{\partial^2 h^k}{\partial t^2} \quad (1)$$

where (z, θ) are the axial and circumferential coordinates on the plane of the bearing. The film thickness h^k is

$$h^k = u_p^k + C_P + e_x \cos \theta + e_y \sin \theta + \left(\xi_{piv}^k - r_p \right) \cos \left(\theta - \theta_p^k \right) + \left(\eta_{piv}^k - R_d \delta_p^k \right) \sin \left(\theta - \theta_p^k \right) \quad (2)$$

where C_P and C_B are the pad machined radial clearance and bearing assembly clearance, respectively, and $r_P = (C_P - C_B)$ is the pad mechanical preload. Above $R_d = R_P + t$ is the sum of the pad machined radius and pad thickness at the pivot location. The pad surface deflection field ($u_p^k > 0$) increases the film thickness.

¹ The model includes for flow turbulence and thermal energy transport effects. The analysis (equations) for these flow conditions are not included for brevity.

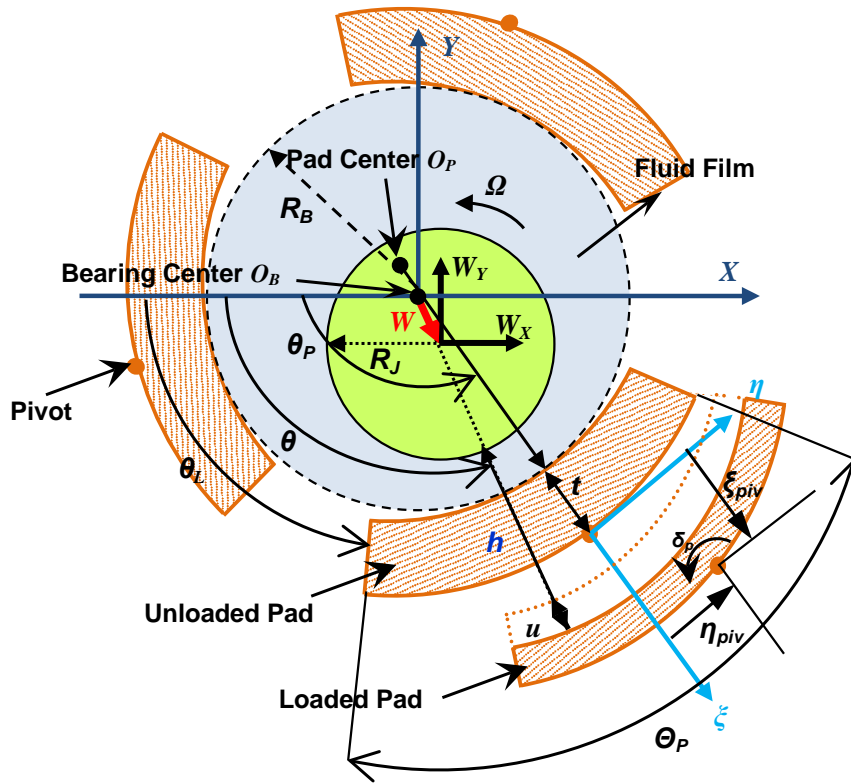


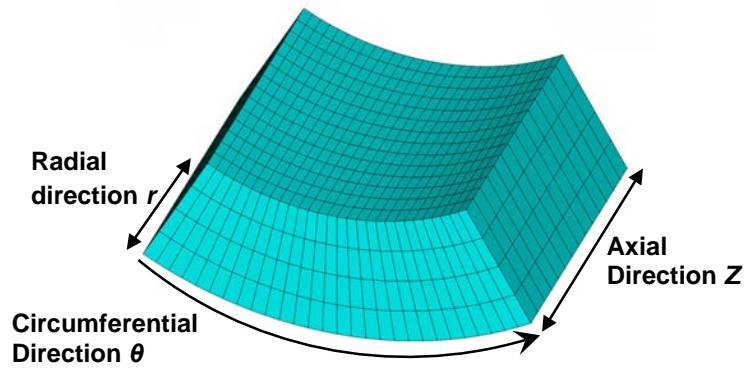
Figure 1 Schematic view of an idealized TPJB. Film thickness (h), pad deflection (u_p), pad rotation (δ_p) and pivot deflections (ξ_{piv}, η_{piv}) greatly exaggerated.

ANALYSIS FOR PAD MECHANICAL DEFORMATIONS

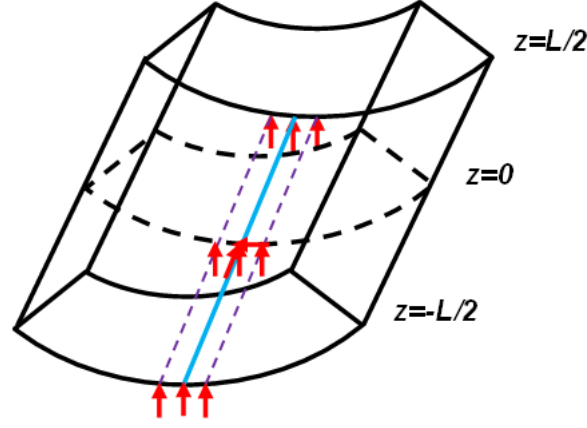
A structural FE analysis predicts the displacements of the k^{th} pad upper surface caused by the fluid film pressure field (P). Figure 2(a) depicts a typical pad assembling a number of brick-like finite elements. The equation for the deflection field (\mathbf{u}^G) due to an applied load (\mathbf{F}^G) is [14]:

$$\mathbf{K}^G \mathbf{u}^G = \mathbf{F}^G + \mathbf{S}^G \quad (3)$$

where \mathbf{K}^G is a global stiffness matrix and \mathbf{S}^G is a vector of surface tractions. Desbordes *et al.* [14] introduce appropriate boundary conditions for an ideal tilting pad, i.e., one with infinite pivot stiffness. Figure 2(b) depicts in graphical form the lines where boundary conditions are specified. The solid line denotes the pivot (line contact) and all FE nodes are constrained to a null displacement; $u_r = u_\theta = u_z = 0$, along the radial, circumferential, and axial directions. The two dashed lines parallel to the line contact denote nodes with no radial displacement, $u_r = 0$ only. On these lines, the nodes can take circumferential (transverse) and axial displacements.



(a) Pad assembled by finite elements



(b) Boundary conditions at pivot line

Figure 2(a) Typical FE model and mesh for a bearing pad, and (b) boundary conditions on pad as modeled in Ref. [14]. u_r , u_θ , u_z are the nodal displacements along the radial, angular, axial directions, respectively.

Once the pivot boundary conditions are assigned, the global system of equations reduces to

$$\bar{\mathbf{K}}^G \bar{\mathbf{u}}^G = \bar{\mathbf{F}}^G \quad (4)$$

where $\bar{\mathbf{K}}^G$ is a reduced (non-singular) stiffness matrix, and $\bar{\mathbf{u}}^G$ and $\bar{\mathbf{F}}^G$ are the vectors of global displacements and forces. The external load generated by the film pressure acts on the (upper) surface of the pad. Thus, further manipulation to reduce Eq. (4) is warranted. Write the vectors of displacements and generalized force in terms of active and inactive degrees of freedom, i.e.,

$$\bar{\mathbf{u}}^G = [\mathbf{u}_p \quad \bar{\mathbf{u}}]^T; \quad \bar{\mathbf{F}}^G = [\mathbf{f}(P) \quad \mathbf{0}]^T \quad (5)$$

where \mathbf{u}_p denotes the vector of radial displacements on the pad upper surface which are active DOFs, and $\bar{\mathbf{u}}$ is the vector of displacements of other nodes, $\mathbf{f}(P) = (\mathbf{A}\mathbf{P})$ is the vector of nodal forces generated by the pressure field \mathbf{P} with \mathbf{A} as a square matrix containing element surfaces. The reduced global stiffness matrix $\bar{\mathbf{K}}^G$ can be partitioned as

$$\bar{\mathbf{K}}^G = \begin{bmatrix} \bar{\mathbf{K}}_p & \bar{\mathbf{K}}_s \\ \bar{\mathbf{K}}_s & \bar{\mathbf{K}}_{na} \end{bmatrix} \quad (6)$$

with these definitions, Eq. (4) can be written as

$$\begin{aligned} \bar{\mathbf{K}}_p \mathbf{u}_p + \bar{\mathbf{K}}_s \bar{\mathbf{u}} &= \mathbf{f}(P) \\ \bar{\mathbf{K}}_s \mathbf{u}_p + \bar{\mathbf{K}}_{na} \bar{\mathbf{u}} &= \mathbf{0} \end{aligned} \quad (7)$$

From Eq. (7a), $\bar{\mathbf{u}} = -[\bar{\mathbf{K}}_{na}]^{-1} \bar{\mathbf{K}}_s \mathbf{u}_p$ and the displacements on the pad surface are obtained from

$$\mathbf{K}_p \mathbf{u}_p = \mathbf{f}(P) \quad (8)$$

where $\mathbf{K}_p = \left(\bar{\mathbf{K}}_p - \bar{\mathbf{K}}_s [\bar{\mathbf{K}}_{na}]^{-1} \bar{\mathbf{K}}_s \right)$ is a positive definite symmetric matrix, easily decomposed into its lower and upper triangular forms, $\mathbf{K}_p = \mathbf{L}\mathbf{L}^T$. Hence, Eq. (8) is rewritten as

$$\mathbf{L}(\mathbf{L}^T \mathbf{u}_p) = \mathbf{f}(P) \quad (9)$$

Let $\mathbf{u}_p^* = \mathbf{L}^T \mathbf{u}_p$; a backward substitution procedure solves first $\mathbf{L}\mathbf{u}_p^* = \mathbf{f}(P)$ to give \mathbf{u}_p^* ; and next, a forward substitution procedure solves $\mathbf{L}^T \mathbf{u}_p = \mathbf{u}_p^*$ to determine \mathbf{u}_p , i.e., the vector of radial displacement at the pad surface. The vector \mathbf{u}_p is used to update the film thickness (h), Eq. (2), for solution of the Reynolds Eq. (1) to find the pressure field (P). Note that the FE structural pad model and its end result, the \mathbf{L} matrix, needs to be performed only once, preferably outside of the main computational program.

PERTURBATION ANALYSIS OF THE THIN FILM FLOW EQUATION

At a constant shaft speed (Ω), the static load $\mathbf{W}_0 = (W_{X_0}, W_{Y_0})^T$ displaces the journal to its equilibrium position $\mathbf{e}_0 = (e_{X_0}, e_{Y_0})^T$ with the generated fluid film pressure (P_0^k) acting on each pad surface. The k^{th} pad reaches its equilibrium position $(\delta_{p_0}^k, \xi_{piv_0}^k, \eta_{piv_0}^k)^T$ and the deflection of the pad upper surface is $u_{p_0}^k$.

An external dynamic force, $\Delta \mathbf{W} = (\Delta W_X, \Delta W_Y)^T e^{i\omega t}$ with excitation frequency (ω) acts on the journal and causes the journal center to displace to $\Delta \mathbf{e} = (\Delta e_X, \Delta e_Y)^T e^{i\omega t}$ away from \mathbf{e}_0 , i.e., $\mathbf{e}_{(t)} = \mathbf{e}_0 + \Delta \mathbf{e} e^{i\omega t}$ [18]. The journal motion leads to changes in the pad pivot displacements and the pad surface deformation as

$$(\delta_p^k, \xi_{piv}^k, \eta_{piv}^k)^T = (\delta_{p_0}^k, \xi_{piv_0}^k, \eta_{piv_0}^k)^T + (\Delta \delta_p^k, \Delta \xi_{piv}^k, \Delta \eta_{piv}^k)^T e^{i\omega t} \quad (10a)$$

$$u_p^k = u_{p_0}^k + \Delta u_p^k e^{i\omega t}, k = 1, \dots, N_{pad} \quad (10b)$$

On the k^{th} pad, the changes in journal center position and pad displacements cause a change in the film thickness as

$$h^k = h_0^k + \Delta h^k e^{i\omega t}, k = 1, \dots, N_{pad} \quad (11a)$$

where

$$\Delta h^k = \{h_x^k \Delta e_X + h_y^k \Delta e_Y + h_\xi^k \Delta \xi_{piv}^k + h_\eta^k \Delta \eta_{piv}^k + h_\delta^k \Delta \delta_p^k + \Delta u_p^k\} \quad (11b)$$

with $h_x^k = \cos \theta$, $h_y^k = \sin \theta$, $h_\xi^k = \cos(\theta - \theta_p^k)$, $h_\eta^k = \sin(\theta - \theta_p^k)$, $h_\delta^k = -R_d h_\eta^k$ [18]. The fluid film pressure on a pad is

$$P^k = P_0^k + \Delta P^k e^{i\omega t}, k = 1, \dots, N_{pad} \quad (12a)$$

where the change in fluid film pressure caused by the perturbations in displacements is

$$\Delta P^k = \{P_X^k \Delta e_X + P_Y^k \Delta e_Y + P_\xi^k \Delta \xi_{piv}^k + P_\eta^k \Delta \eta_{piv}^k + P_\delta^k \Delta \delta_p^k\} \quad (12b)$$

Let $\mathbf{g}(P^k) = [\mathbf{K}_p^k]^{-1} \mathbf{f}(\mathbf{P}^k)$. Hence, the pad deformations caused by the equilibrium pressure field (\mathbf{P}_0^k) and the perturbed pressure field ($\Delta \mathbf{P}^k$) are

$$\mathbf{u}_{p_0}^k, \Delta \mathbf{u}_p^k = \mathbf{g}(P_0^k), \mathbf{g}(\Delta P^k) \rightarrow [\mathbf{K}_p^k]^{-1} \mathbf{A} \mathbf{P}_0^k, [\mathbf{K}_p^k]^{-1} \mathbf{A} \Delta \mathbf{P}^k \quad (13)$$

Substituting Eq. (12b) into Eq. (13) yields the change in pad surface deformation as

$$\Delta u_p^k = u_x^k \Delta e_x + u_y^k \Delta e_y + u_\delta^k \Delta \delta_p^k + u_\xi^k \Delta \xi_{piv}^k + u_\eta^k \Delta \eta_{piv}^k \quad (14)$$

Thus,

$$\begin{aligned} \Delta h^k &= (h_x^k + u_x^k) \Delta e_x + (h_y^k + u_y^k) \Delta e_y + (h_\xi^k + u_\xi^k) \Delta \xi_{piv}^k \\ &\quad + (h_\eta^k + u_\eta^k) \Delta \eta_{piv}^k + (h_\delta^k + u_\delta^k) \Delta \delta_p^k \end{aligned} \quad (15)$$

That is, the film thickness changes due to physical displacements of the journal and pad as well as due to the deformation induced by a change or perturbation in pressure.

Define the following linear operators [19],

$$\ell(*) = \frac{\partial}{R \partial \theta} \left[\frac{h_0^3}{12\mu} \frac{\partial(*)}{R \partial \theta} \right] + \frac{\partial}{\partial z} \left[\frac{h_0^3}{12\mu} \frac{\partial(*)}{\partial z} \right] = \vec{\nabla} \left(\frac{h^3}{12\mu_{(T)}} \vec{\nabla} (*) \right) \quad (16)$$

$$\Re(*) = \left[\frac{\Omega}{2} \frac{\partial(*)}{\partial \theta} + i\omega(*) - \frac{\rho h_0^2}{12\mu_{(T)}} \omega^2(*) \right] - \vec{\nabla} \left(\frac{3h_0^2}{12\mu_{(T)}} (*) \vec{\nabla} P_0 \right) \quad (17)$$

Substitution of h^k and P^k into the extended Reynolds Eq. (1) gives:

$$\ell(P_0) = \frac{\Omega}{2} \frac{\partial h_0}{\partial \theta} \quad (18)$$

and

$$\ell\{P_\sigma^k\} = \Re\{h_\sigma^k + u_\sigma^k\} \quad \alpha, \beta = X, Y, \delta, \xi, \eta \quad (19)$$

Note that the first-order or perturbed pressure fields due to a pad rotation and pivot radial and transverse displacements are a linear combination of P_X and P_Y [18,19], i.e.,

$$\begin{aligned} P_\delta^k &= -R_J P_\eta^k \\ P_\xi^k &= P_X^k \cos \theta_p^k + P_Y^k \sin \theta_p^k, \quad k = 1, \dots, N_{pad} \\ P_\eta^k &= -P_X^k \sin \theta_p^k + P_Y^k \cos \theta_p^k \end{aligned} \quad (20)$$

Since the pad deformations are a linear function of the applied pressure, i.e.,

$u_\sigma^k = g(P_\sigma^k)$ with $\alpha, \beta = X, Y, \delta, \xi, \eta$, then

$$\begin{aligned} u_\delta^k &= -R_J u_\eta^k \\ u_\xi^k &= u_X^k \cos \theta_p^k + u_Y^k \sin \theta_p^k, \quad k = 1, \dots, N_{pad} \\ u_\eta^k &= -u_X^k \sin \theta_p^k + u_Y^k \cos \theta_p^k \end{aligned} \quad (21)$$

The analysis above reveals that the perturbed pressure fields due to pad rotation or pivot transverse displacements can be readily gathered from the fields determined for changes in the journal eccentricity ($\Delta e_x, \Delta e_y$). Furthermore, the changes in pad deformation also follow immediately after the perturbed displacements (u_x^k, u_y^k) are found. The process is computationally fast and efficient. The only caveat is that Eq. (19) is solved iteratively, as $u_\sigma^k = g(P_\sigma^k)$ $\alpha, \beta = X, Y$.

Integration of the pressure fields, obtained after solving the first order Eqs. (19) on the k^{th} pad, renders 25 fluid film dynamic complex stiffness coefficients ($\tilde{Z}_{\alpha\beta}^k$) [18]

$$\tilde{Z}_{\alpha\beta}^k = \int_{-L/2}^{L/2} \int_{\theta_1^k}^{\theta_2^k} P_\beta^k h_\alpha^k R_j d\theta^k dz \quad \alpha, \beta = X, Y, \sigma, \zeta, \eta \quad (22)$$

Reduced frequency force coefficients for lateral displacements, $Z_{\alpha\beta} = (K_{\alpha\beta} + i\omega C_{\alpha\beta})$ $\alpha, \beta = X, Y$, are extracted from the complete sets of 25 Z 's by assuming all pads move with the same frequency ω . For details on the reduction process, see Refs. [18-20].

COMPARISON OF PREDICTIONS AGAINST TEST DATA

Gaines [1] produced test data for three TPJB sets, each having three pads of unequal thickness, to quantify the effect of pad flexibility on the bearings force coefficients, in particular damping, over a range of loads and rotational speed conditions. This section presents exhaustive comparisons of predictions from the current model to the test data in Ref. [1]. Gaines and Childs [17] summarize the comprehensive measurements and data analysis in Ref. [1].

Table 1 lists the geometry of the TPJBs, lubricant properties and operation conditions and Figure 3 depicts the load between pads configuration of a test bearing and a photograph of one pad with a rocker back pivot type. The bearing assembly operates with a flooded lubricant supply condition, all measurements conducted with a flow rate of 31 LPM [1]. The arc length (90°) and inner radius of the pads for the three bearing configurations are identical. A pad is composed of two materials: a main body made of steel and a Babbitt layer, 1.5 mm thick. The elastic moduli for steel and Babbitt are 200 GPa and 50 GPa, respectively [21]. The pad thicknesses (t) quoted includes the Babbitt thickness and is measured at the location noted in the photograph. Presently, in the discussion of results, the pads are referred as thin ($t=8.5$ mm), medium thickness ($t=10$ mm), and thick ($t=11.5$ mm).

Table 1 Geometry, lubrication properties, and operating conditions of three test TPJBs [1].

Number of pads, N_{pad}	3		
Configuration	LBP		
Rotor diameter, D	101.6 mm		
Pad axial length, L	61 mm		
Pad arc angle, θ_p	90°		
Pivot offset	50%		
Nominal preload, \bar{r}_p	0.25		
Pad thickness, t	8.5mm	10 mm	11.5mm

Cold bearing clearance, C_B	69 μm	70 μm	70 μm
Cold pad clearance, C_P	92 μm	93 μm	93 μm
Lubricant type	ISO VG 46		
Supply lubricant temperature	49 °C		
Supply lubricant pressure	2.2 bar		
Lubricant density	854 kg/m^3		
Viscosity at 49 °C, μ_0	0.0269 $\text{Pa}\cdot\text{s}$		
Viscosity temperature coefficient, α	0.0319 $1/^\circ\text{C}$		
Specific heat capacity at 70 °C	1830 $\text{J}/(\text{kg}\cdot\text{K})$		
Specific load, $W/(LD)$	172 kPa -1724 kPa		
Journal speed, Ω	6,000-12,000 rpm		
Surface speed, ΩR	32-64 m/s		

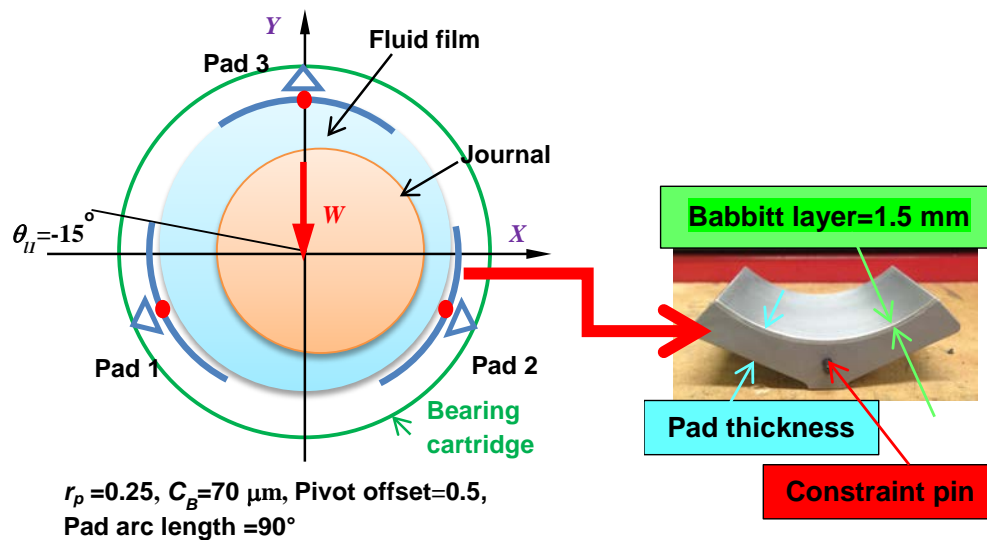


Figure 3 Load configuration for three pad TPJB and photograph of one pad.

To measure the pivot stiffness of a single pad, Gaines [1] assembled the bearing in a load on pad (LOP) orientation, as shown in Fig. 4. A hydraulic cylinder and spring pull on the bearing cartridge and displace it against the rigid rotor (journal). The applied load on the shaft is trough contact pressure over the whole pad arc extent. Eddy current

sensors record the relative displacements between the bearing casing assembly and the journal or shaft.

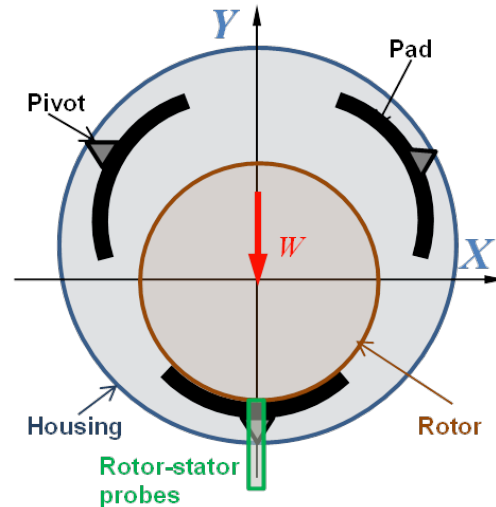


Figure 4 Illustration of three pad bearing and set up for measurement of pivot stiffness.

The load versus pivot deflection measurements conducted on pads (of differing thickness) evidence a nonlinear stiffening effect with mechanical hysteresis. For applied loads $47 \text{ kPa} < W/(LD) < 1,100 \text{ kPa}$, Table 2 lists the average pivot stiffness² (K_{piv}) used in the following predictions.

Table 2 Estimated average pivot stiffness for each pad set configuration in Ref. [1]

	Pivot stiffness [MN/m]
Thin pad ($t=8.5 \text{ mm}$)	573
Medium thickness pad ($t=10 \text{ mm}$)	675
Thick ($t=11.5 \text{ mm}$)	775

The current predictive model includes the shaft and pad thermal expansion due to rises in film temperature³ that determine the *hot* bearing clearance. The predictive model assumes that the lubricant carries away all the heat generated in the bearing, i.e., an adiabatic heat flow process. The inlet heat carry over coefficient (λ) varies for differing

² It is not clear why the pads have differing stiffness as they have identical back surface configurations and contact areas at their pivot location.

³ The shaft and pads, both made of steel, have a thermal expansion coefficient of $1.2 \times 10^{-5}/^\circ\text{C}$ [22].

rotor speeds⁴; at $\Omega=6$ krpm, $\lambda=0.8$; while at $\Omega=12$ krpm, $\lambda=0.98$. Predictions with and without pad flexibility are presented.

Figure 5 shows the predicted and measured journal eccentricity (e/C_p) along the load direction (Y) for operation at 6 krpm and 12 krpm. The predictions include curves for pad flexibility and without it. At both operational speeds ($\Omega=6$ krpm, 12 krpm), the predictions with pad flexibility correlate best with the test data as the applied load increases. Predictions solely considering pivot flexibility deliver a larger static eccentricity. Pad flexibility tends to reduce the predicted journal eccentricity, in particular for operation at the high rotor speed ($\Omega=12$ krpm).

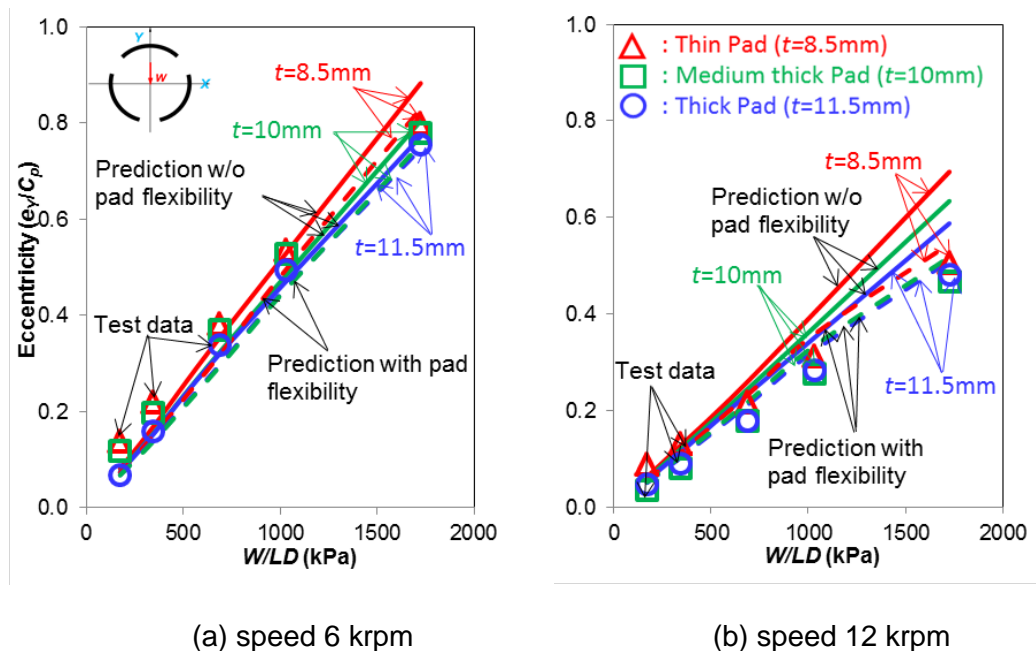


Figure 5 Journal eccentricity (e_y) along the load direction versus unit load W/LD . Journal speed $\Omega=6$ krpm and 12 krpm. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.

Figure 6 depicts the predicted maximum film temperature rise ($\Delta T_{max}/T_{in}$) versus unit load for two journal speeds. In the tests [1], a pad (~ 5 mm below) sub-surface

⁴ According to Tao [20], the inlet heat carry over coefficient (λ) varies with journal speed; a larger λ should be used for a high rotor speed ($\Omega>10$ krpm).

temperature is recorded at 75% of the pad arc length. Predictions account for the heat transfer conducted through a pad and the heat convection in the back of a pad. The predictions show the bulk fluid film temperature in pad 2. Generally, the maximum temperatures are underestimated, in particular at the high speed and largest load; $\Omega=12$ krpm, $W/(LD)=1,726$ kPa.

The significant differences in temperatures are due to the test bearing being supplied with a fixed flow rate, irrespective of the load and journal speed condition. For test with a shaft speed of 6 krpm, the actual supplied flow rate (31 LPM) is greater than the one predicted (17 LPM) while the recorded peak pad surface temperatures are much higher. Note the oil is pushed through three holes and not routed efficiently into the bearing pads. To support the assertion, note that the recorded lubricant outlet temperature is much lower than the peak temperatures measured [1], and also lower than the predicted lubricant temperature at the bearing exit plane. Hence, excessive churning of the lubricant on the bearing sides contributes to the distinctive differences. Most of the cold supply flow rate does not enter the bearing pads, thus causing the bearing pads to heat excessively.

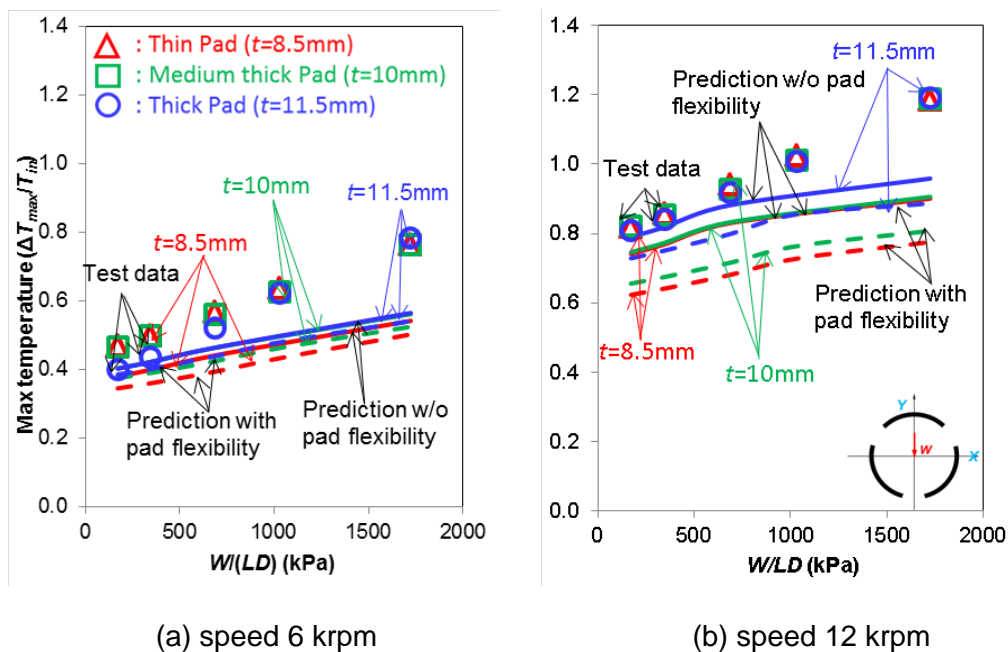
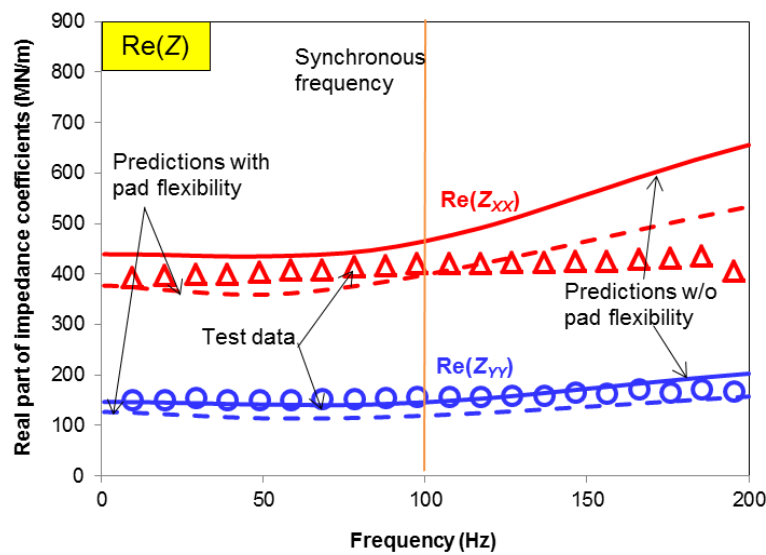


Figure 6 Maximum pad-surface temperature versus unit load $W/(LD)$. Journal speed $\Omega=6$ krpm and 12 krpm. Predictions (without and with pad flexibility) and

test data from Gaines [1]. Inlet oil temperature, $T_{in} = 49^{\circ}\text{C}$. Results shown for thin, medium and thick pads.

For the largest applied static load, Figs. 7 and 8 show the real part of the bearing direct complex stiffnesses, $\text{Re}(Z)$, obtained at two shaft speeds ($\Omega=6$ krpm and 12 krpm) versus excitation frequency ($0 < \omega < 200$ Hz). Note $\text{Re}(Z_{YY})$, along the load direction, is lesser than $\text{Re}(Z_{XX})$. This peculiar behavior is distinctive for the three pad bearing, each pad having a large (90°) arc extent. In general, the predicted $\text{Re}(Z_{YY})$ correlates best with the test data, whereas $\text{Re}(Z_{XX})$ is overestimated at high frequencies ($\omega > 100$ Hz). Note the experimental $\text{Re}(Z)$ show little frequency dependency, yet the predictions forward a stiffening $\text{Re}(Z_{YY})$ as frequency increases, and in particular for operation at the low shaft speed (6 krpm). Including pad flexibility reduces the dynamic stiffness, $\text{Re}(Z)$; the effect being more pronounced on the thin pad.



(a) Thin pad, $t=8.5$ mm

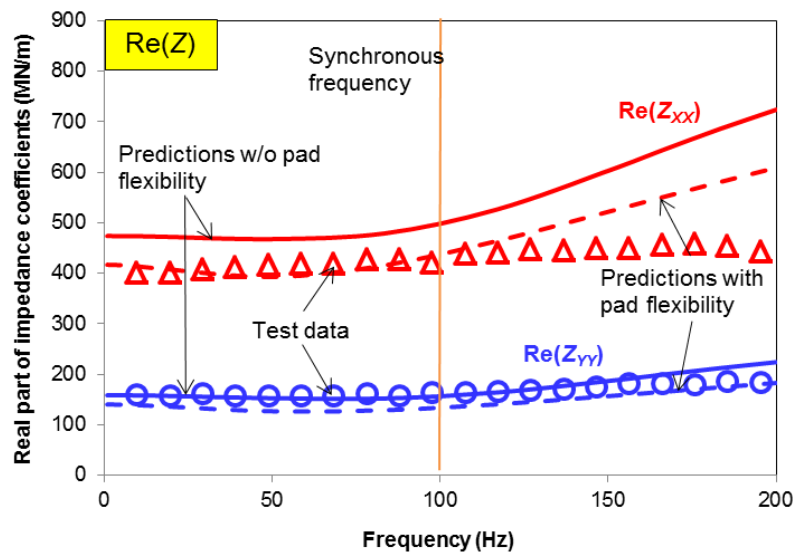
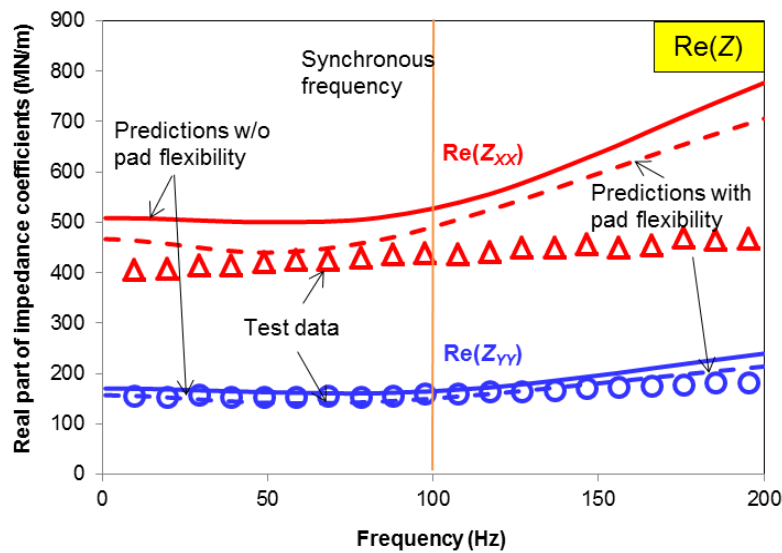
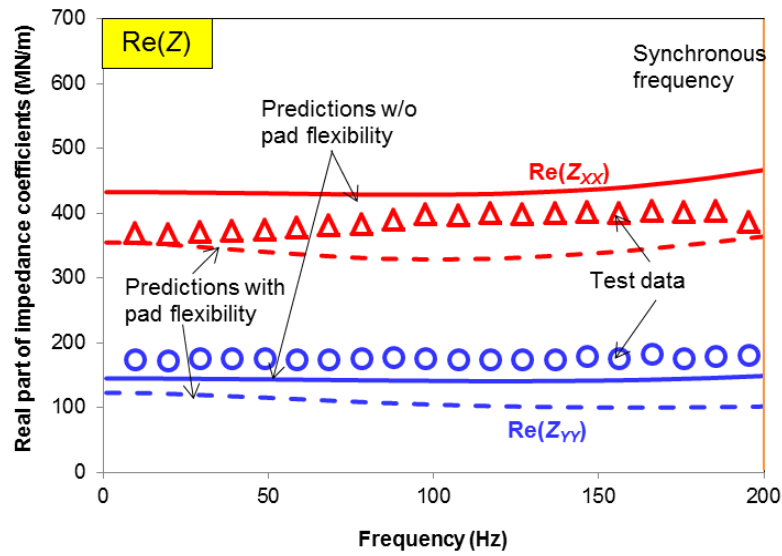
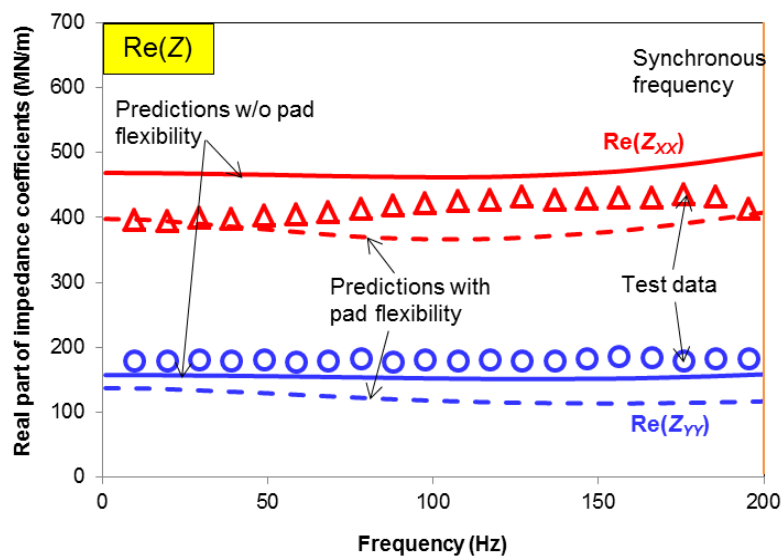
(b) Medium thickness pad, $t=10$ mm(c) Thick pad, $t=11.5$ mm

Figure 7 Real part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=6$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).

(a) Thin pad, $t=8.5$ mm(b) Medium thickness pad, $t=10$ mm

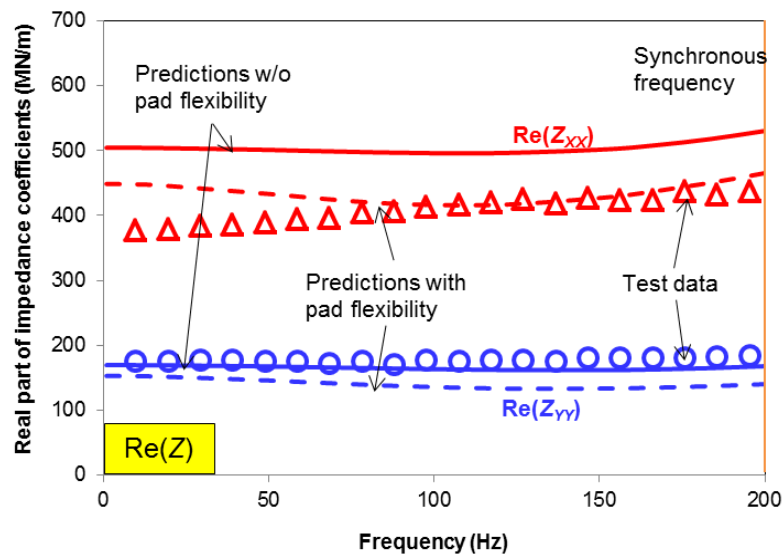
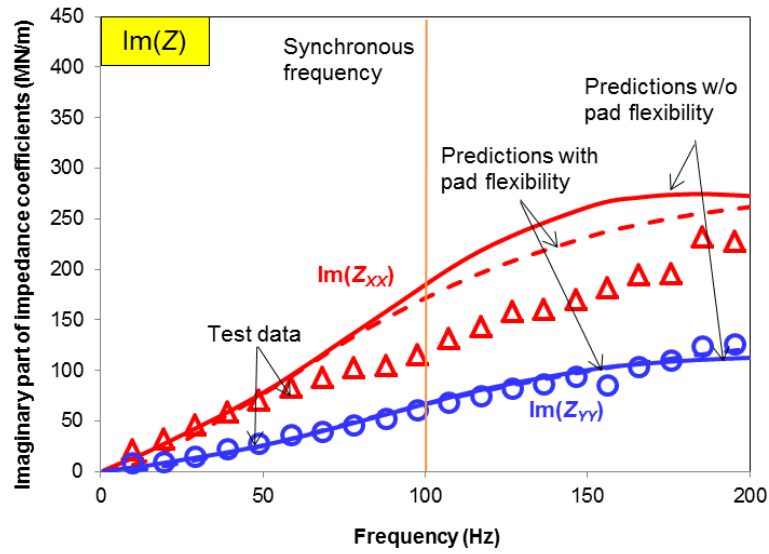
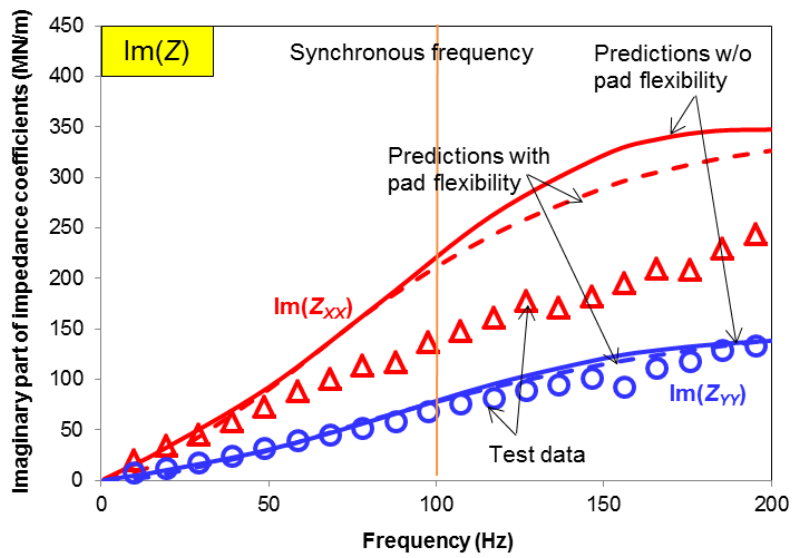
(c) Thick pad, $t=11.5$ mm

Figure 8 Real part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=12$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).

For the same static load condition and two journal speeds, Figs. 9 and 10 depict the experimental and predicted imaginary part of the bearing complex stiffnesses, $\text{Im}(Z)$. In general, the bearing damping coefficient (C) is the slope of $\text{Im}(Z) \sim \omega C$. Both $\text{Im}(Z_{XX})$ and $\text{Im}(Z_{YY})$ from the experiments show a linear growth with frequency, i.e., a frequency independent C . Note $\text{Im}(Z_{XX}) > \text{Im}(Z_{YY})$. The predictions are in very good agreement with the experimental results for operation with the high shaft speed (12 krpm). On the other hand, for operation at 6 krpm, the predicted $\text{Im}(Z_{XX})$ is larger than the test results and evidences a reduction in growth on the high side of the excitation frequency range ($\omega > 1.5\Omega$).

(a) Thin pad, $t=8.5 \text{ mm}$ (b) Medium thickness pad, $t=10 \text{ mm}$

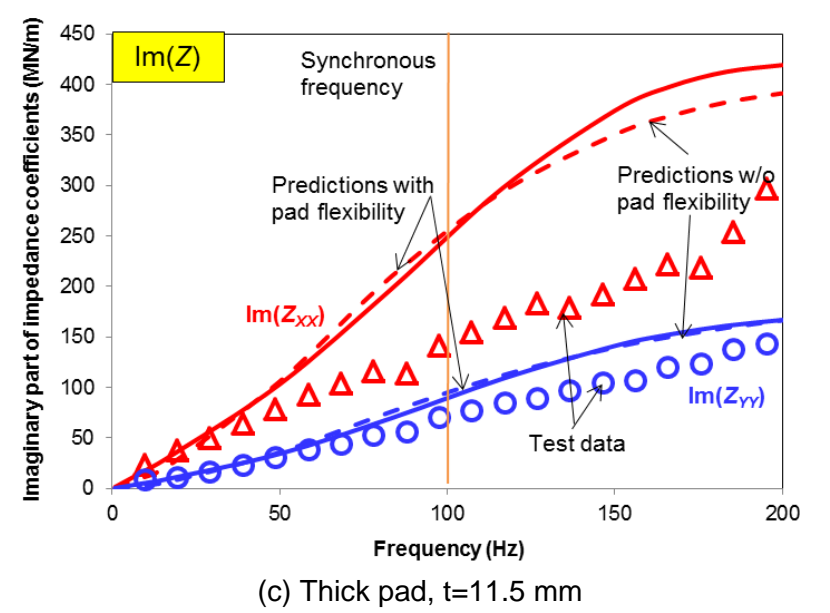
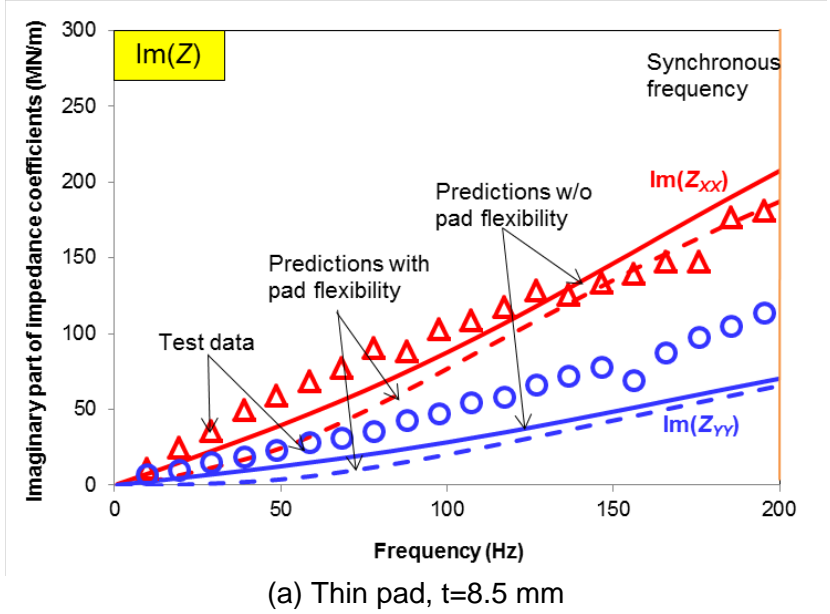


Figure 9 Imaginary part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=6$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).



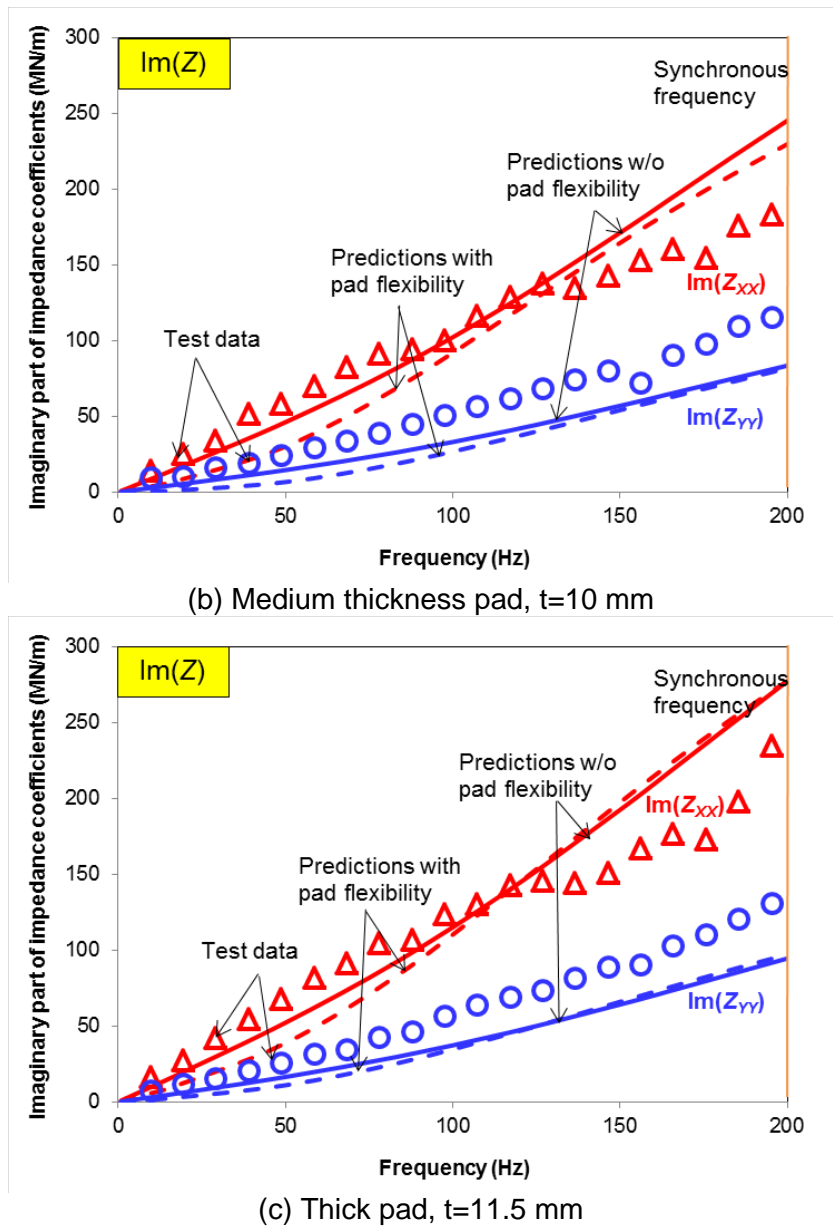


Figure 10 Imaginary part of complex stiffnesses for TPJBs with pads of thickness (a) $t=8.5$ mm (b) $t=10$ mm (c) $t=11.5$ mm. Shaft speed $\Omega=12$ krpm and unit load $W/(LD)=1,726$ kPa. Test data from Gaines [1] and predictions (with and without pad flexibility).

Childs [23], supported by a myriad of experimental results from numerous test bearings [1,3-7], stresses that the (issue on) frequency dependency of force coefficients in TPJBs is settled with a simple $[K,C,M]$ model. That is the curve fits $\text{Re}(Z) \leftarrow (K - \omega^2 M)$,

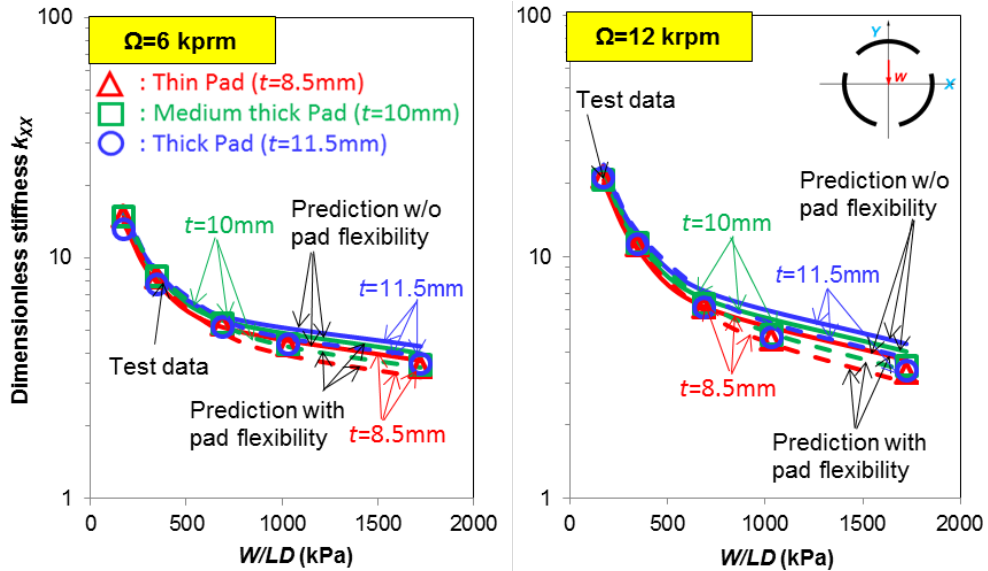
$\text{Im}(Z) \leftarrow (\omega C)$ deliver accurate stiffness (K), damping (C), and virtual mass (M) coefficients for confident use in rotor-bearing system stability analysis, i.e., a reliable estimation of the system logarithmic decrement.

Presently, define dimensionless dynamic force coefficients as [24],

$$k_{ij} = \frac{K_{ij} C_p}{W}, c_{ij} = \frac{C_{ij} \Omega C_p}{W}, m_{ij} = \frac{M_{ij} \Omega^2 C_p}{W} \quad i,j=X,Y \quad (23)$$

Above C_p is the cold pad radial clearance, Ω is the rotor speed, and W is the static load applied to the bearing.

Figure 11 depicts the TPJB stiffness coefficients ($k_{XX} > k_{YY}$) versus unit load as identified (curve fits) from the measured and predicted complex stiffnesses (Z). The predicted stiffnesses correlate well with the test data at low loads, $W/(LD) < 1,032$ kPa, but are underestimated at the highest load, $W/(LD) = 1,724$ kPa. Note that pad flexibility increases the predicted k_{XX} and k_{YY} at low loads, $W/(LD) < 689$ kPa, whereas it reduces the predicted k_{XX} and k_{YY} for high loads, $W/(LD) > 689$ kPa. Predicted direct stiffnesses accounting for pad flexibility are up to 20% smaller than those assuming a rigid pad. As the pad thickness decreases from 11.5 mm to 8.5 mm, the predicted k_{XX} decreases by 21%.



(a) k_{XX}

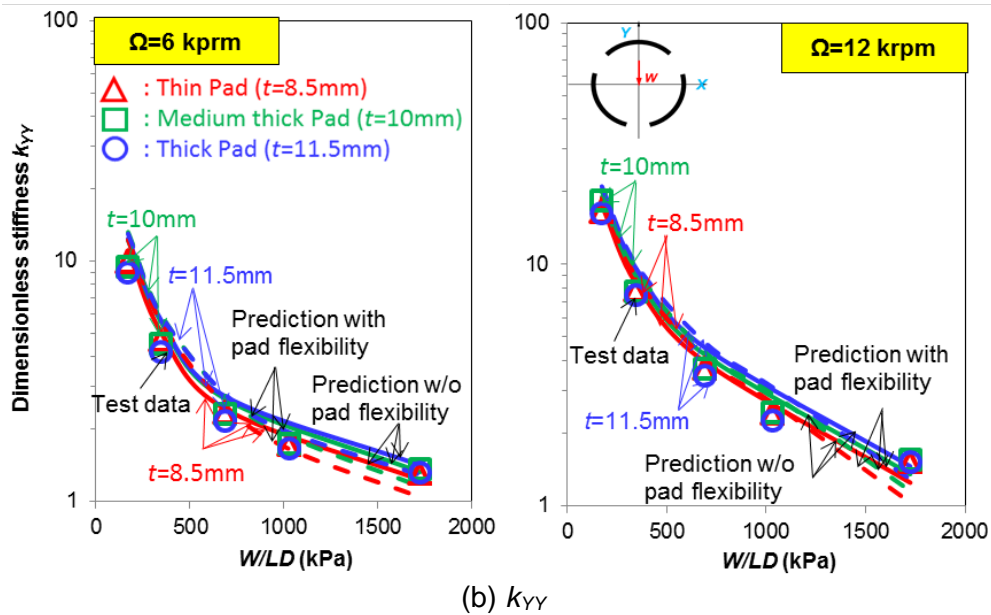


Figure 11 Direct stiffnesses (k_{XX} and k_{YY}) versus unit load and two shaft speeds. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.

Interestingly enough, the direct stiffness (k_{YY}) along the static load direction ($-Y$) is significantly lower than the stiffness k_{XX} , in particular as the unit load increases. Fig. 12 depicts the film thickness and hydrodynamic pressure at the bearing mid-plane ($z=1/2 L$). Both the minimum film thickness and the maximum pressure are quite close to the X axis ($\theta=180^\circ$), thus causing a large stiffness along the unloaded direction (X). That is, the stiffening effect is a result of the long arc extent of the bearing pads, 90° .

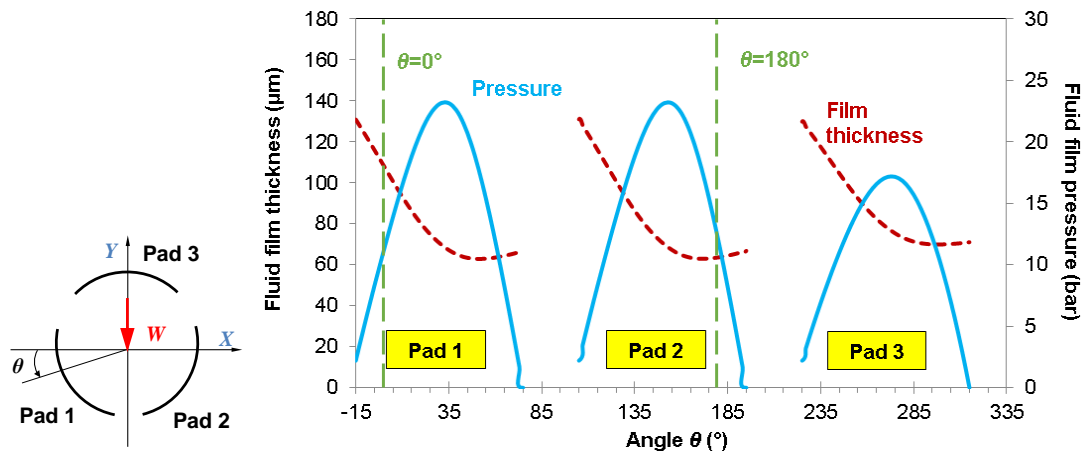


Figure 12 Predicted film pressure and film thickness at bearing mid plane. Operation with unit load $W/(LD)=172$ kPa and shaft speed $\Omega=6$ krpm. Location of the maximum film pressure for each pad: $\theta_1=33^\circ$ (pad 1), $\theta_2=153^\circ$ (pad 2) and $\theta_3=273^\circ$ (pad 3). Location of the minimum film thickness for each pad: $\theta_1=53^\circ$ (pad 1), $\theta_2=173^\circ$ (pad 2) and $\theta_3=301^\circ$ (pad 3).

Figure 13 depicts the damping coefficients ($c_{xx} > c_{yy}$) versus unit load and for two shaft speeds. Pad flexibility reduces the predicted damping over the entire load range, $172 \text{ kPa} < W/(LD) < 1,726 \text{ kPa}$. The experimental results show less differences for the three pad thicknesses than the model otherwise predicts. Predictions including pad flexibility deliver damping coefficients up to 46% lower than similar coefficients obtained with a rigid pads model. Reducing the pad thickness from 11.5 mm to 8.5 mm produces also a reduction of up to 50% in predicted direct damping. Note that the test results appear to agree with the predictions including pad flexibility for operation at the low speed of 6 krpm. The opposite argument applies to the test results at the high speed of 12 krpm, as they appear to agree with the predictions without pad flexibility.

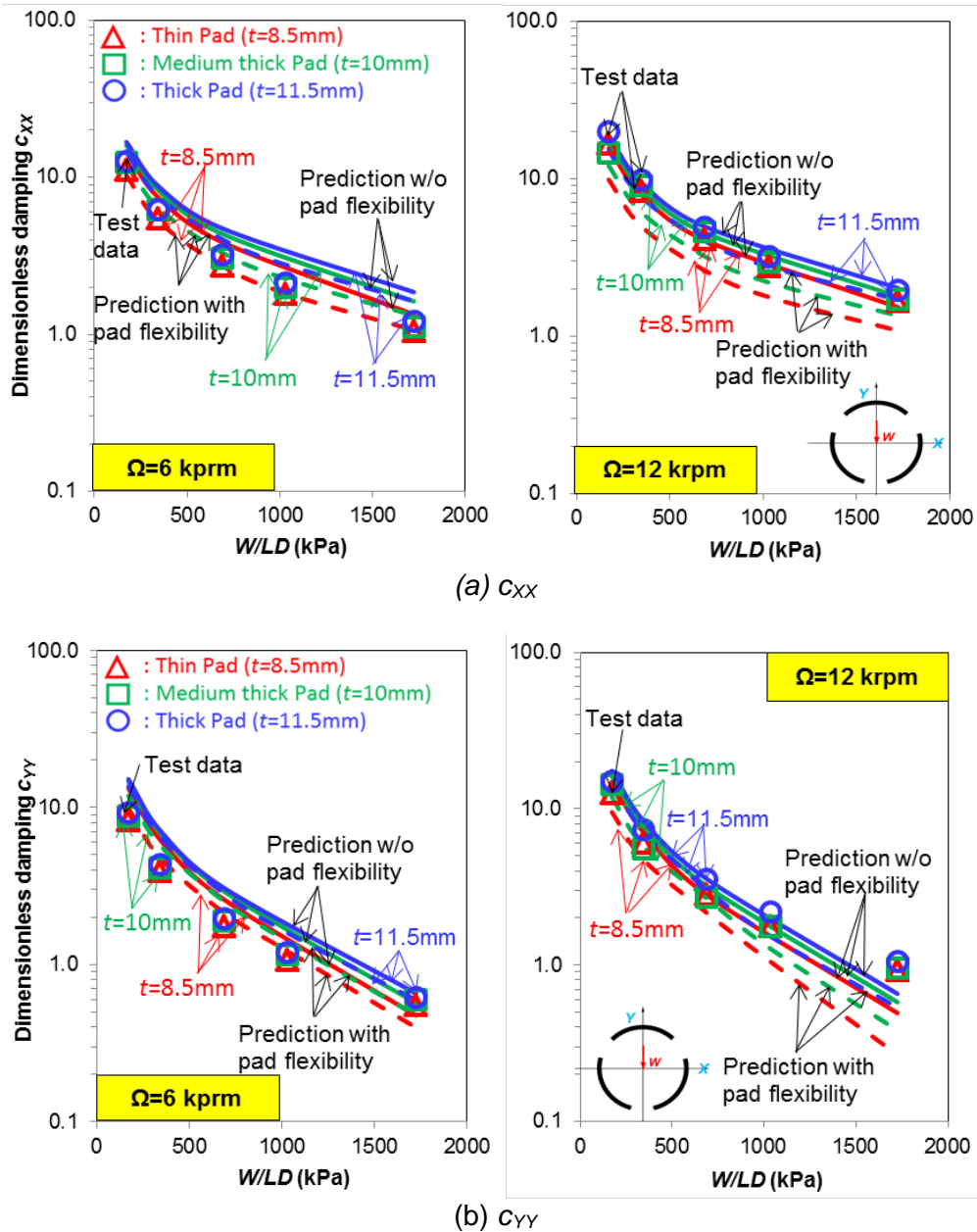


Figure 13 Direct damping coefficients (c_{xx} and c_{yy}) versus unit load and two shaft speeds. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.

For completeness, Fig. 14 displays the virtual mass coefficients (m_{xx} , m_{yy}) versus unit load and operation at the low shaft speed of 6 krpm. The test results evidence lesser magnitudes for the added masses than the predictions otherwise show. The negative

values denote the bearing dynamic stiffness hardens slightly as frequency increases, see Fig. 7. Most importantly, as the unit load increases, note (m_{xx}, m_{yy}) approach null values, thus indicating the real part of the complex stiffness (Z) does not show a frequency dependency. Similar results follow for operation at 12 krpm, hence those predictions are omitted for brevity. Note that in the sub-synchronous frequency range ($\omega < \Omega$), the virtual mass coefficients have a negligible impact on the dynamic stiffnesses ($\text{Re}(Z)$).

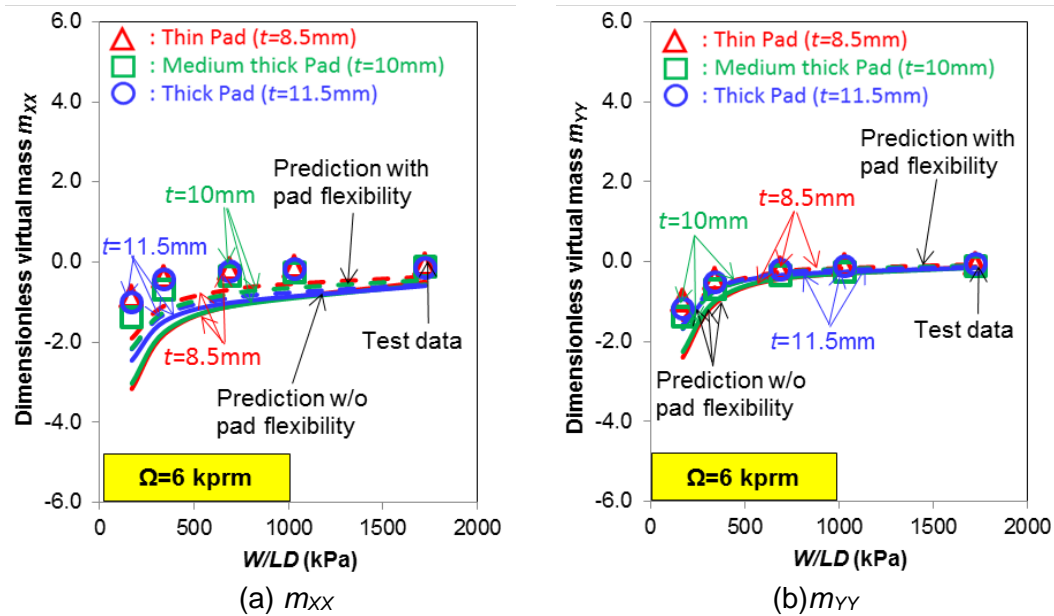


Figure 14 Direct virtual mass coefficients (m_{xx} and m_{yy}) versus unit load and shaft speed= 6krpm. Predictions (without and with pad flexibility) and test data from Gaines [1]. Results shown for thin, medium and thick pads.

PARAMETRIC STUDY ON THE EFFECT OF PAD FLEXIBILITY ON TPJB PERFORMANCE

Nicholas [2] and Someya [24] produced useful design performance data for TPJBs. However, the original data do not include either pad or pivot flexibility. San Andrés *et al.* [9] complement the design data by accounting for pivot flexibility. This section further supplements TPJB performance characteristics with consideration of both pad and pivot flexibility for a three-pad TPJB operating under both LOP and LBP configurations. The bearing configurations selected reproduces the geometry of the bearings in Ref. [1] and also include variations in pad preload (\bar{r}_p). Table 3 presents the operating conditions and geometrical parameters for the TPJB model. The TPJB performance parameters are shown as a function of the Sommerfeld number (S), defined as

$$S = \frac{\mu_s NLD}{W} \left(\frac{R}{C_p} \right)^2 \quad (24)$$

where $N = \Omega/2\pi$ is the rotational speed and μ_s is the lubricant viscosity at the supply or inlet temperature.

Table 3 Cases to assess effect of pad flexibility on the performance of a TPJB.

Static specific load, $W/(LD)$	689 kPa
Journal speed, Ω	500 rpm – 12,000 rpm
Pad preload, \bar{r}_p	0, 0.25, 0.5
Pad thickness, t	Rigid pad, 8.5 mm, 11.5 mm
Pivot stiffness, K_{piv}	750 MN/m

Branagan [25] introduces several equations to estimate a pad elastic deformation due to a bending moment (M). He suggests as adequate,

$$u_t = -\frac{M}{AE} (r-1) \frac{4[1-r^2(1-2\log r)]}{[r^2-1]^2 - (2r \cdot \log r)^2} \quad (25)$$

with $r = (R+t)/R$ and A is the cross-sectional area of a pad. Nilsson [10] already notes that a pad with a long arc length is more flexible; alas, Branagan's equation does not account

for pad arc length. In addition, Eq. (25) neglects the Babbitt layer that makes a pad more flexible.

Prior to presenting predictions, let's define a suitable pad bending stiffness (K_{pad}). Since a typical bearing pad integrates a thick base material (steel) and a Babbitt layer, define an equivalent elastic modulus as [26]

$$E_{eq} = \frac{E_1 \left(\frac{1}{2}t_1 - R_n\right)t_1 + E_2 \left(\frac{1}{2}t_2 - R_n\right)t_2}{\left(\frac{1}{2}t - R_{eq}\right)t} \quad (26)$$

where E_1 and E_2 are the elastic moduli of materials 1 and 2, t_1 and t_2 are thicknesses, and $t=t_1+t_2$ is the pad thickness. Above, R_n is the neutral axis of a pad (beam) of two materials, whereas R_{eq} is the neutral axis of a pad made of one material,

$$R_n = \frac{E_1 t_1 + E_2 t_2}{E_1 \ln\left(1 + \frac{t_1}{R}\right) + E_2 \ln\left(1 + \frac{t_2}{R+t_1}\right)}, R_{eq} = \frac{t}{\ln\left(1 + \frac{t}{R}\right)} \quad (27)$$

A simple elasticity analysis, based on information gathered from Refs. [27-30], takes a pad as a curved beam and apply a uniform pressure (\bar{p}) to determine the radial deformation at its edge (u_t),

$$\frac{u_t}{\bar{p}} = \frac{1}{2} \cdot \frac{LR_{eq}^4}{E_{eq}I} \left[\left(\cos \frac{\Theta_P}{2} - 2 \right) \cos \frac{\Theta_P}{2} + 1 \right] \quad (28)$$

where Θ_P is the pad arc extent and $I=Lt^3/12$ is the area moment of inertia. See Ref. [31] for details on the derivation of the equations above. Define a dimensionless pad stiffness as

$$k_{pad} = \frac{C_p}{W/(LD)} K_{pad} = \frac{C_p}{u_t} \cdot \frac{\bar{p}LD}{W} = \frac{E_{eq}I}{R_{eq}^3} \cdot \frac{D}{R_{eq}} \frac{2}{\left[\left(\cos \frac{\Theta_P}{2} - 2 \right) \cos \frac{\Theta_P}{2} + 1 \right]} \left(\frac{C_p}{W} \right) \quad (29)$$

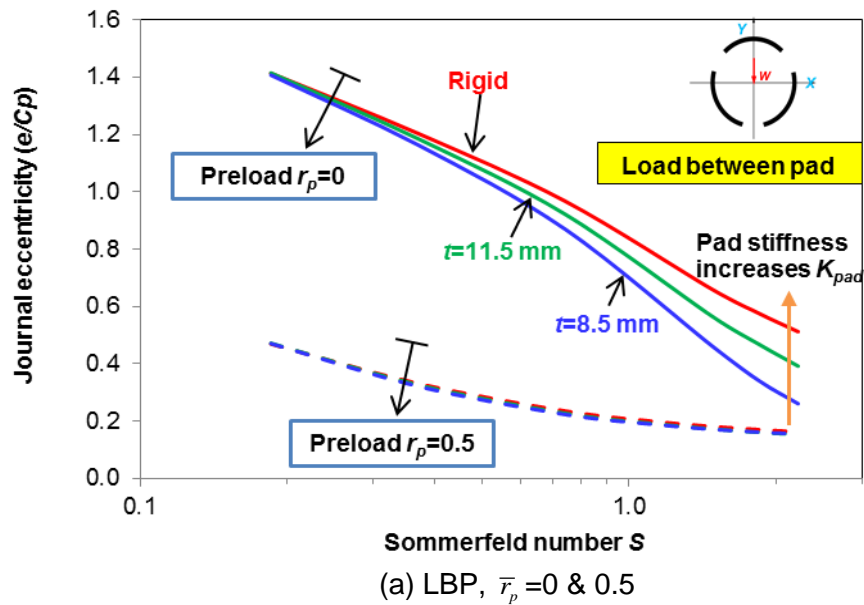
Thus, $k_{pad} = \infty$, 3.15 and 7.33, for a rigid pad, a pad with a thickness of 8.5 mm, and another pad with a thickness of 11.5 mm, respectively⁵. Note that the k_{pad} is nearly

⁵ Recall a pad has a 1.5 mm Babbitt layer.

proportional to the cube of a pad thickness, $\sim t^3$. Incidentally, in the analysis, the dimensionless pivot stiffness (k_{piv}) is, with $K_{piv}=750$ MN/m,

$$k_{piv} = \frac{K_{piv} C_p}{W_{max}} = 16 \quad (30)$$

For the LBP and LOP configurations and varying pad preload, $\bar{r}_p = 0, 0.25, 0.5$, Fig. 15 depicts the journal eccentricity versus Sommerfeld number (S). At a large S and as the pad flexibility increases (k_{pad} decreases), the journal eccentricity decreases greatly for the LBP bearing with a null pad preload and the LOP bearing with 0.25 preload, in particular.



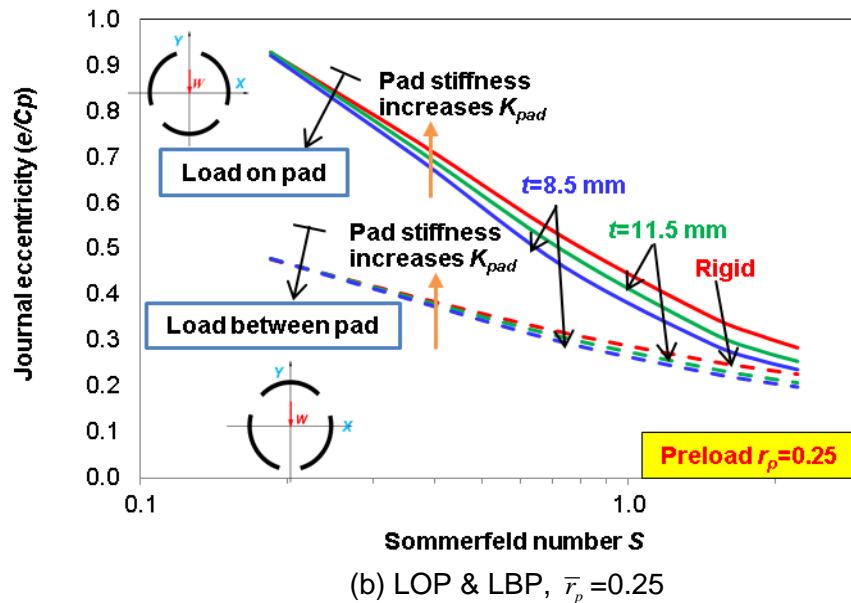


Figure 15 Three-pad TPJB journal eccentricity (e/C_p) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload varies: LBP and LOP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.

Figure 16 depicts the drag friction coefficient, $f=Torque/(RJW)$ increasing proportionally with S for both the LBP and LOP configurations, the bearing with the largest preload having more drag. Pad flexibility has no effect on f , hence has no influence on the bearing drag power losses.

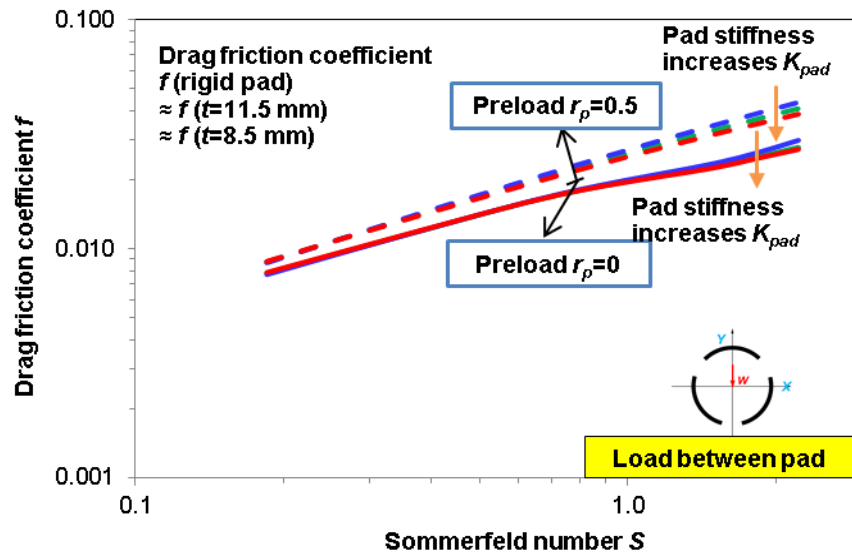
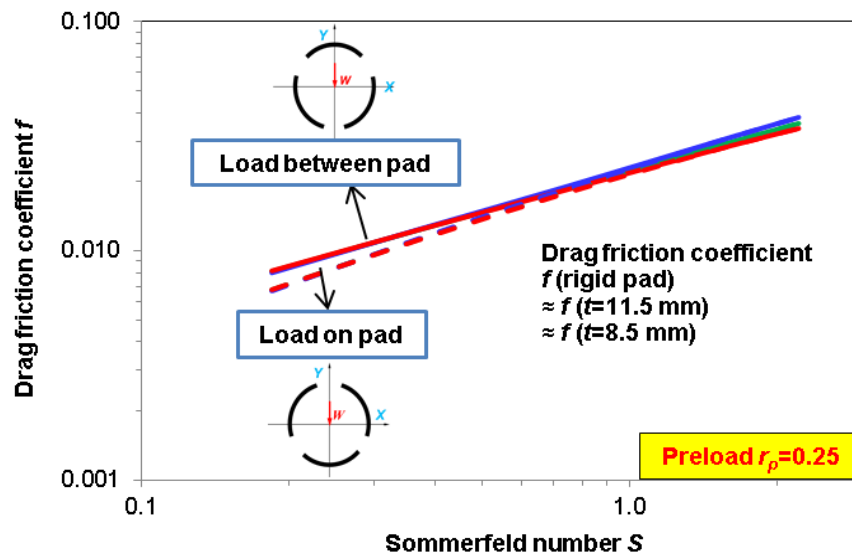
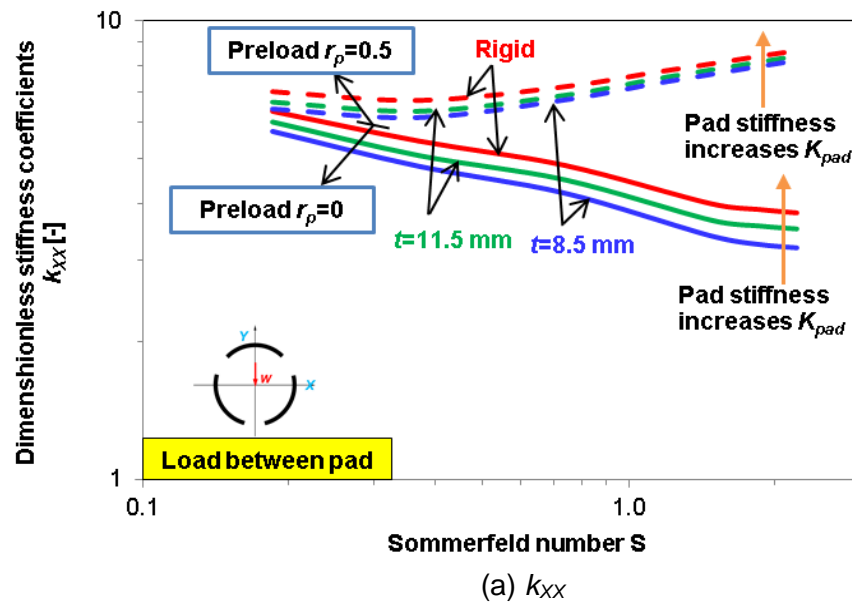
(a) LBP, $\bar{r}_p = 0$ & 0.5 (b) LOP & LBP, $\bar{r}_p = 0.25$

Figure 16 Three-pad TPJB drag friction coefficient (f) vs. Sommerfeld number. Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload varies: LBP and LOP configurations. Specific load $W/(LD)=689 \text{ kPa}$, rotor speed $\Omega=500 \text{ rpm}$ to $12,000 \text{ rpm}$.

Figures 17 and 18 show pad flexibility reduces the stiffness coefficients (k_{xx}, k_{yy}). For the LBP bearing, pad flexibility increases slightly k_{yy} by 3% at $S > 1$. Recall that Fig.

11 depicts an increase in bearing stiffnesses due to pad flexibility at a smaller unit load. For the LOP TPJB, the stiffness along the unloaded direction (k_{xx}) is one order of magnitude lesser than the stiffness along the load direction (k_{yy}) at a Sommerfeld number less than 1 ($S < 1$). For both the TPJBs under LBP and LOP configurations, pad flexibility slightly increases k_{yy} at a large Sommerfeld number ($S > 0.8$).

For the LBP TPJB with 0.5 preload and the LOP TPJB with 0.25 preload, k_{xx} and k_{yy} increase with S . For the LPB TPJB with null preload, k_{xx} and k_{yy} decrease with an increase in S .



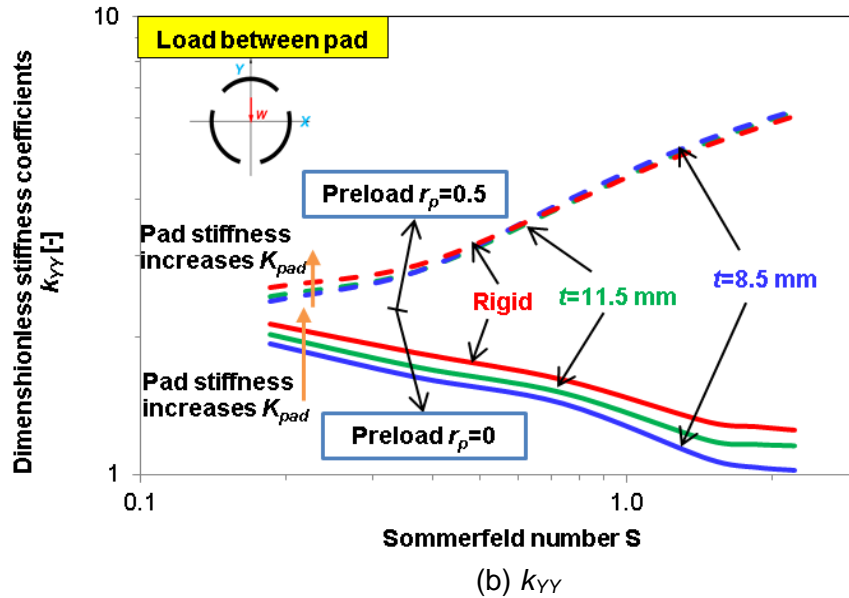
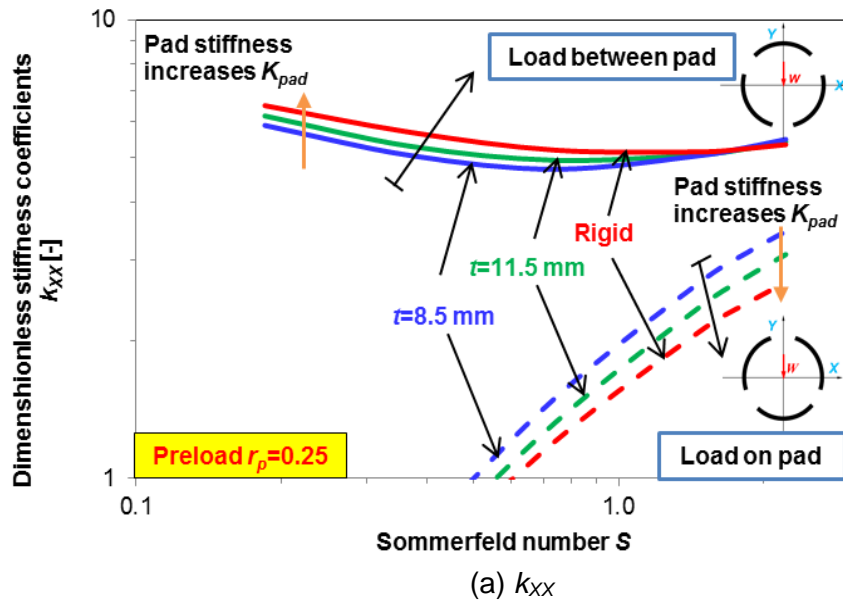


Figure 17 Three-pad TPJB stiffness coefficients (k_{xx} , k_{yy}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0, 0.5$: LBP configuration. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.



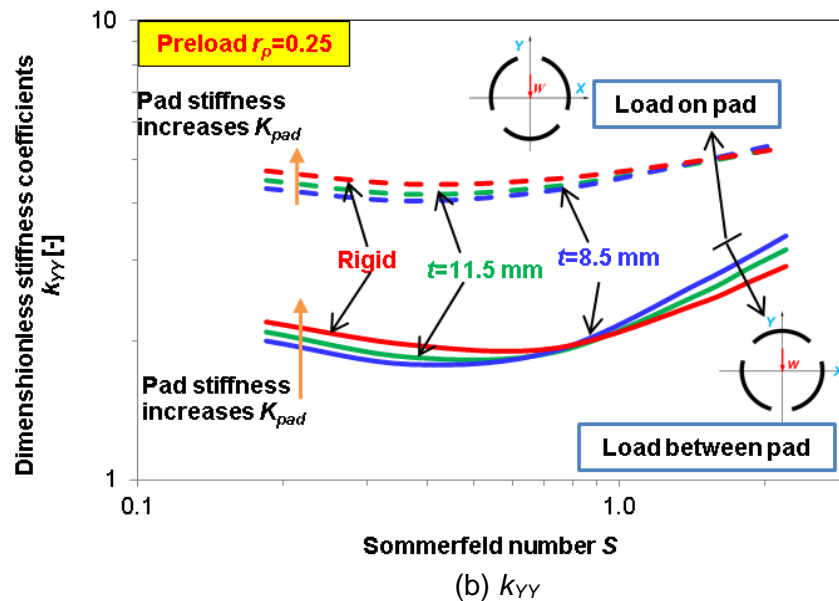


Figure 18 Three-pad TPJB stiffness coefficients (k_{xx} , k_{yy}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0.25$: LOP and LBP configurations. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.

Figures 19 and 20 show the damping coefficients (c_{xx} , c_{yy}) versus Sommerfeld number (S) for both LBP and LOP bearings and with pad preload $\bar{r}_p = 0, 0.25, 0.5$. The dimensionless damping coefficients increase with S ; however, for the LBP TPJB with null preload, c_{yy} decreases on $0.74 < S < 1.08$. As pad flexibility increases, the damping coefficients decrease dramatically by up to 33%, in particular for large Sommerfeld number ($S > 1.0$). Pad flexibility has a more pronounced effect on the damping coefficients of a LBP TPJB with null pad preload ($\bar{r}_p = 0$).

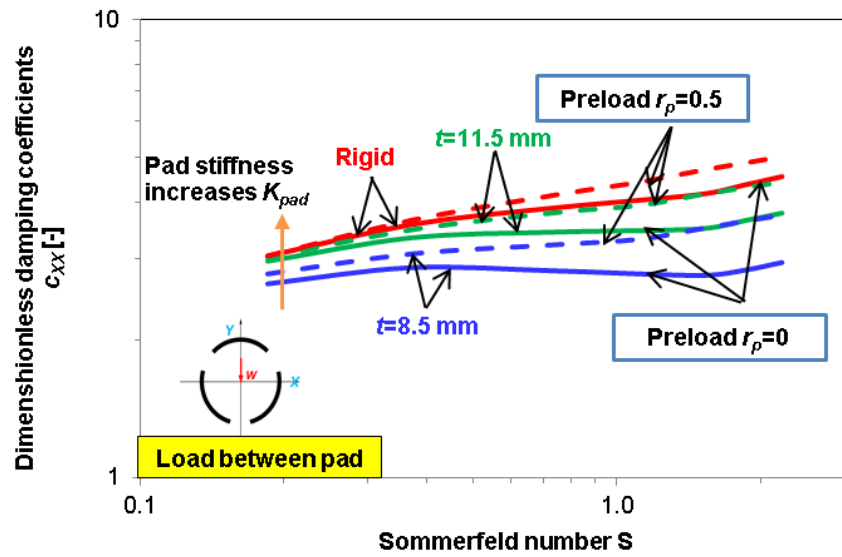
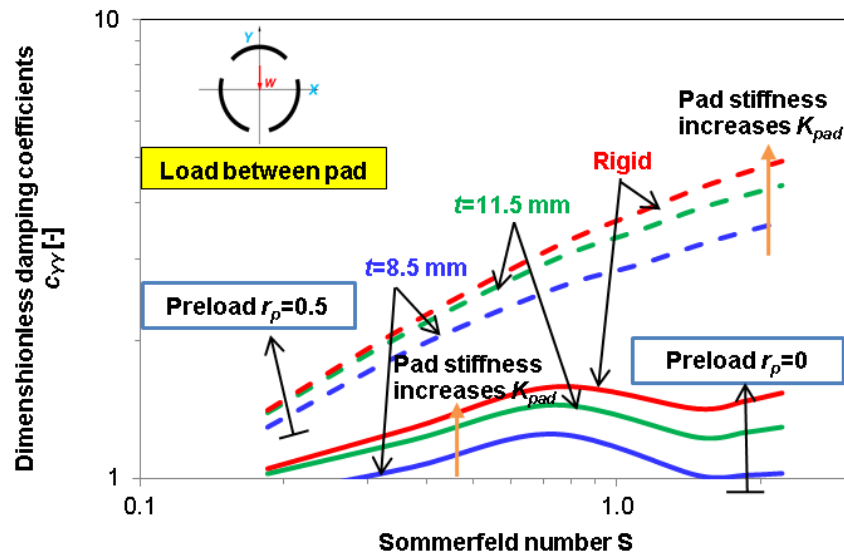
(a) c_{XX} (b) c_{YY}

Figure 19 Three-pad TPJB damping coefficients (c_{XX} , c_{YY}) vs. Sommerfeld number. Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0, 0.5$: LBP configuration. Specific load $W/(LD)=689$ kPa, rotor speed $\Omega=500$ rpm to 12,000 rpm.

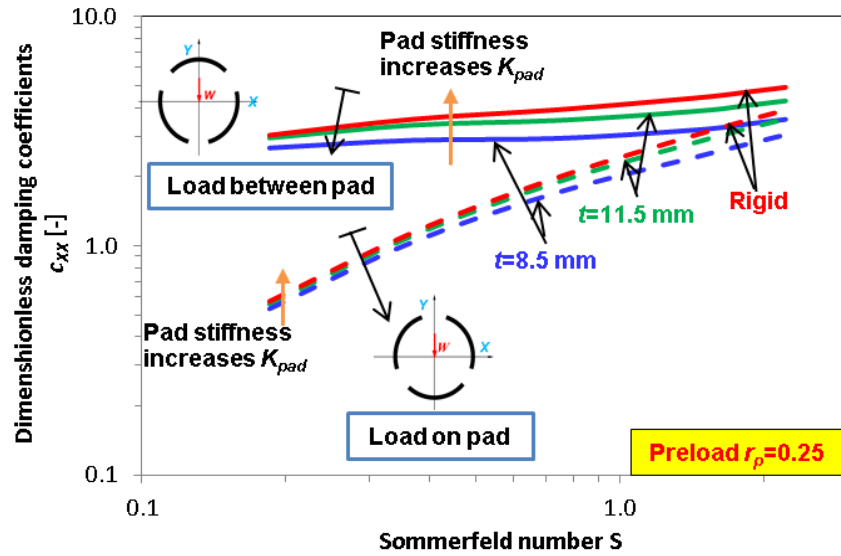
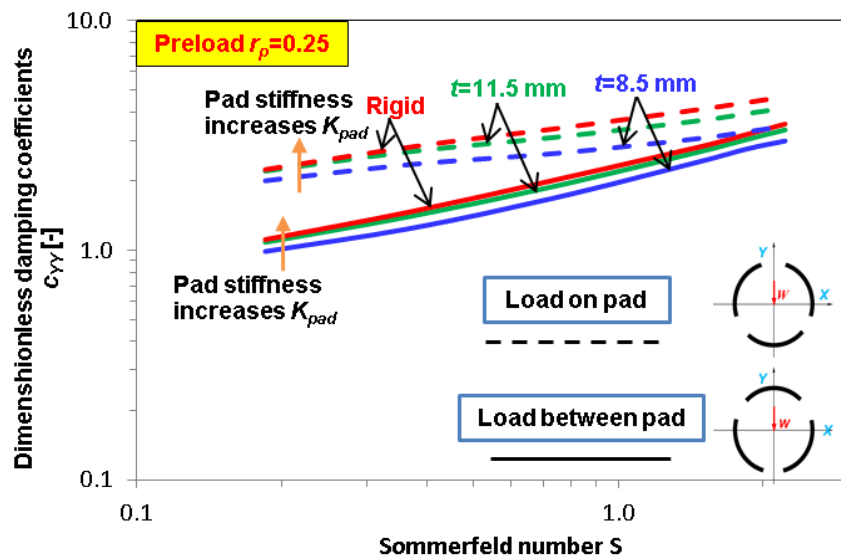
(a) c_{xx} (b) c_{yy}

Figure 20 Three-pad TPJB damping coefficients (c_{xx} , c_{yy}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0.25$: LOP and LBP configurations. Specific load $W/(LD) = 689$ kPa, rotor speed $\Omega = 500$ rpm to 12,000 rpm.

Figures 21 and 22 depict the virtual mass coefficients (m) versus Sommerfeld number (S). In general $m < 0$ denotes the bearing will stiffen as the excitation frequency increases. Pad flexibility has a more pronounced effect on m_{yy} , in particular for the LBP

bearing with preload equal to 0 and 0.25. Though the coefficients (m) do not approach zero as S increases, the virtual mass coefficients (M_{YY} , M_{YY}) approach to zero at a large S , thus indicating the dynamic stiffness $\text{Re}(Z_{YY})$ is frequency independent.

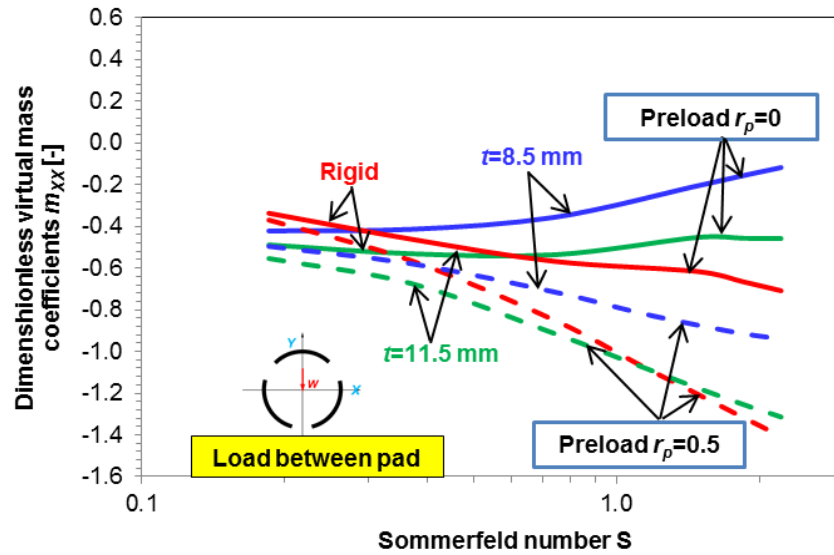
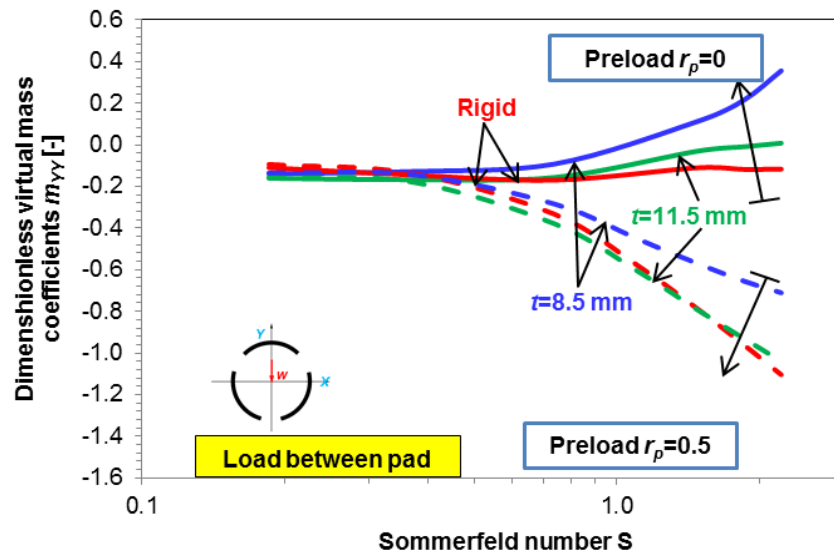
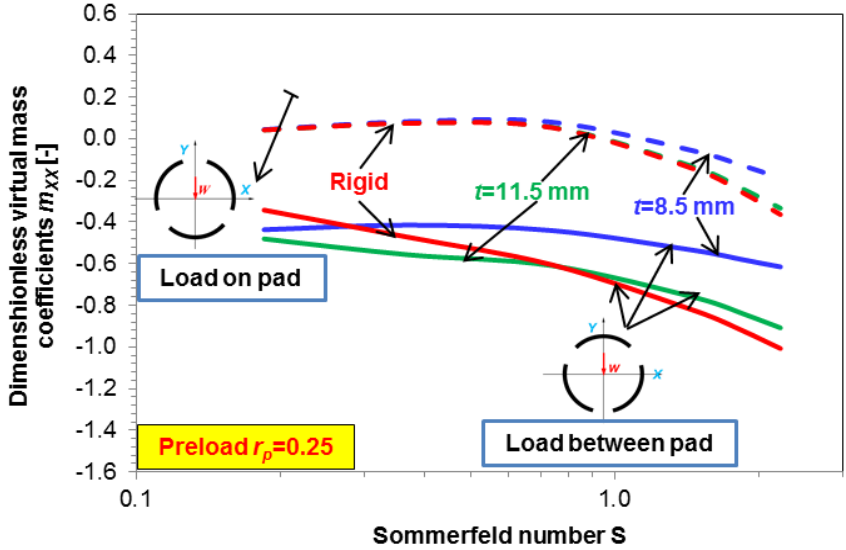
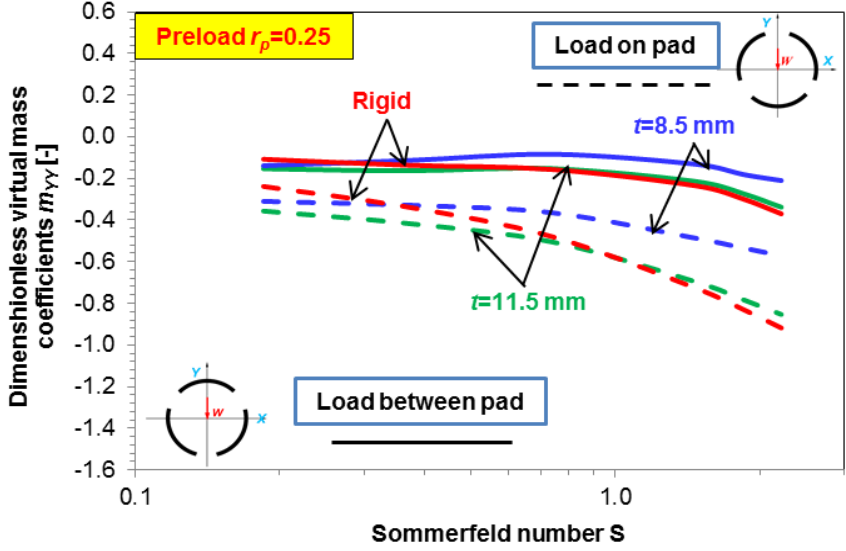
(a) m_{XX} (b) m_{YY}

Figure 21 Three-pad TPJB virtual mass coefficients (m_{XX} , m_{YY}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0, 0.5$: LBP configuration. Specific load $W/(LD) = 689$ kPa, rotor speed $\Omega = 500$ rpm to 12,000 rpm.



(a) m_{xx}



(b) m_{yy}

Figure 22 Three-pad TPJB damping coefficients (m_{xx} , m_{yy}) vs. Sommerfeld number (S). Pad stiffness $k_{pad} = 3.15, 7.33, \infty$ (rigid) and $k_{piv} = 16$. Pad preload $\bar{r}_p = 0.25$: LOP and LBP configurations. Specific load $W(LD) = 689$ kPa, rotor speed $\Omega = 500$ rpm to 12,000 rpm.

Guide to get a pad structural stiffness matrix in ANSYS

Yingkun Li

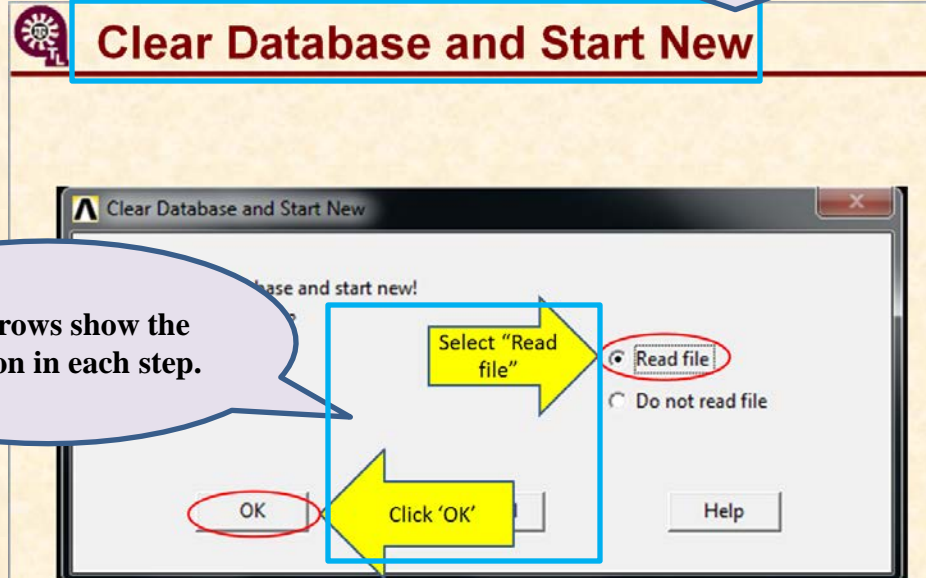
Research Assistant

Dr. Luis San Andrés

Mast-Childs Professor

Those slides give the detail of each step to get the stiffness matrix from ANSYS©.

The title of each slide shows the name of the dialog box shown in the slide



The arrows show the operation in each step.

This guide uses a pad model shown on the right to help to illustrate the operations.

The table below lists the geometry and material properties.

Table Pad geometry and material parameters

Inner radius (R_i)	0.050921m
Outer radius (R_o)	0.076321m
Axial length of the pad (l)	0.0559m
Angular amplitude of pad (β)	58.9°
Pivot position	0.5
Young's modulus (E)	2.0×10^{11} Pa
Poisson's ratio	0.3

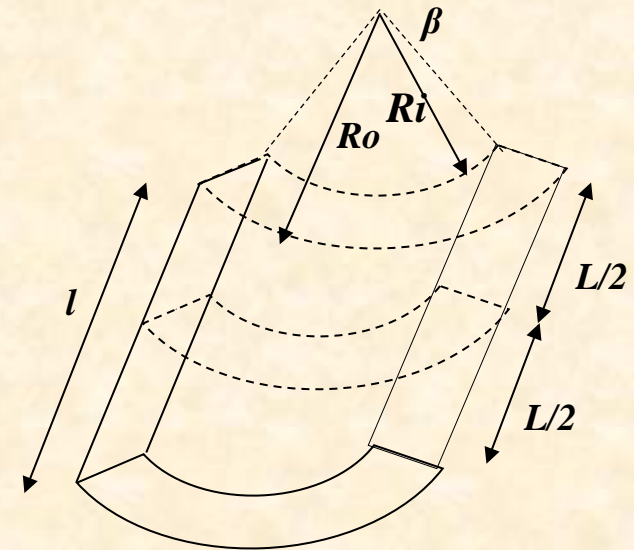
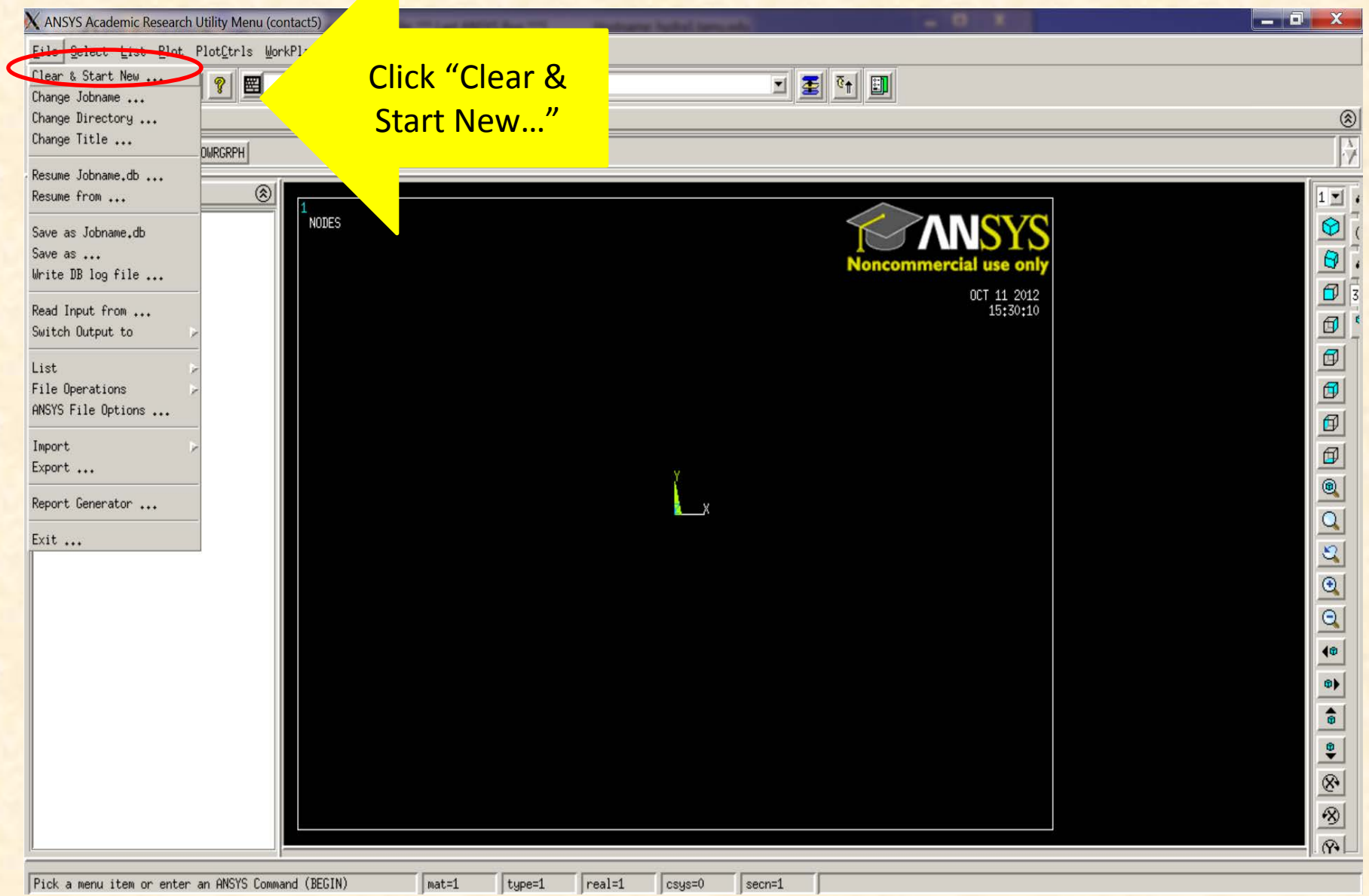


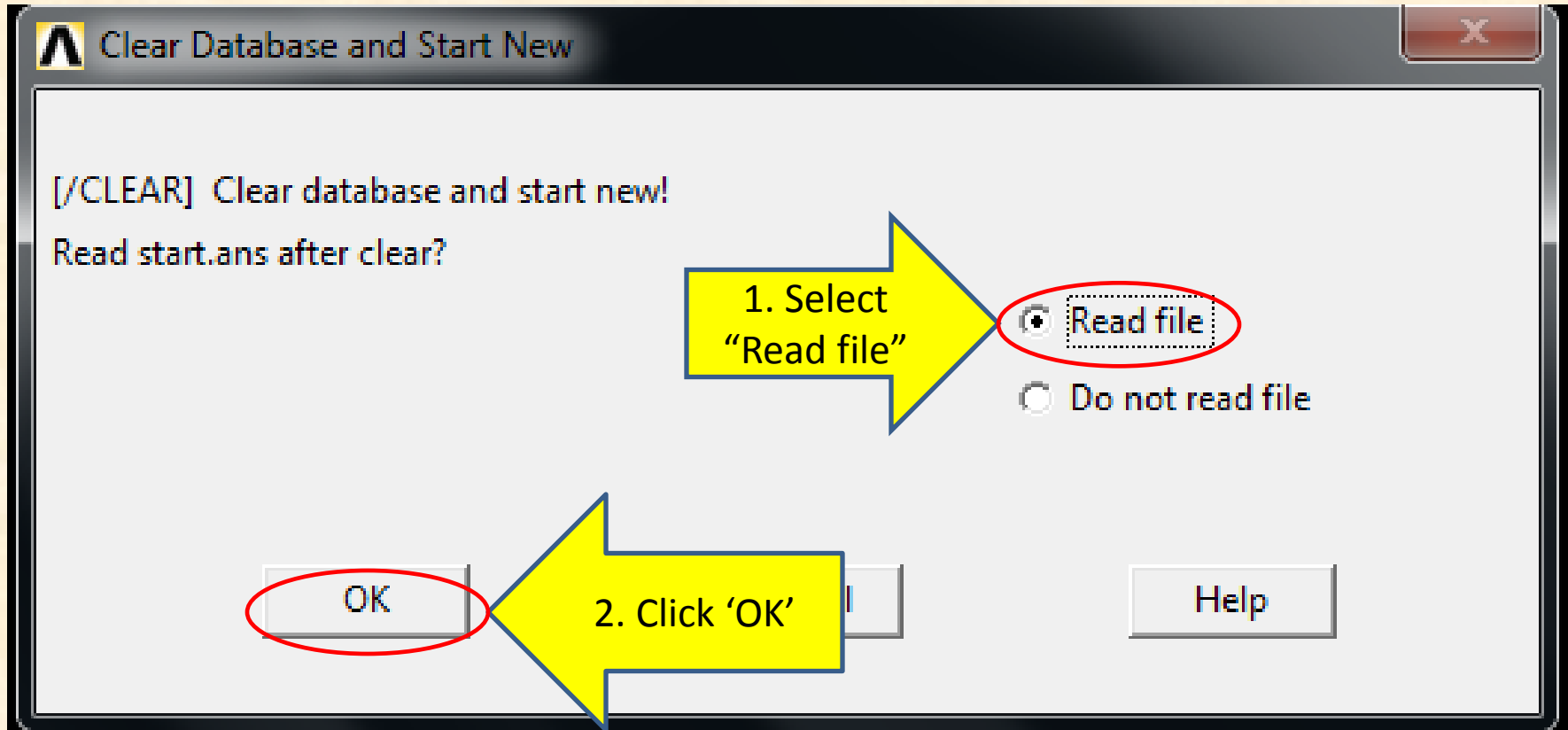
Figure The shape of the pad model

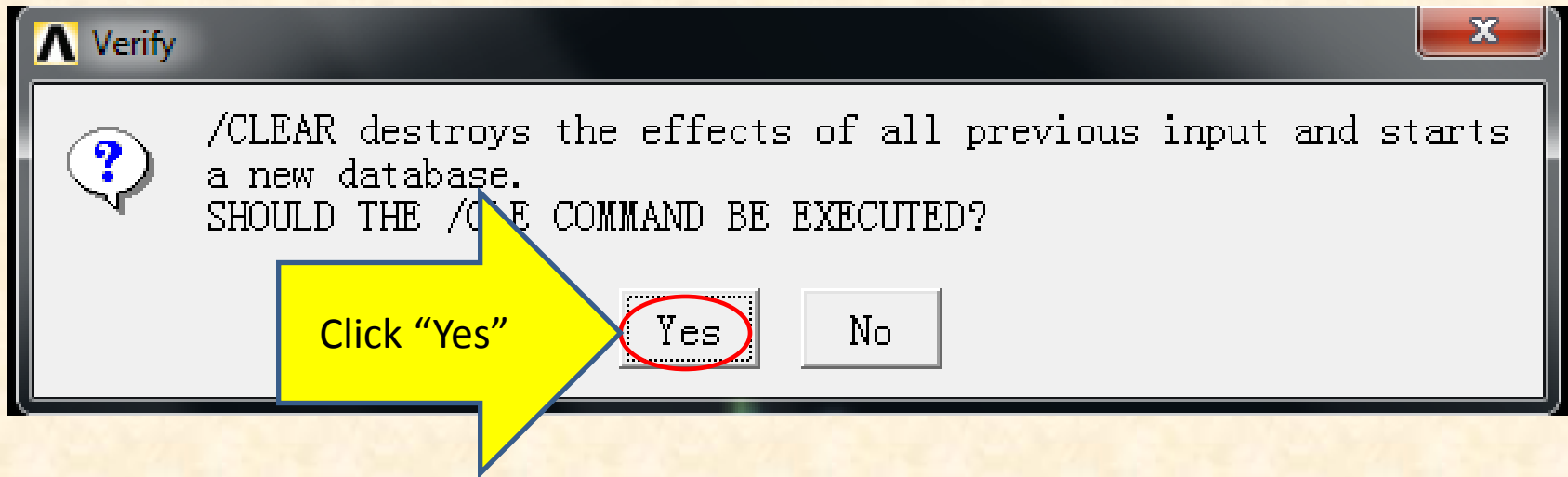
Main interface



Click "Clear & Start New..."

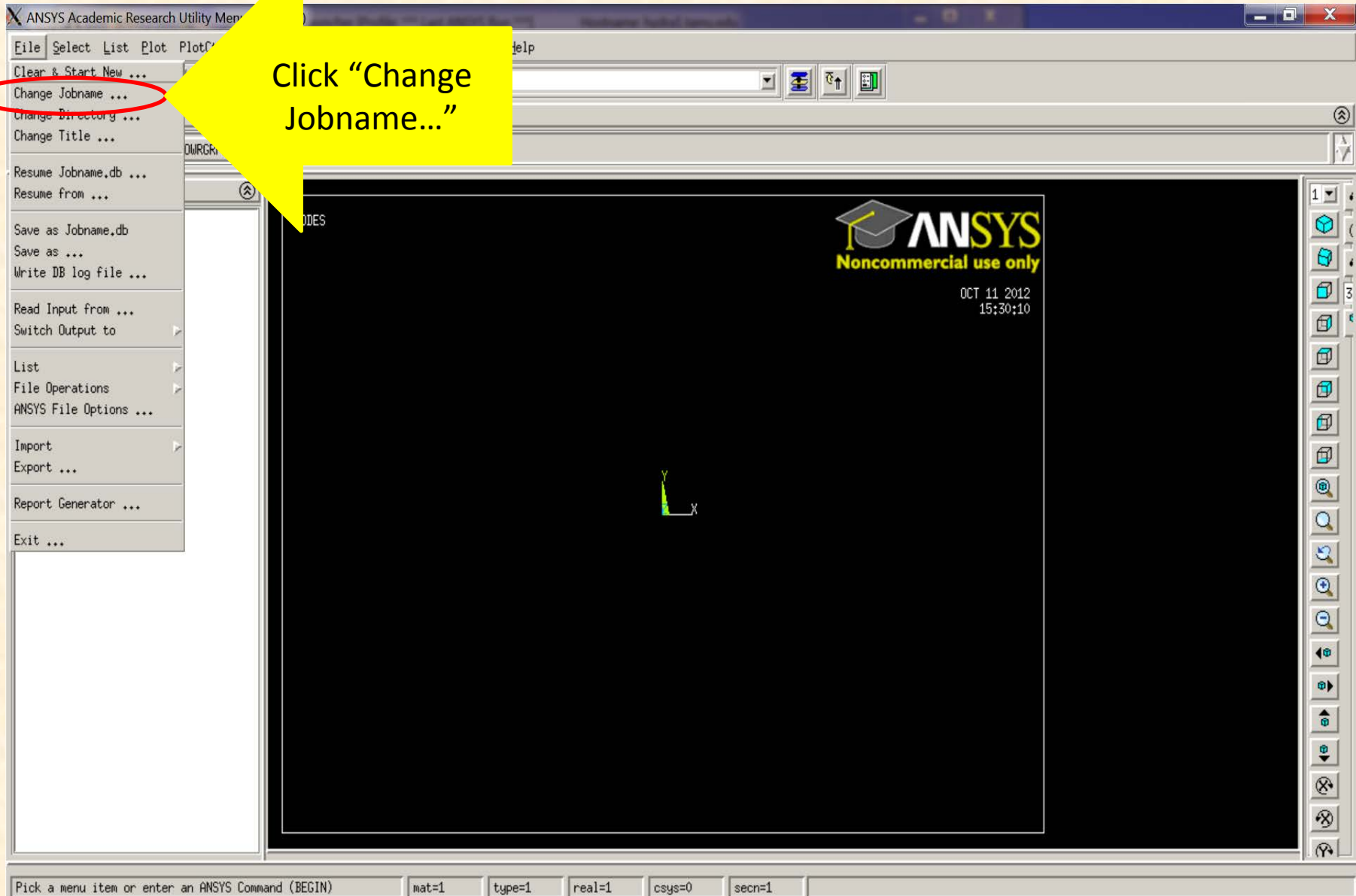
Clear Database and Start New



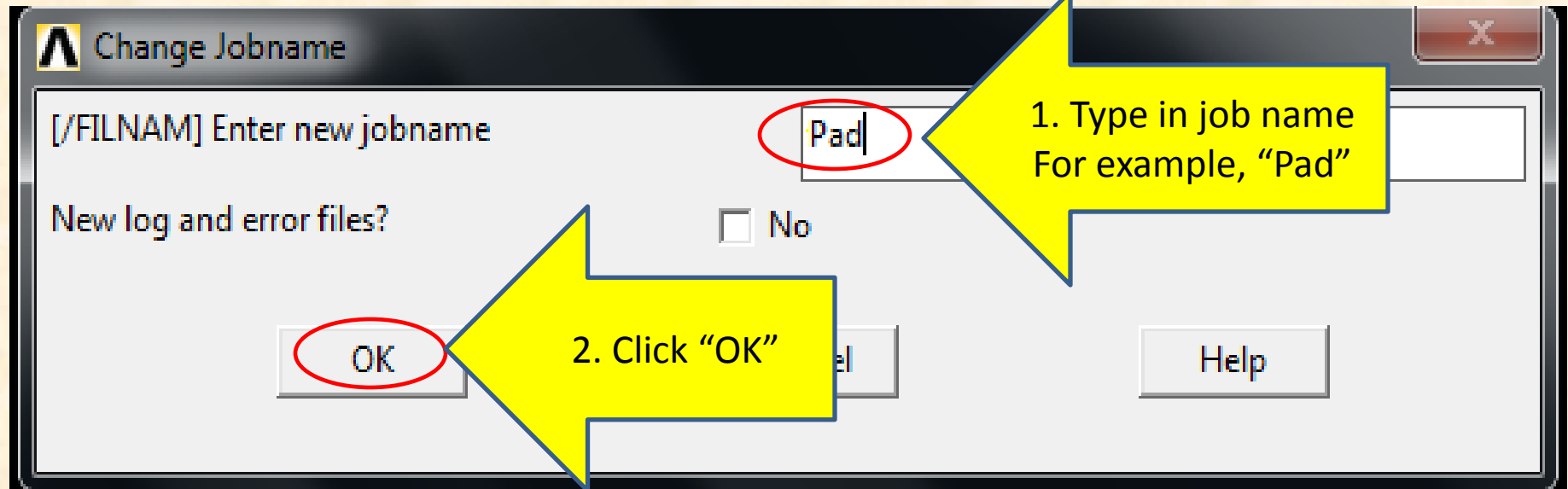


Main interface

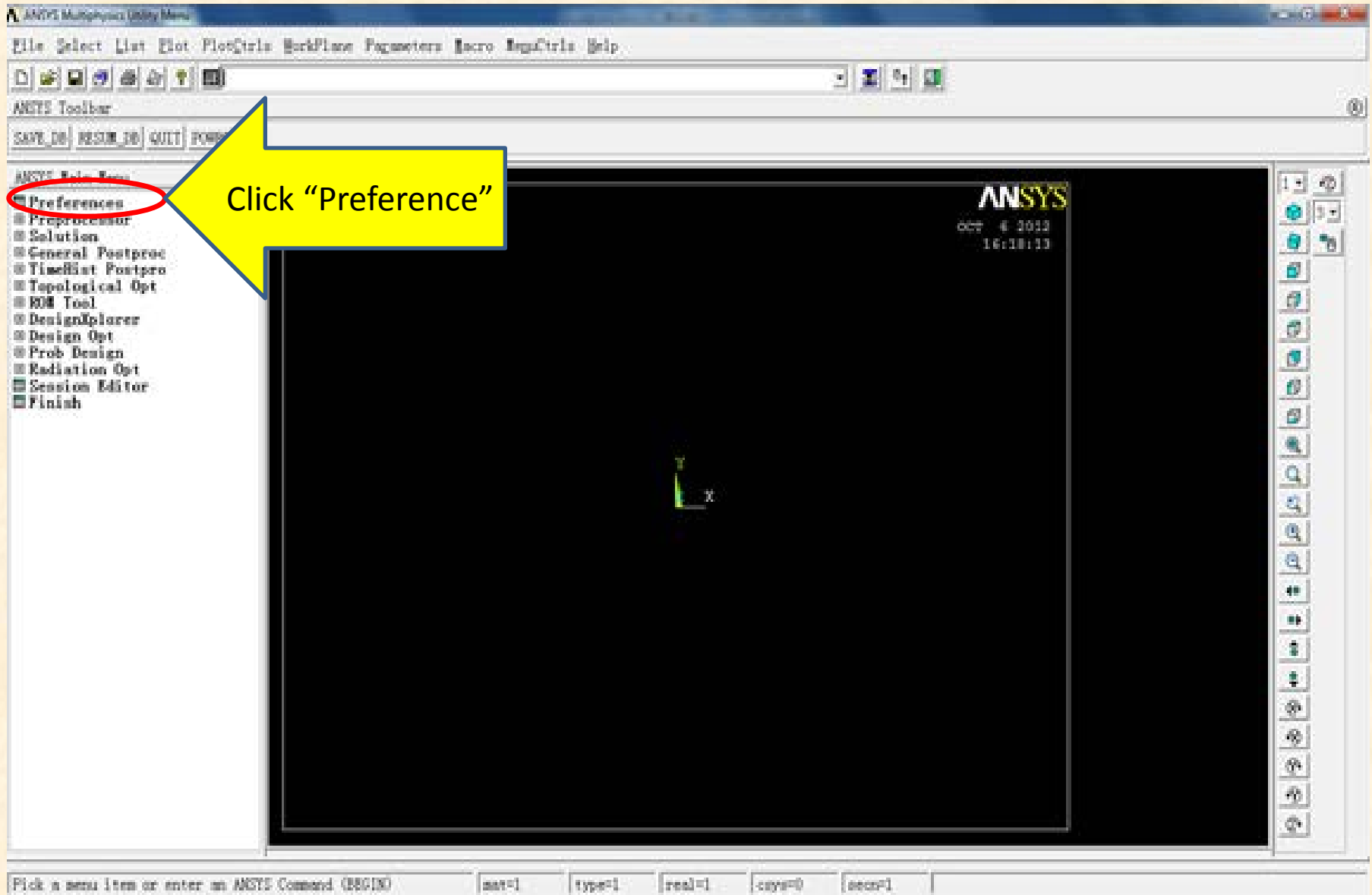
Click "Change Jobname..."



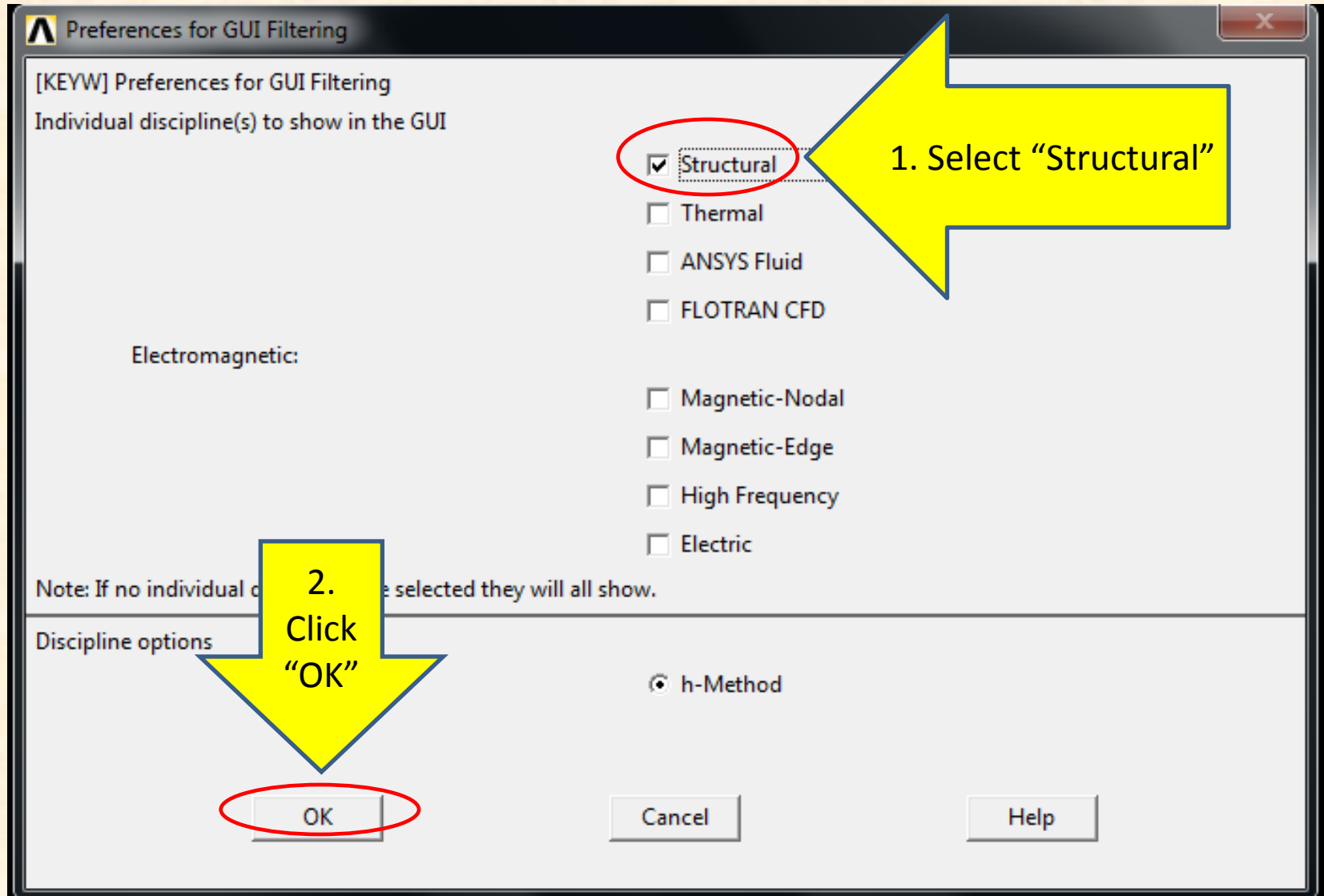
Change jobname



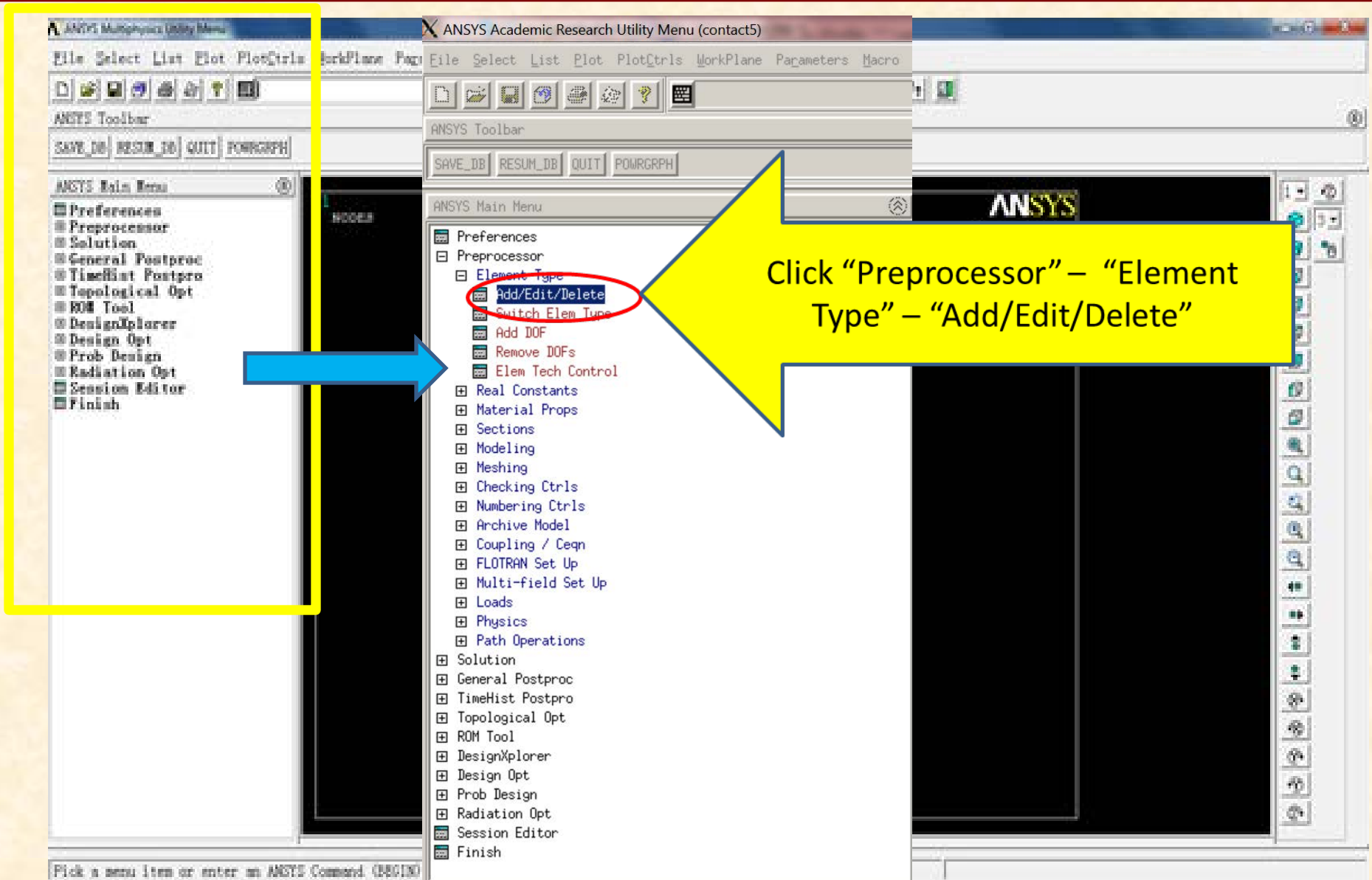
Main interface



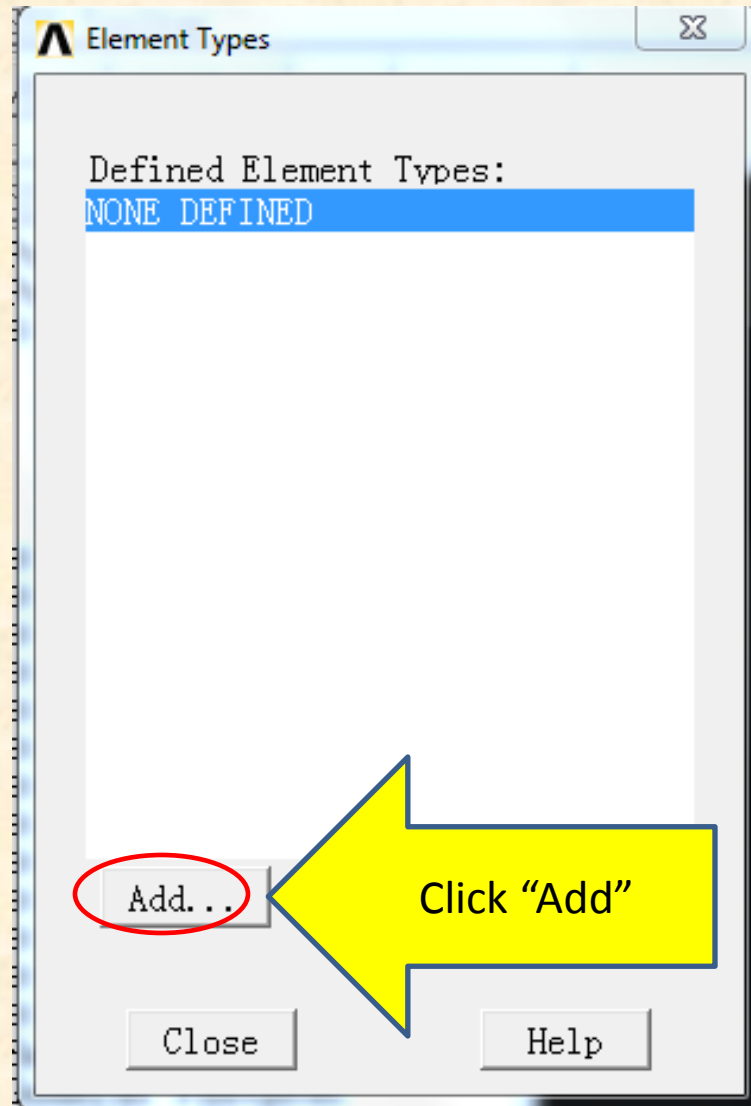
Preference for GUI Filtering



Main interface



Element Types



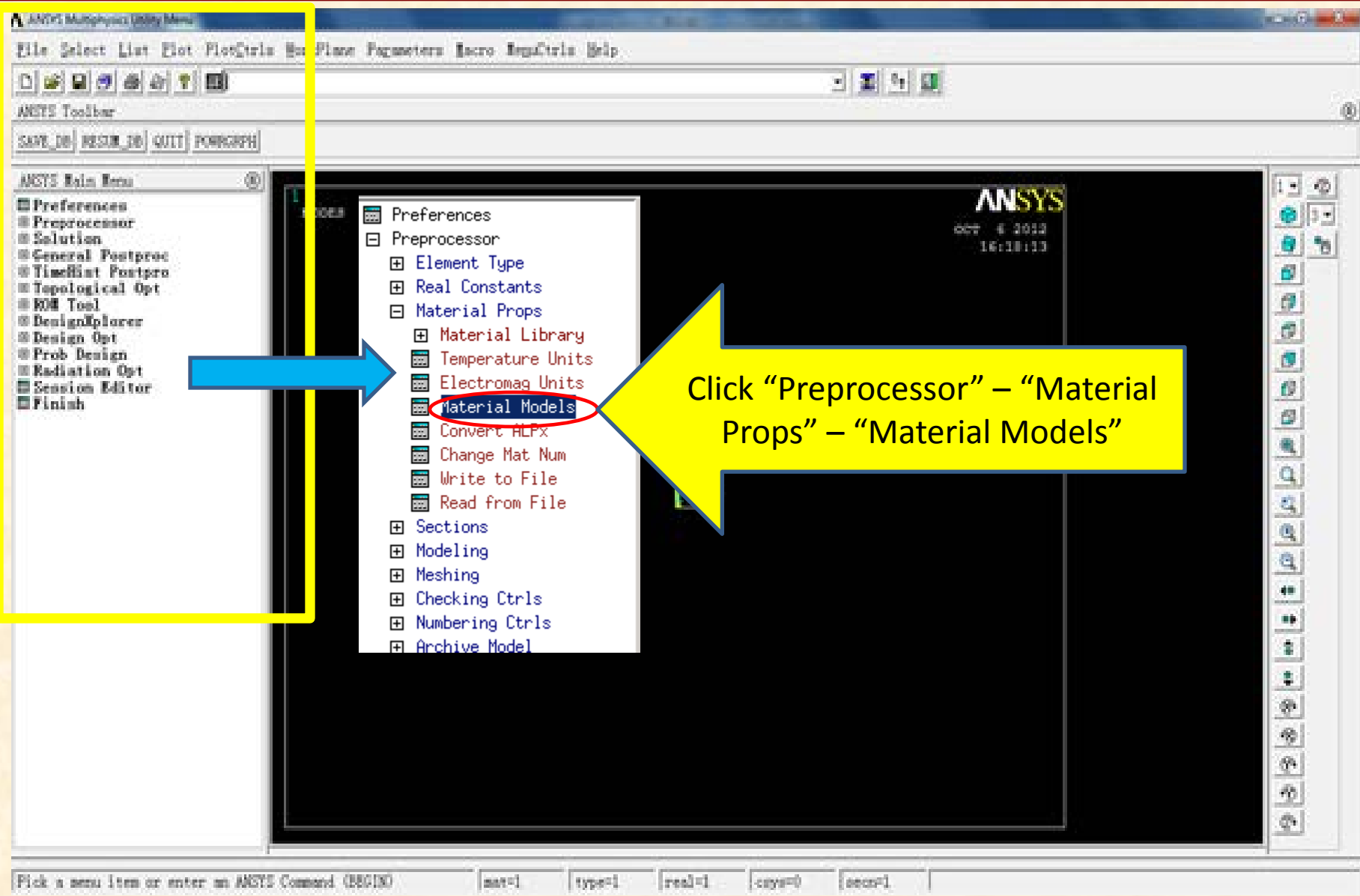
Element Types

The screenshot shows the 'Library of Element Types' dialog box. The title bar reads 'Library of Element Types'. Below the title bar, it says 'Only structural element types are shown' and 'Library of Element Types'. A list of element types is shown: Link, Beam, Pipe, Solid (highlighted in blue), Shell, and Solid-Shell. To the right of this list is another list of element types: 4 node 182, 8 node 183, Brick 8 node 185 (highlighted in blue), 20node 186, and concret 65. Below these lists is a text box containing 'Brick 8 node 185'. At the bottom left, there is a text box labeled 'Element type reference number' containing the value '1'. At the bottom right, there are three buttons: 'OK', 'Cancel', and 'Help'. The 'OK' button is circled in red. Two yellow arrows with black outlines provide instructions: one points to the 'Solid' option in the list, and the other points to the 'OK' button.

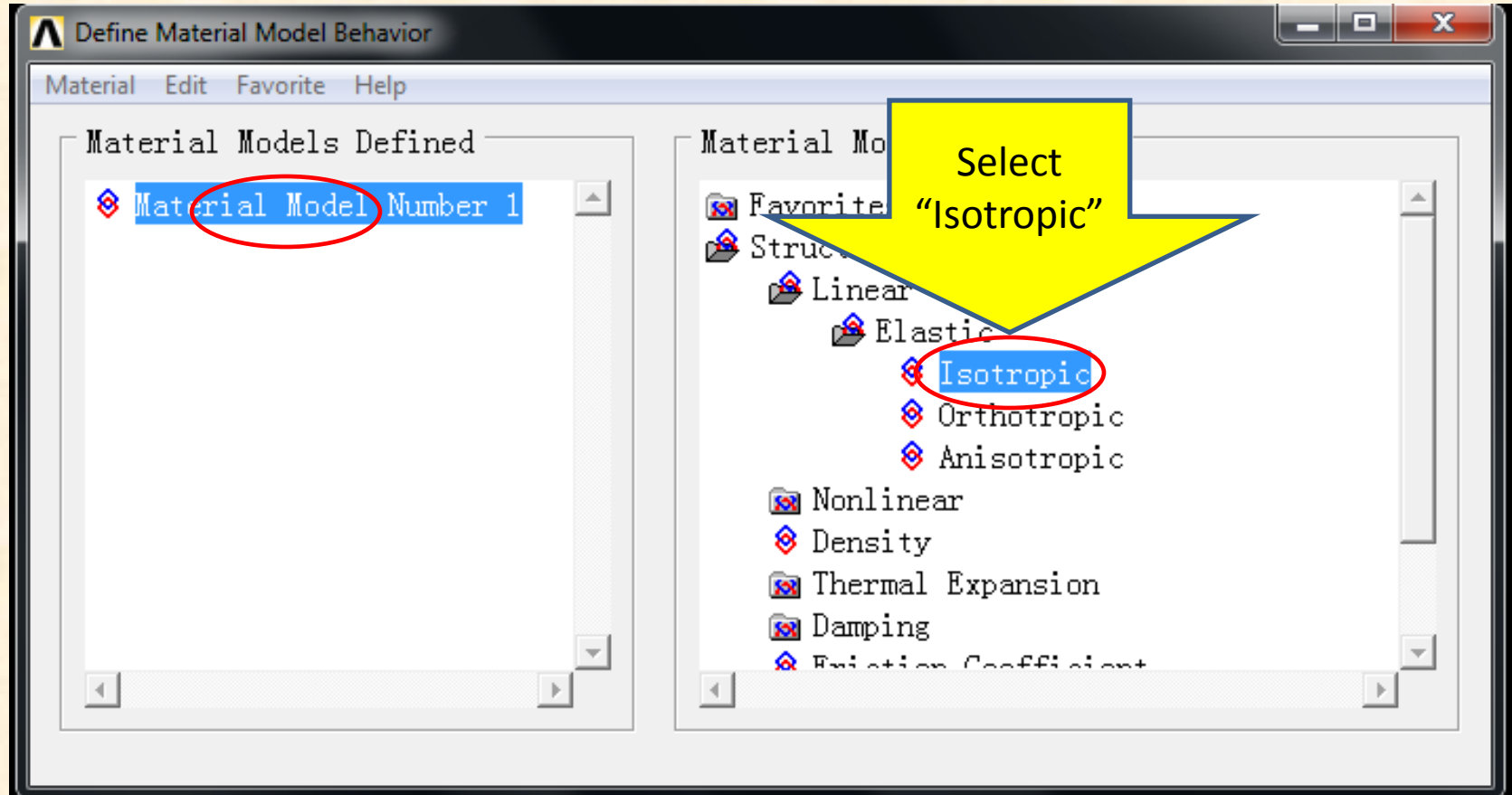
1. Select "Solid"
– "Brick 8 node 185"

2. Click "OK"

Main interface



Define Material Model Behavior



Linear Isotropic Properties

Linear Isotropic Properties for Material Number 1

Linear Isotropic Material Properties for Material Number 1

Temperatures	EY
	2.1e11
	0.3

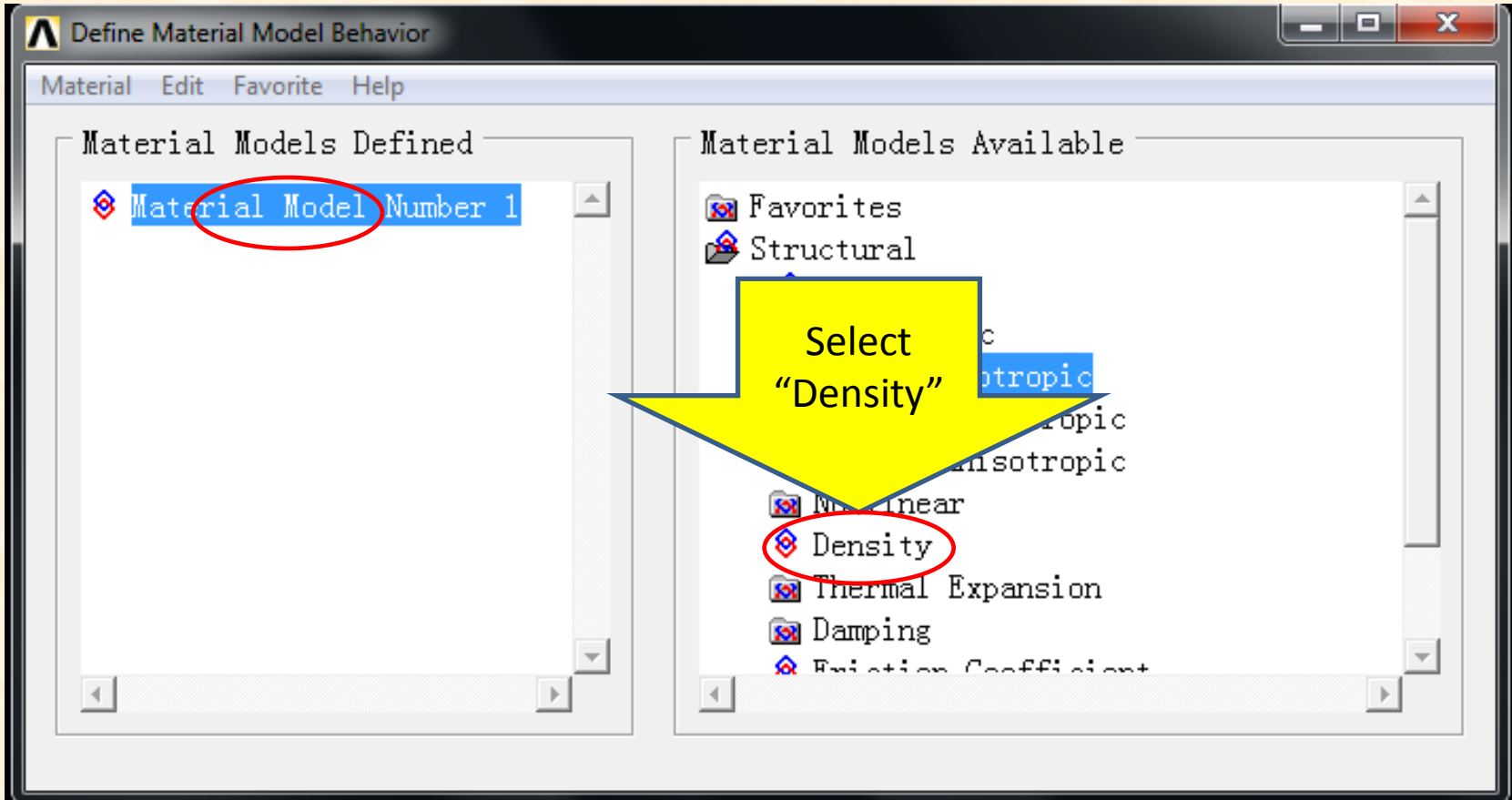
1. Type in Young's Modulus

2. Type in Poisson's ratio

3. Click "OK"

Add Temperature Delete Temperature Graph OK Cancel Help

Define Material Model Behavior



Density for Material Number 1

Density for Material Number 1

Temperatures	T1
DENS	7200

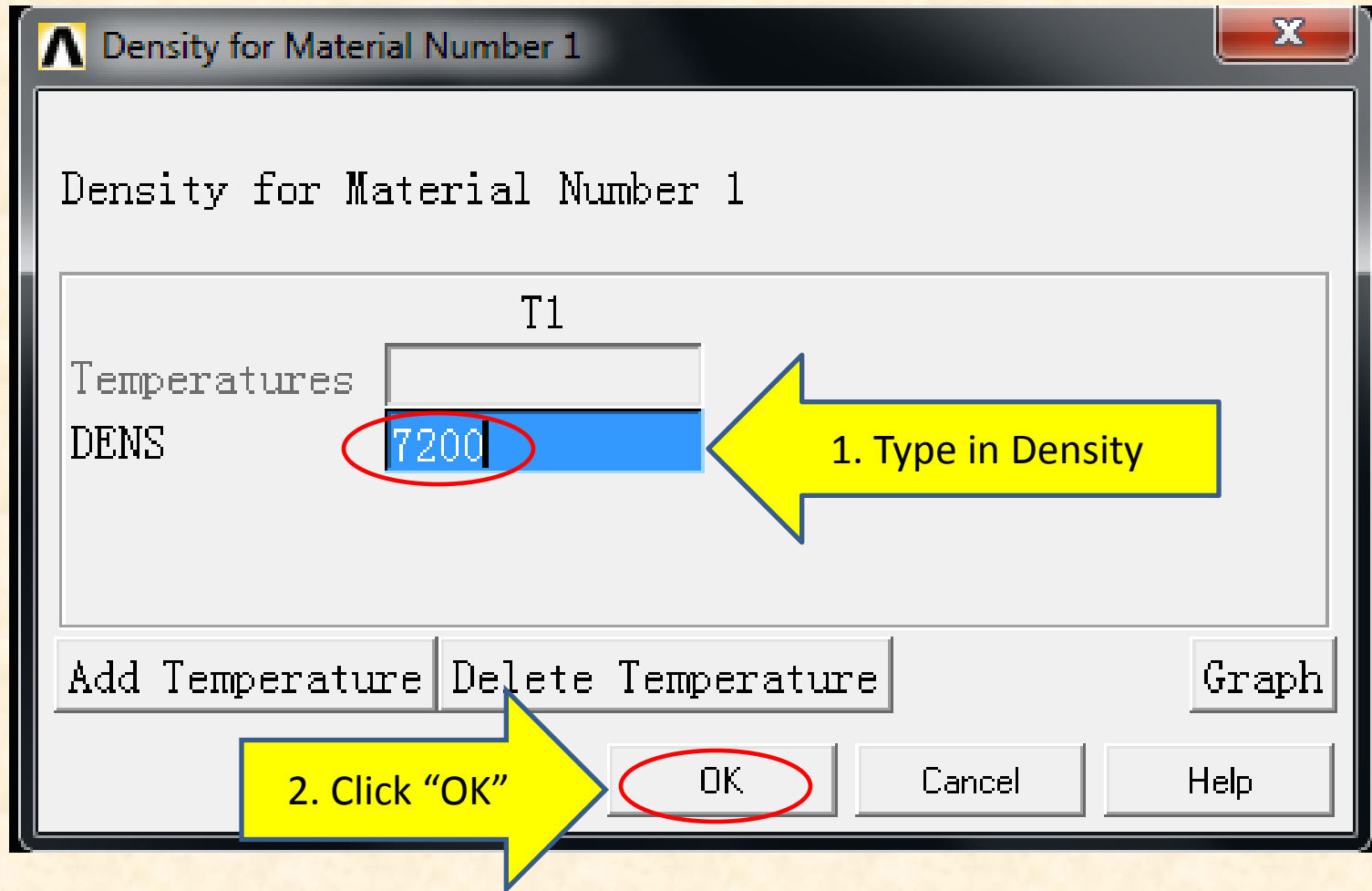
Add Temperature | Delete Temperature | Graph

OK | Cancel | Help

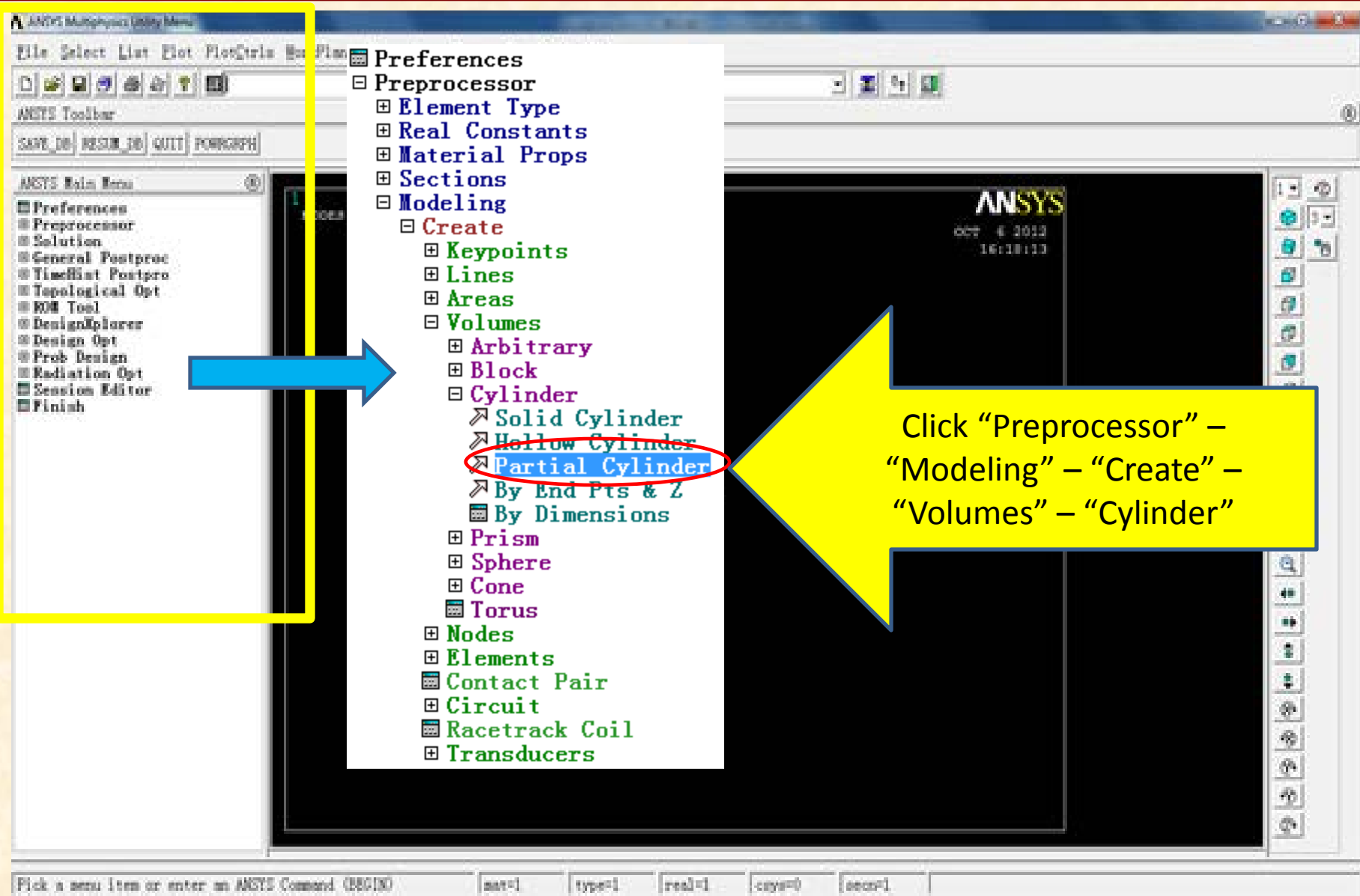
1. Type in Density

2. Click "OK"

Density for Material Number 1



Main interface



Partial Cylinder

<input checked="" type="radio"/> Pick <input type="radio"/> Unpick	
WP X =	
Y =	
Global X =	
Y =	
Z =	
WP X	0
WP Y	0
Rad-1	0.050921
Theta-1	140.55
Rad-2	0.076321
Theta-2	299.45
Depth	0.0559
<input type="button" value="OK"/> <input type="button" value="Apply"/> <input type="button" value="Reset"/> <input type="button" value="Cancel"/> <input type="button" value="Help"/>	

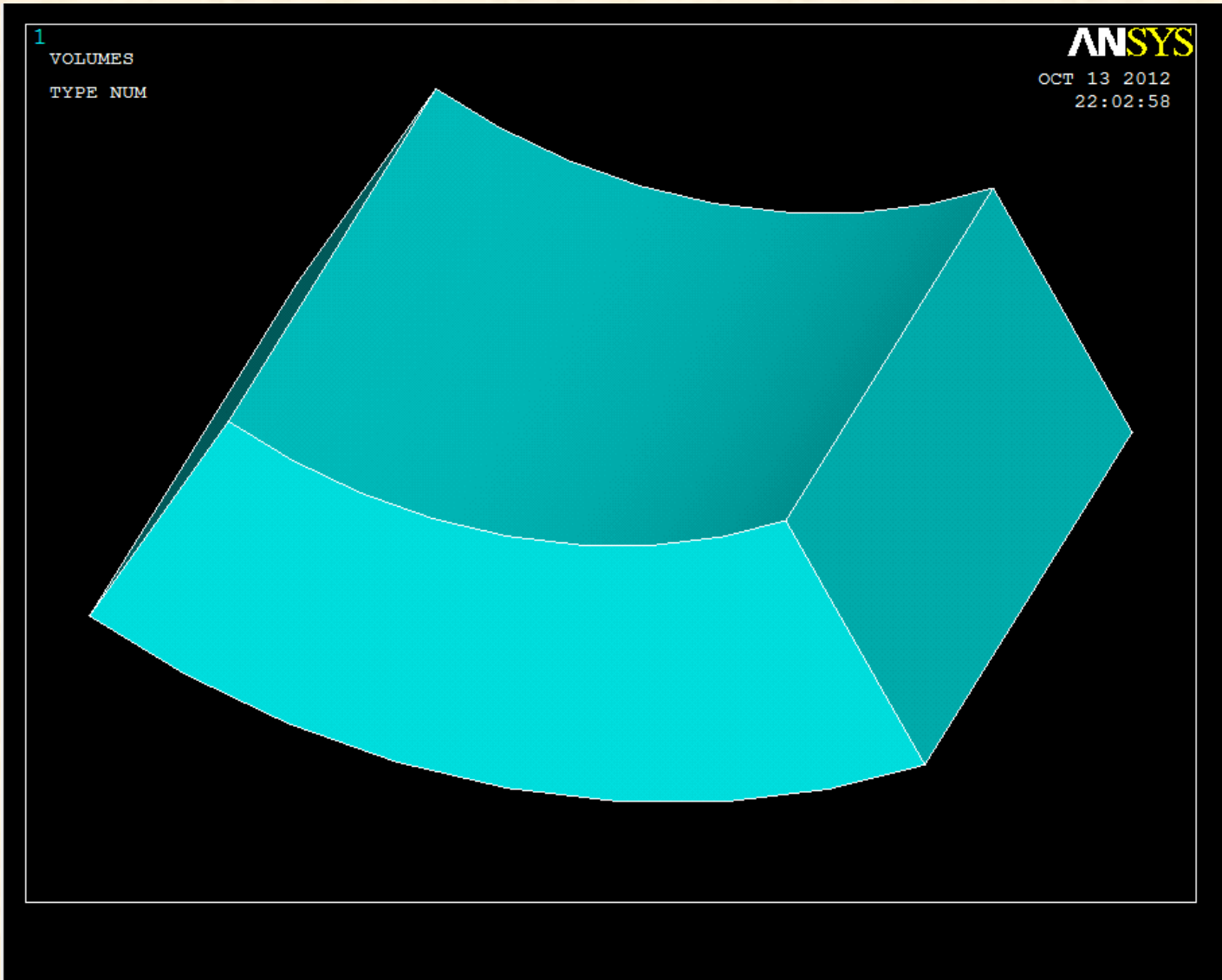
Type in the center coordinate of the bearing

Type in the leading edge angle and trailing edge angle of the pad

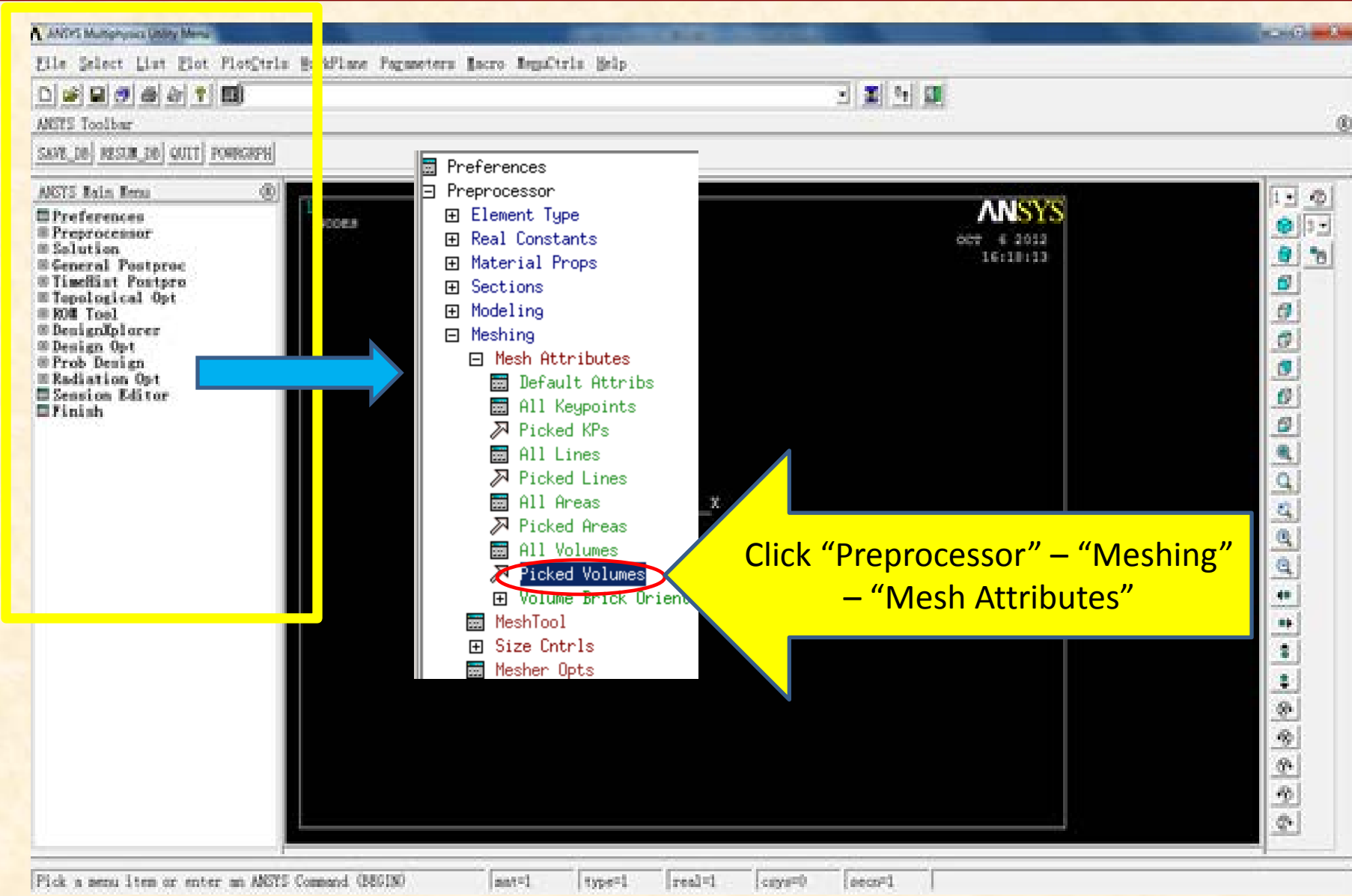
Type in inner and outer radius of the pad

Type in the axial length of the pad

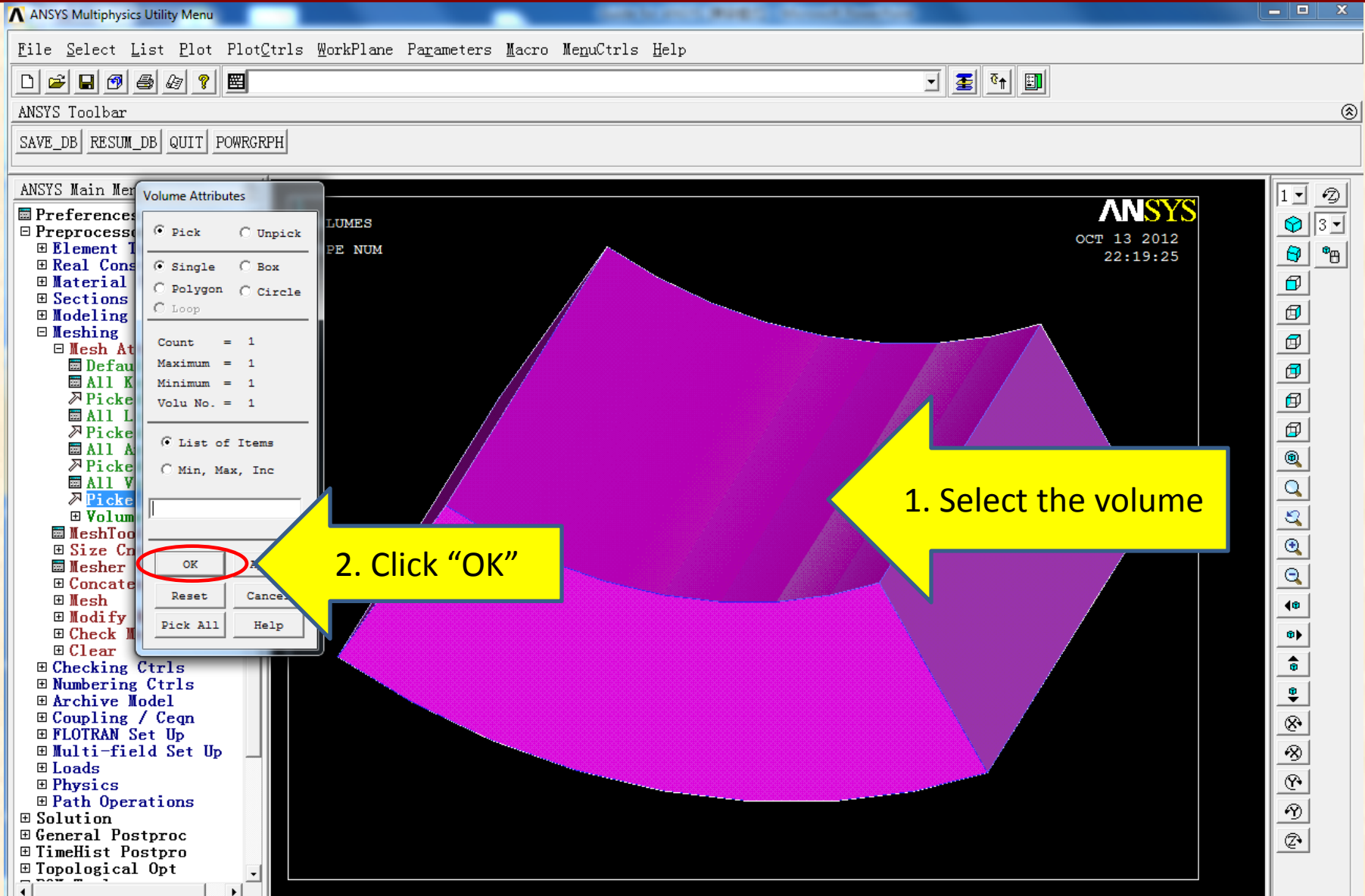
The example pad model



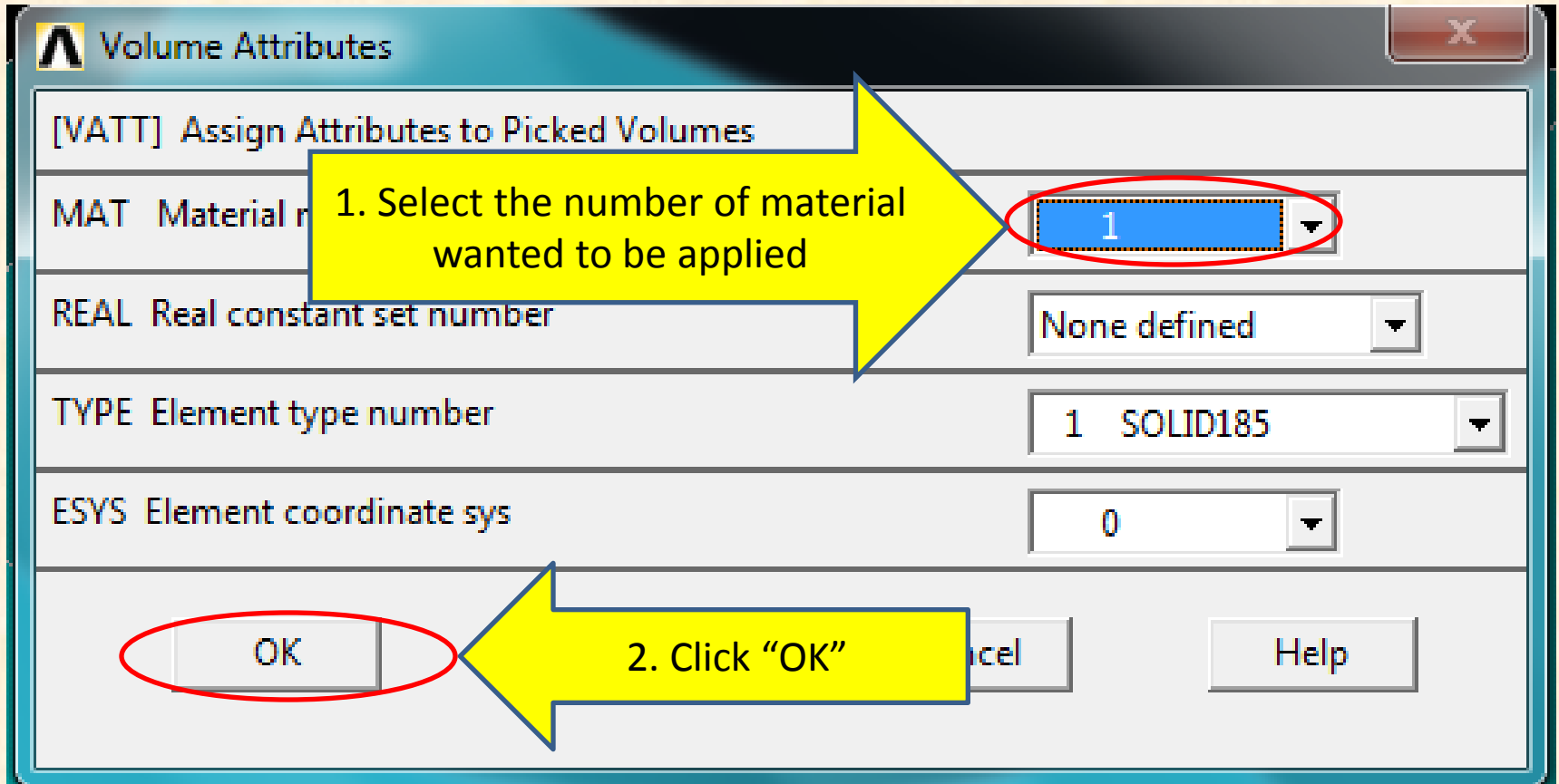
Main interface



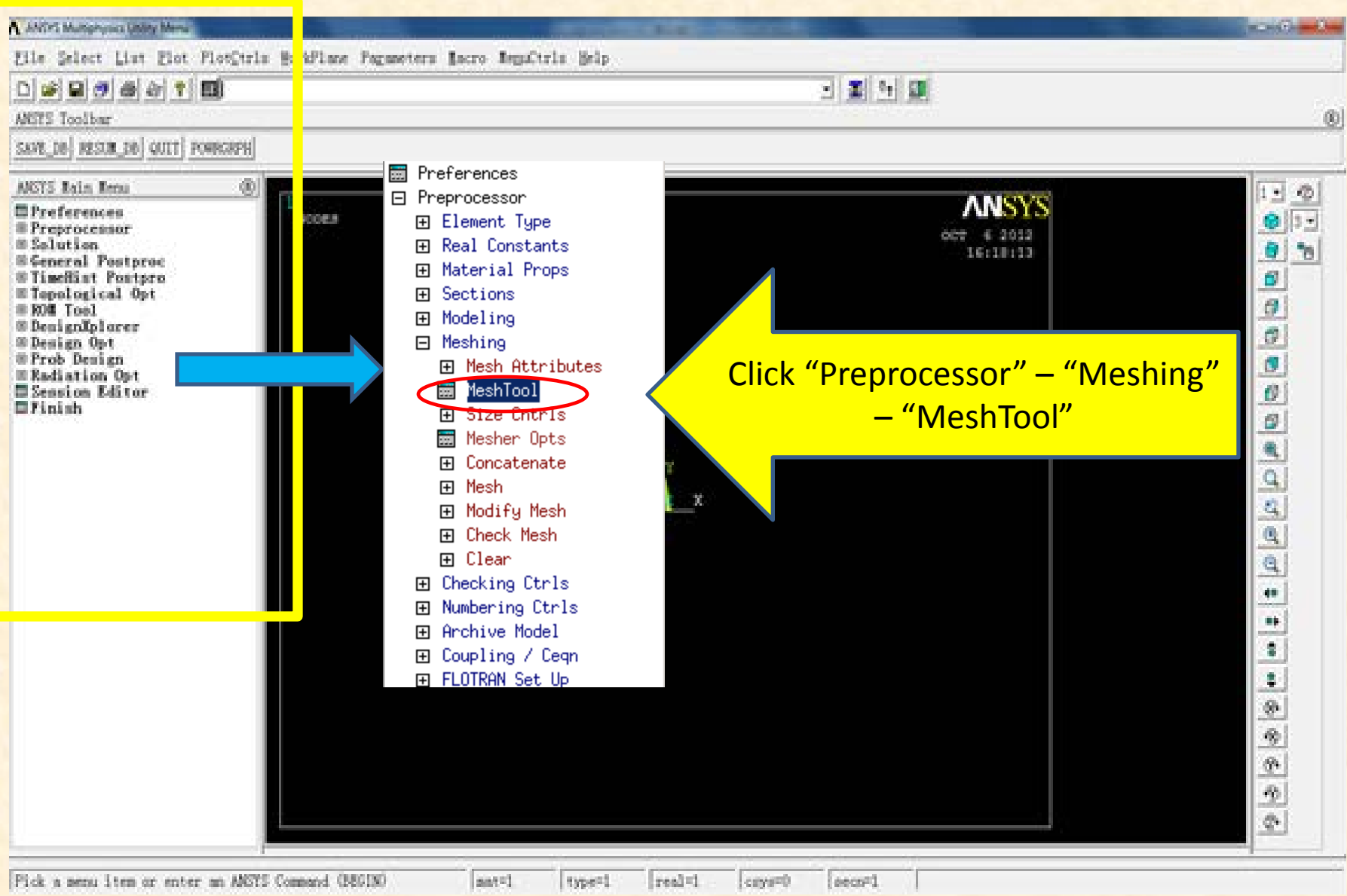
Volume Attributes

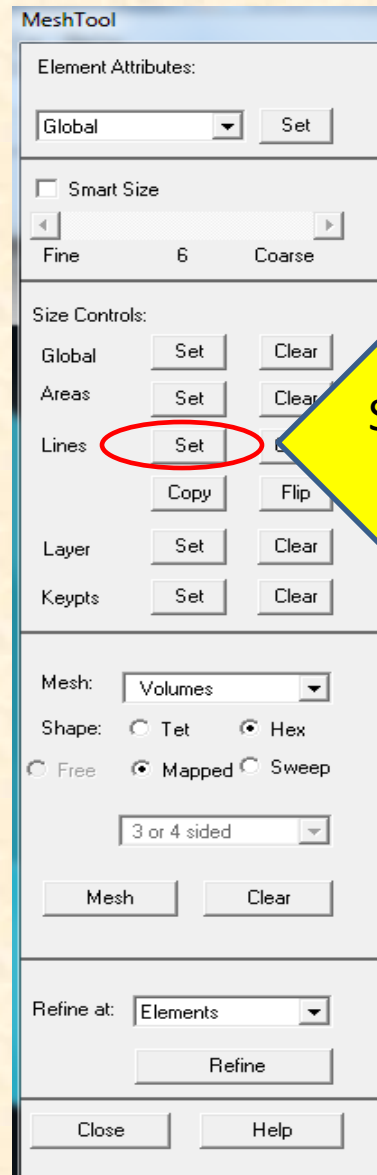


Volume Attributes



Main interface





Set element size on lines

MeshTool

1
VOLUMES
PE NUM

ANSYS
OCT 6 2012
23:28:50

Element Size on Picked Lines

Pick Unpick

Single Box

Polygon Circle

Loop

Count = 4
Maximum = 12
Minimum = 1
Line No. = 4

List of Items
 Min, Max, Inc

OK
Reset
Pick All
Cancel
Help

1. Select lines to define the size of elements on it

2. Click "OK"

stiffness_matrix_surface_1

MeshTool

Element Sizes on Picked Lines

[LESIZE] Element sizes on picked lines

SIZE Element edge length

NDIV No. of elements

(NDIV is used only if SIZE is blank or zero)

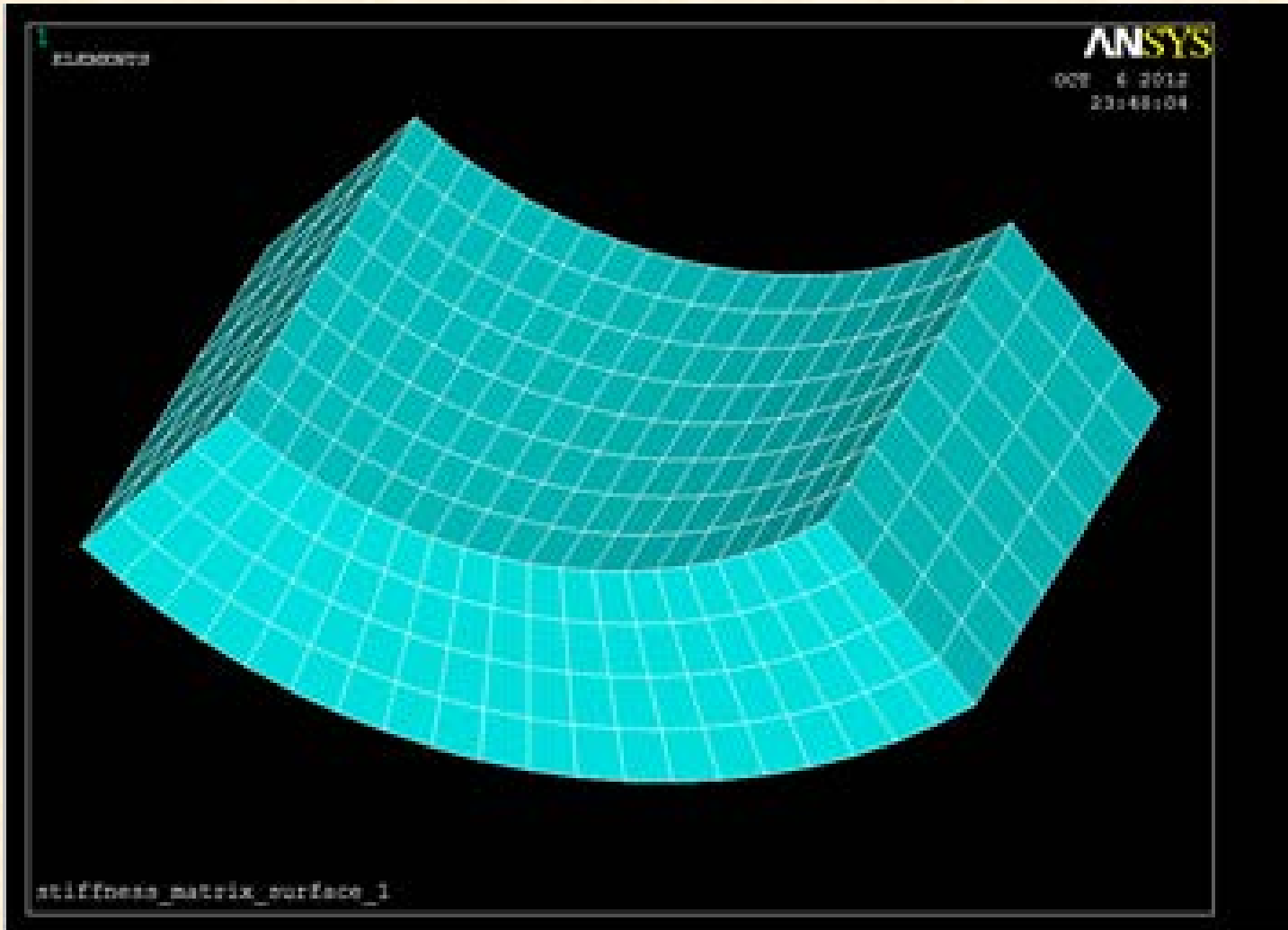
KYNDIV SIZE,NDIV can be changed Yes

SPACE Spacing ratio

ANGSIZ Division arc (degrees)

(use ANGSIZ only if number of divisions (NDIV) and element edge length (SIZE) are blank or zero)

Meshed model

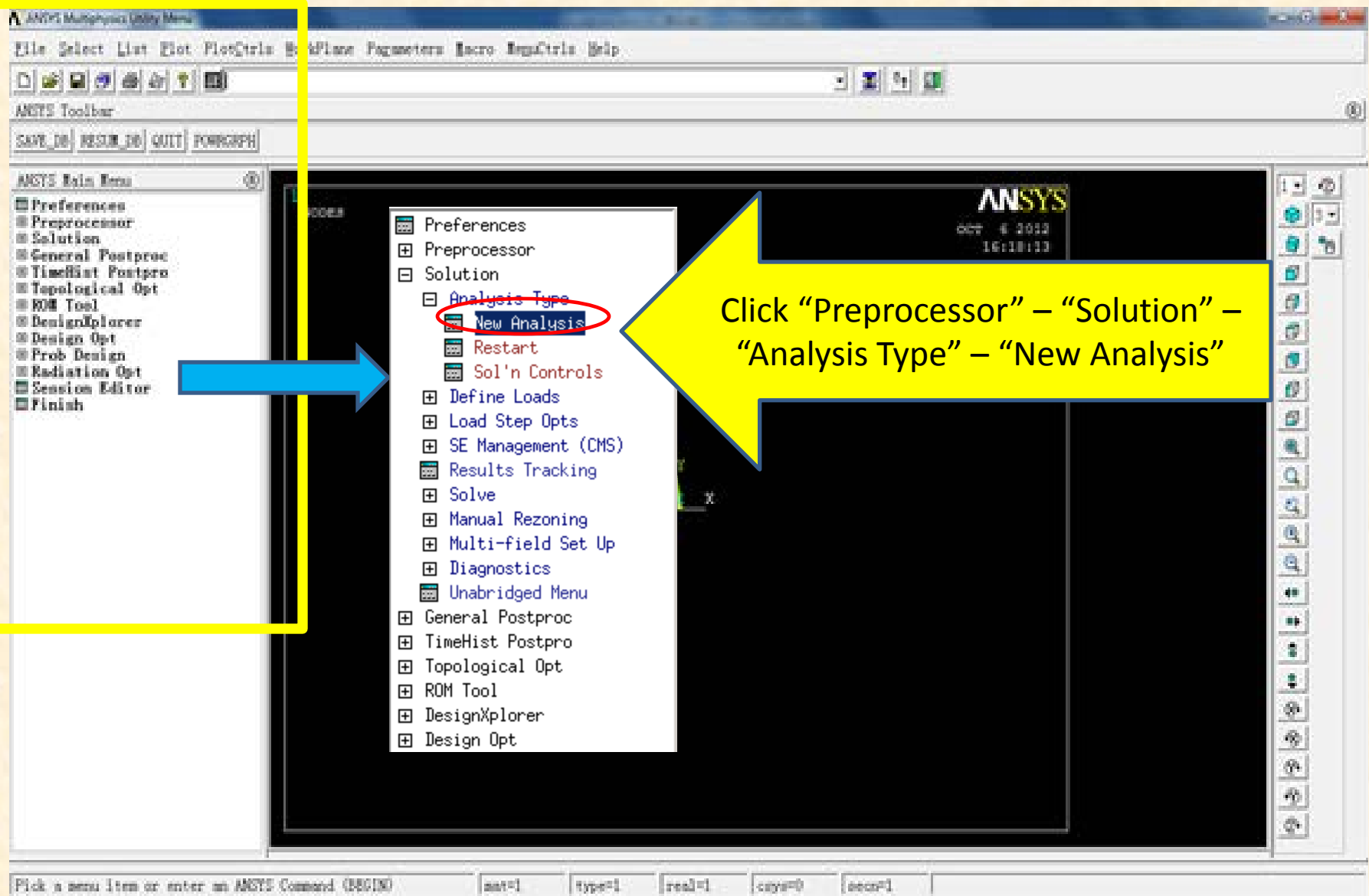


Get stiffness matrix

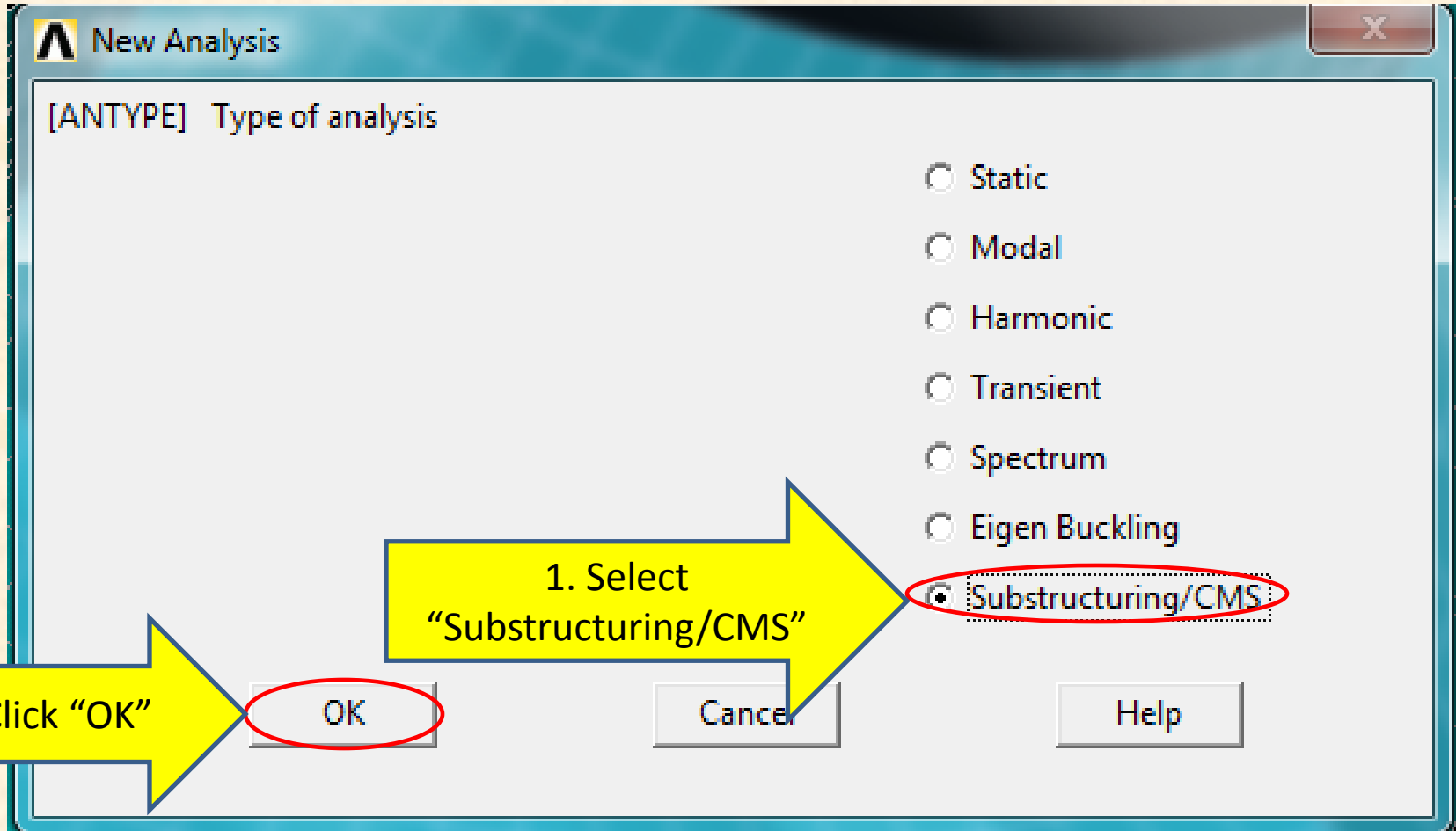
ANSYS© has a function to reduce the stiffness matrix of a structural model.

As not all the nodal displacements are important to the analysis, ANSYS © offers a way to reduce the stiffness matrix into the size desired.

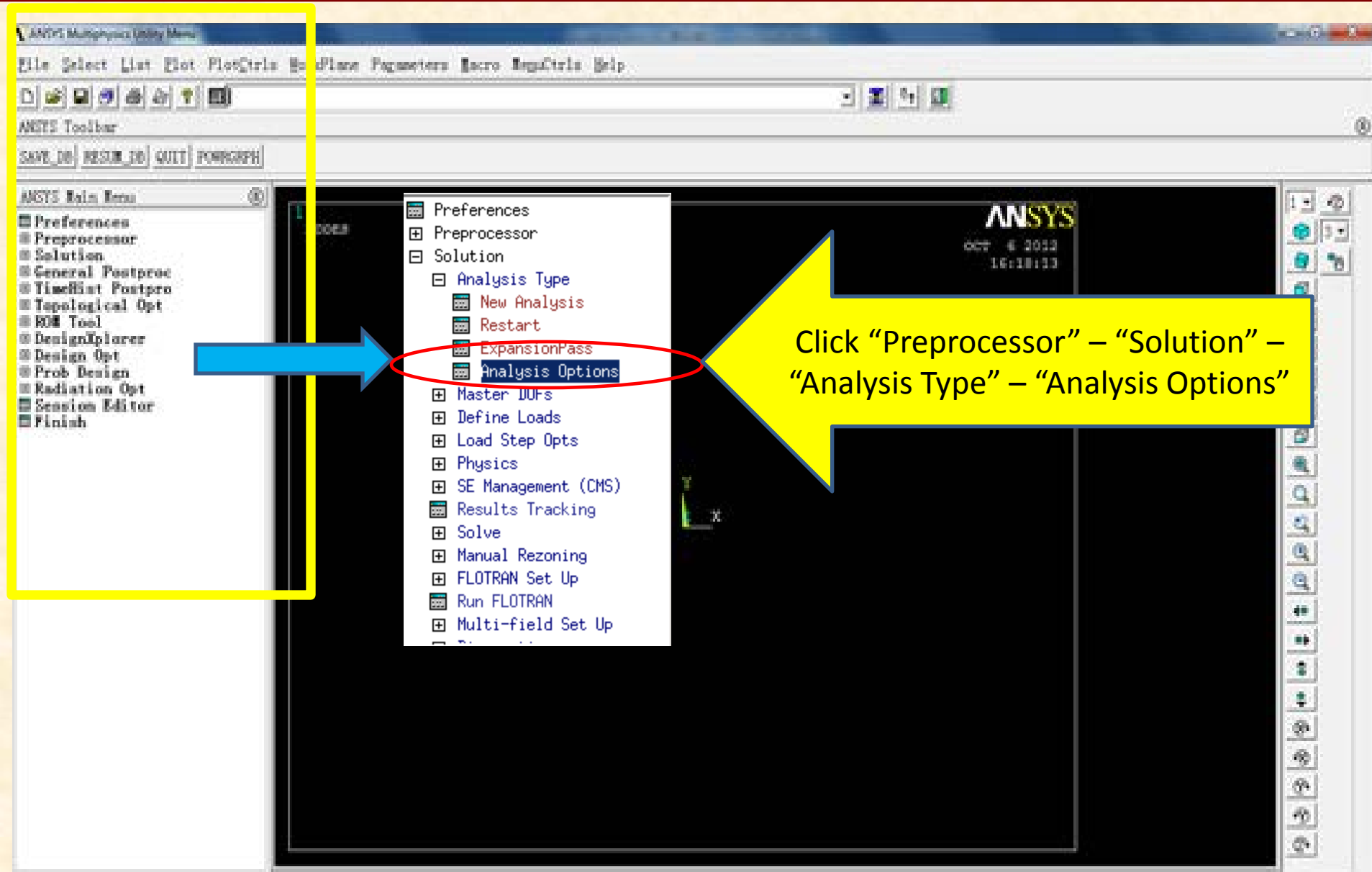
Main interface



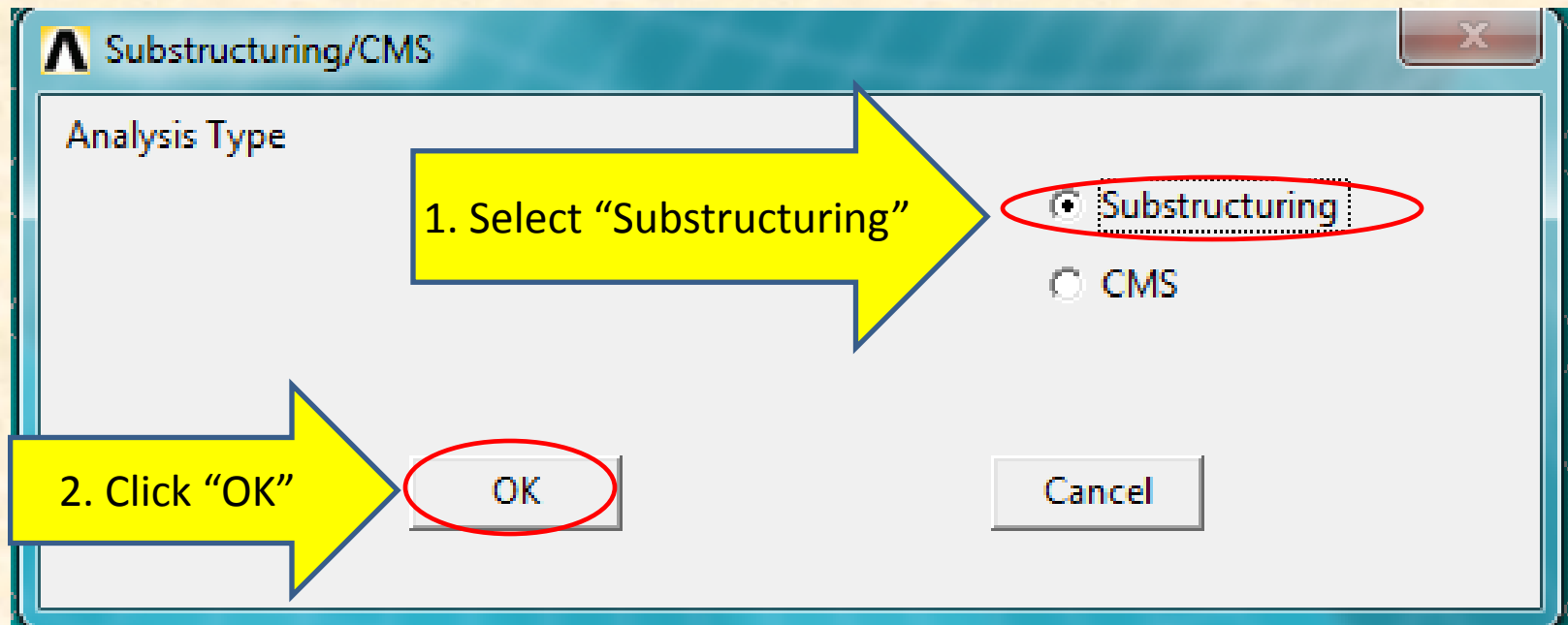
New Analysis



Main interface



Substructuring/CMS



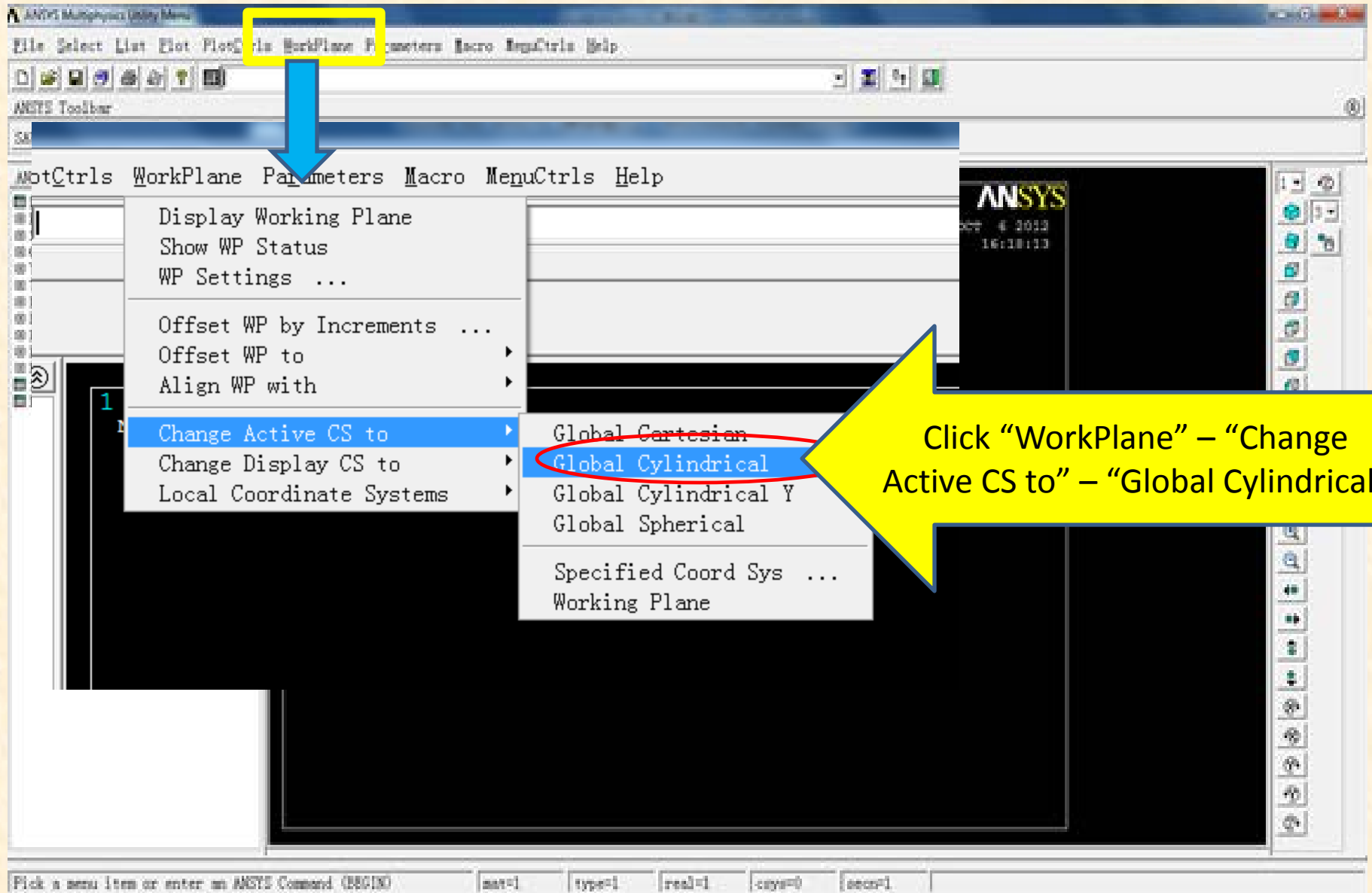
Substructuring Analysis

The screenshot shows a dialog box titled "Substructuring Analysis" with the following options:

- [SEOPT] Substructuring Analysis Options
 - Sename Name to be assigned to - Pad-1995
 - superelement matrix file
 - SEMATR Matrices to be generated - Stiffness
 - SEPR Items to be printed - No printout
 - SESST Save space for [S]? No
([S] is the stress stiffness matrix)
 - EXPMTH Expansion pass method - Backsubstitution
- [LUMPM] Use lumped mass approx? No
(if mass matrix is to be generated)
- [SEGEN] Automatic superelement generation
 - Generate during SOLVE?
 - No
 - Yes
 - Preview

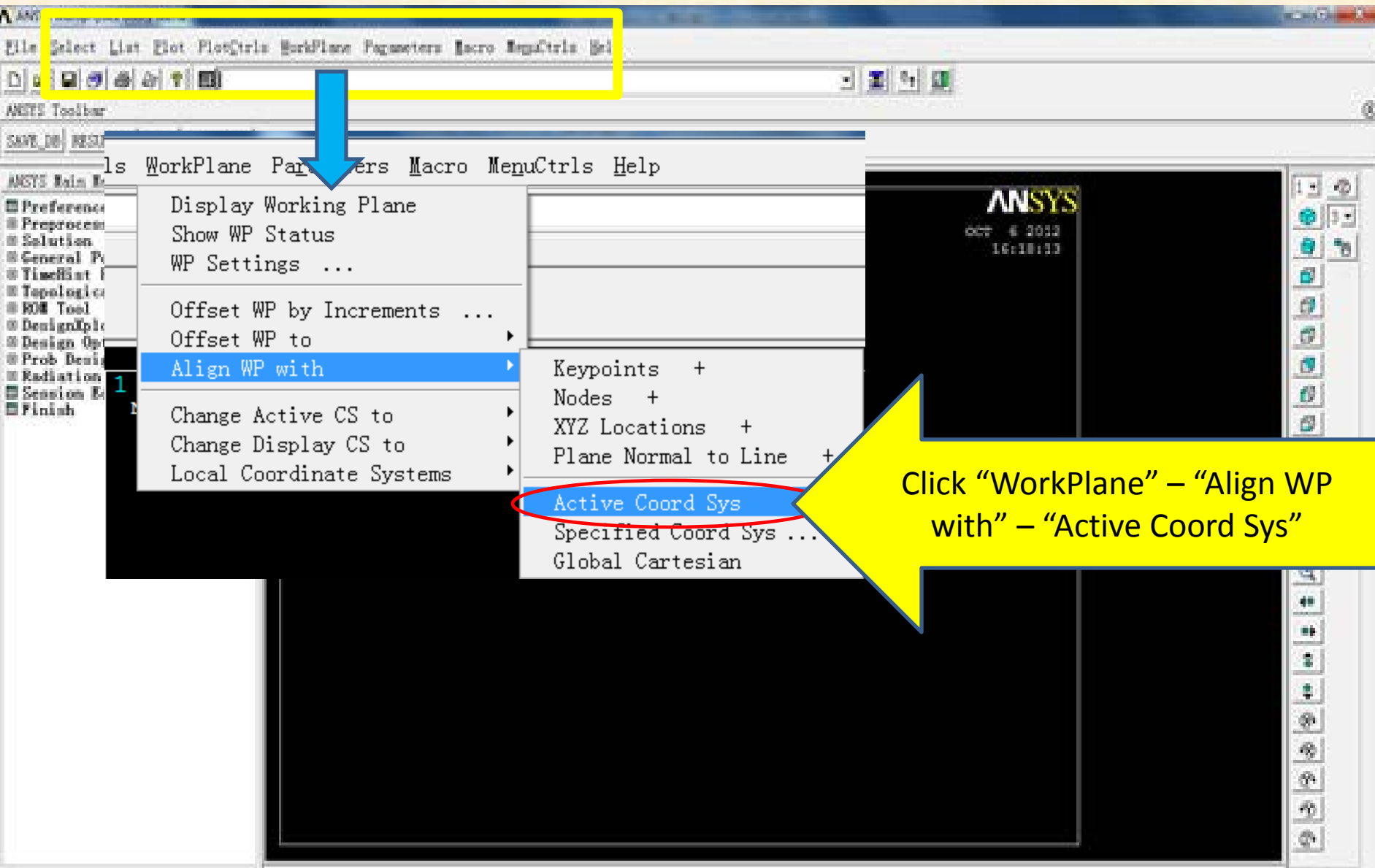
Annotations include a yellow arrow pointing to the "Yes" radio button with the text "1. Select 'Yes'", and another yellow arrow pointing to the "OK" button with the text "2. Click 'OK'". Both the "Yes" radio button and the "OK" button are circled in red.

Main interface



Click "WorkPlane" – "Change Active CS to" – "Global Cylindrical"

Main interface



Click "WorkPlane" – "Align WP with" – "Active Coord Sys"

Main interface

The screenshot displays the ANSYS Multiphysics main interface. On the left, the 'ANSYS Main Menu' is visible, with a yellow box highlighting the 'Preprocessor' section. A blue arrow points from this section to the 'Preprocessor' menu in the main application window. The 'Preprocessor' menu is expanded, showing the following path: 'Modeling' -> 'Create' -> 'Nodes' -> 'Rotate Node CS' -> 'To Active CS'. The 'To Active CS' option is highlighted with a red oval. A large yellow arrow points from this menu path towards the main workspace area on the right, which shows the ANSYS logo and a date/time stamp.

Click "Preprocessor" – "Modeling" –
 "Create" – "Nodes" – "Rotate Node
 CS" – "To Active CS"

Rotate Nodes into CS

Rotate Nodes into CS

Pick Unpick

Single Box

Polygon Circle

Loop

Count = 0

Maximum = 945

Minimum = 1

Node No. =

List of Items

Min, Max, Inc

OK Apply

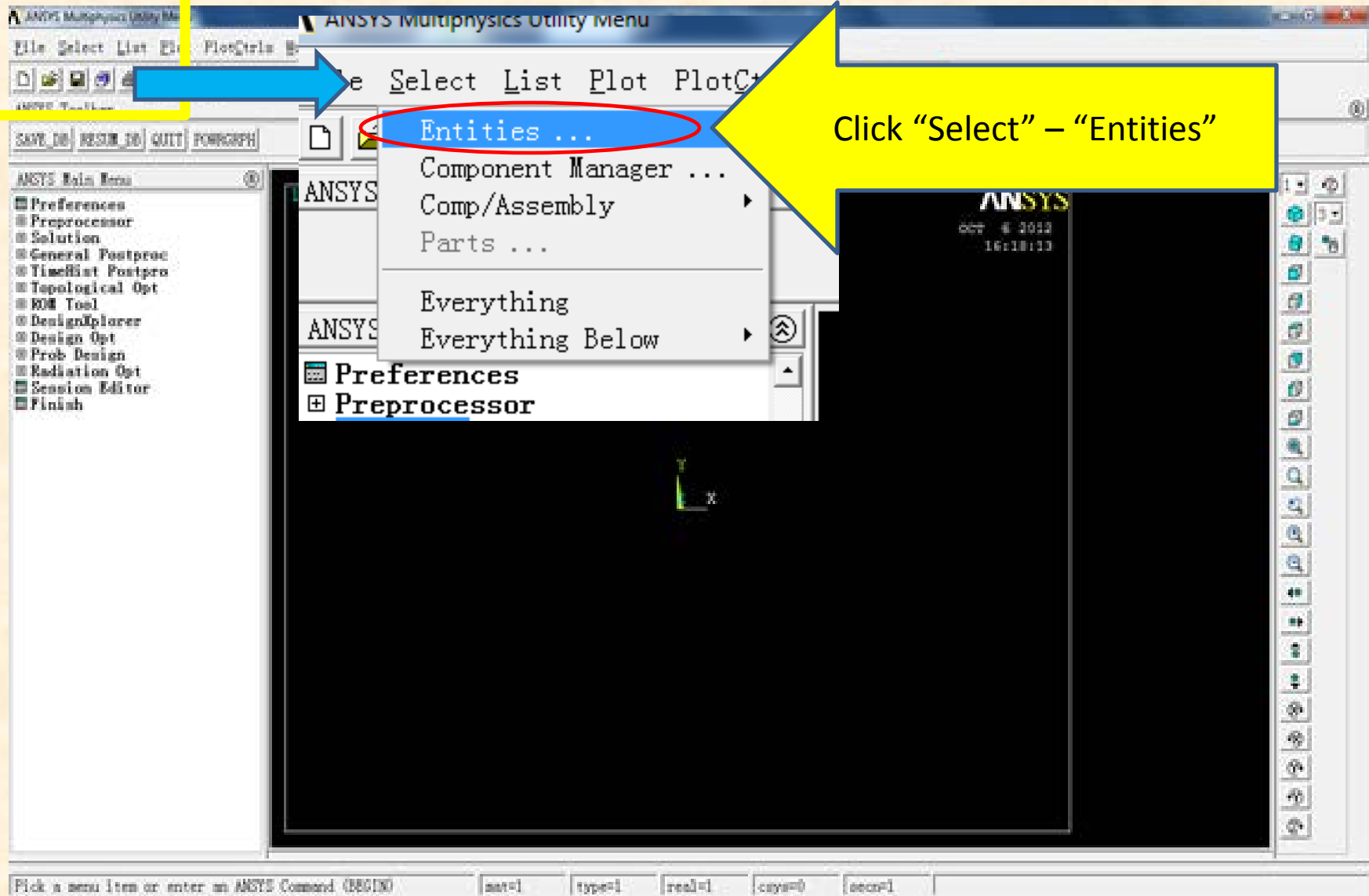
Reset Cancel

Pick All Pick None

2. Click "OK"

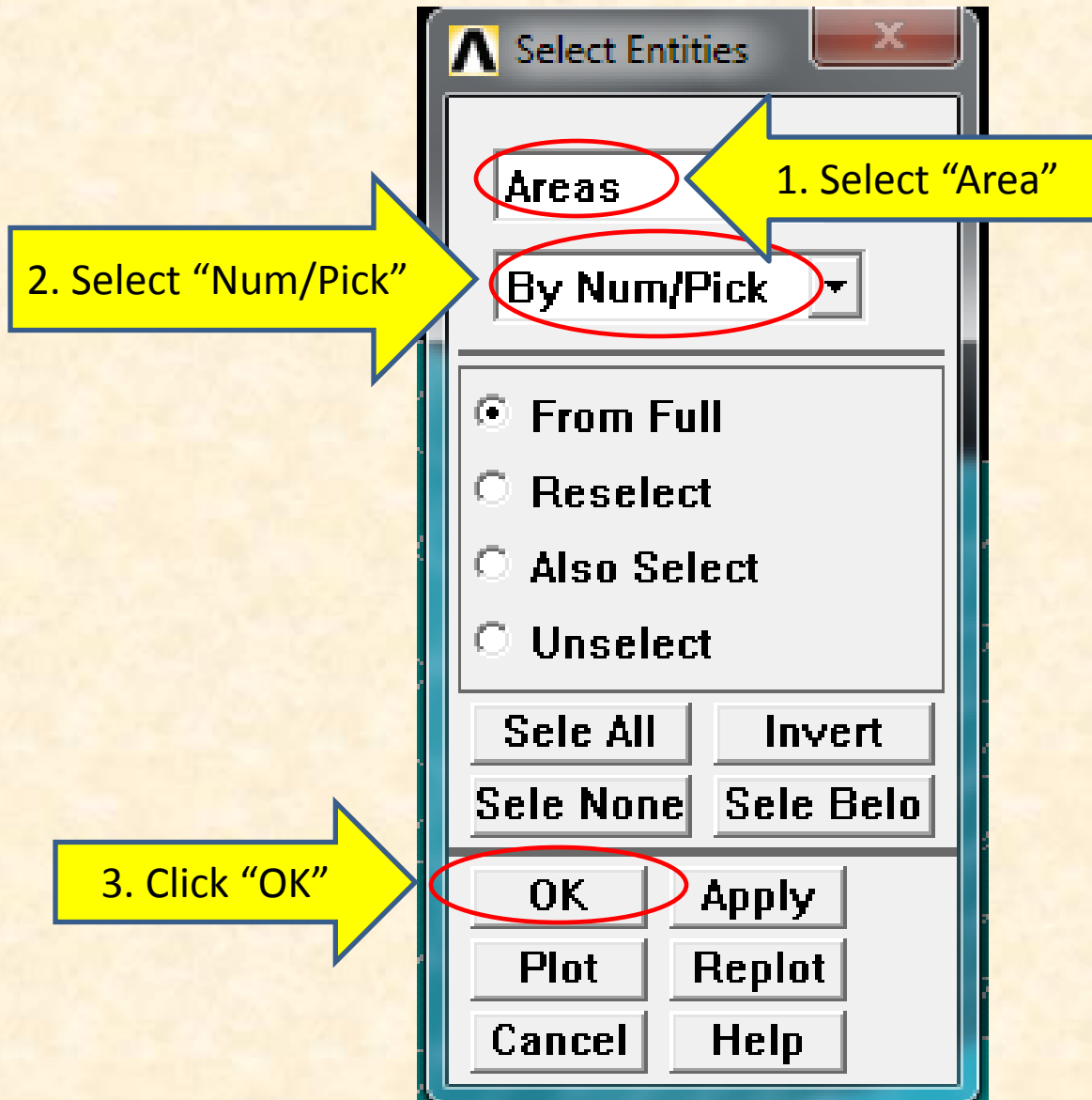
1. Click "Pick All"

Main interface



Click "Select" - "Entities"

Select Entities



Select areas

1

ANSYS

OCT 6 2012
23:48:04

Select areas

Pick Unpick

Single Box

Polygon Circle

Loop

Count = 1
Maximum = 6
Minimum = 1
Area No. = 4

List of Items
 Min, Max, Inc

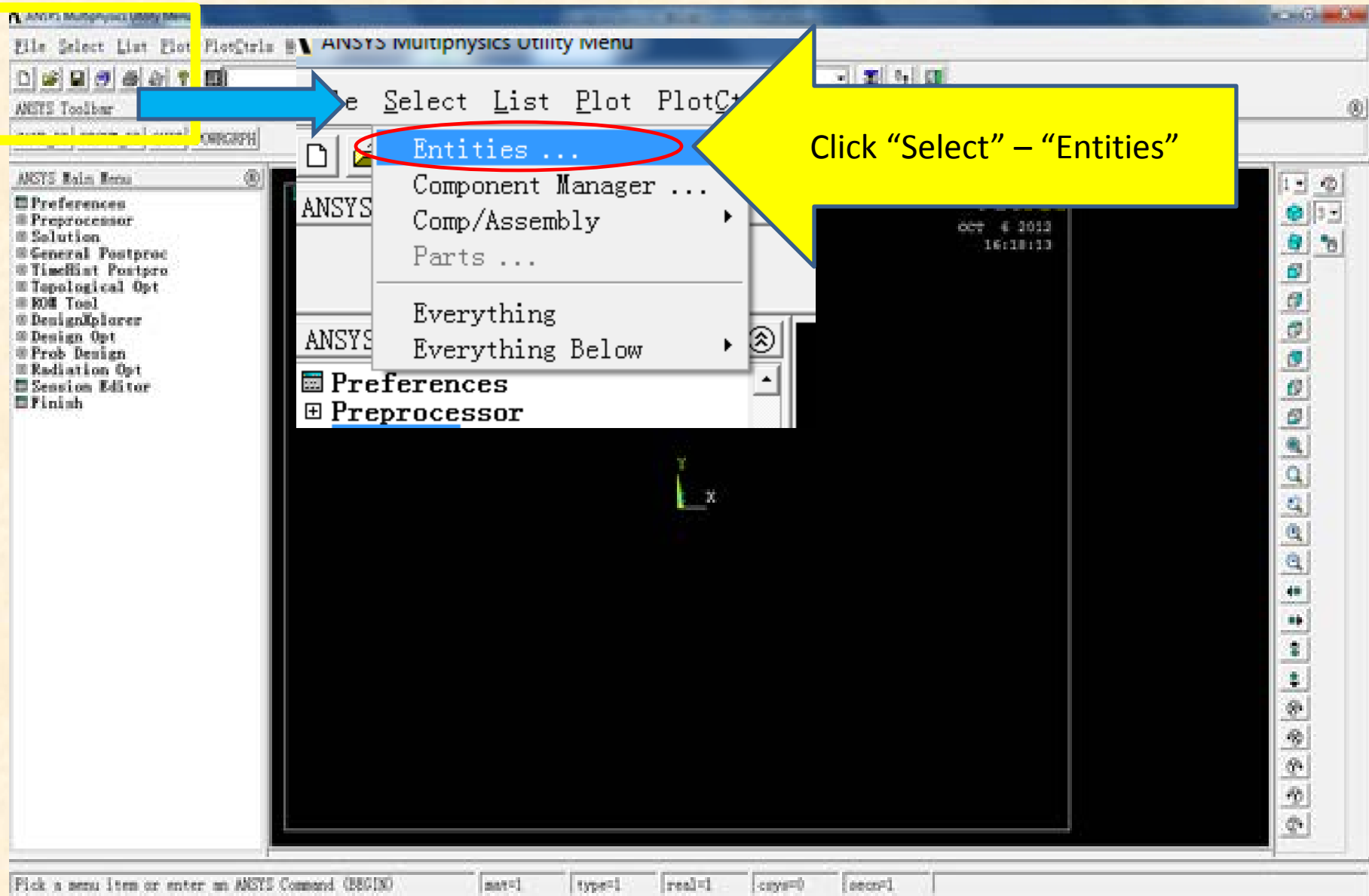
OK
Reset
Pick All

1. Select the upper surface

2. Click "OK"

stiffness_matrix_surface_1

Main interface



Select Entities

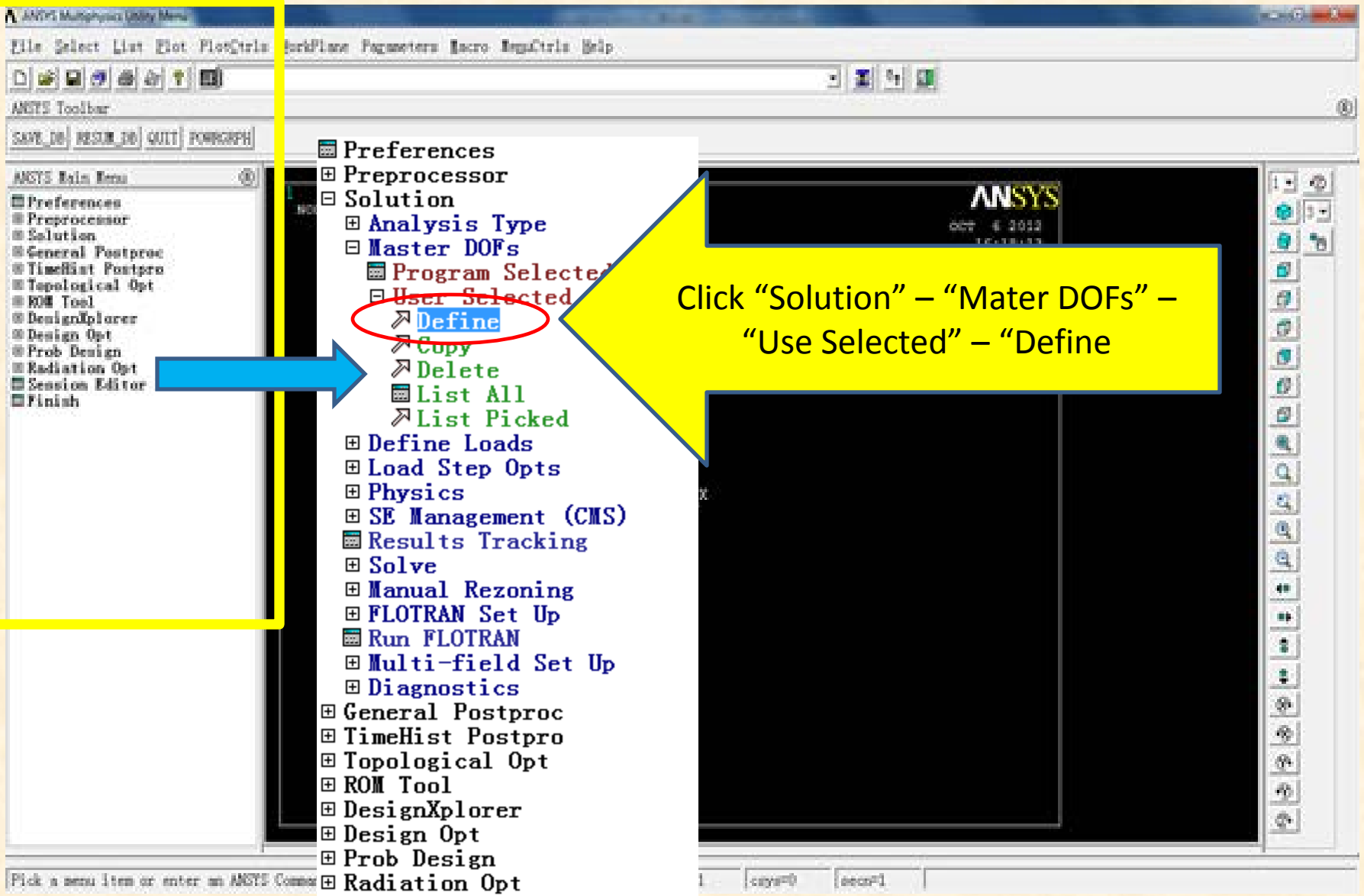
The image shows a 'Select Entities' dialog box with the following elements and callouts:

- 1. Select "Nodes"**: A yellow arrow points to the 'Nodes' text, which is circled in red.
- 2. Select "Attached to"**: A yellow arrow points to the 'Attached to' dropdown menu, which is circled in red.
- 3. Select "Sele All"**: A yellow arrow points to the 'Sele All' button, which is circled in red.
- 4. Select "OK"**: A yellow arrow points to the 'OK' button, which is circled in red.

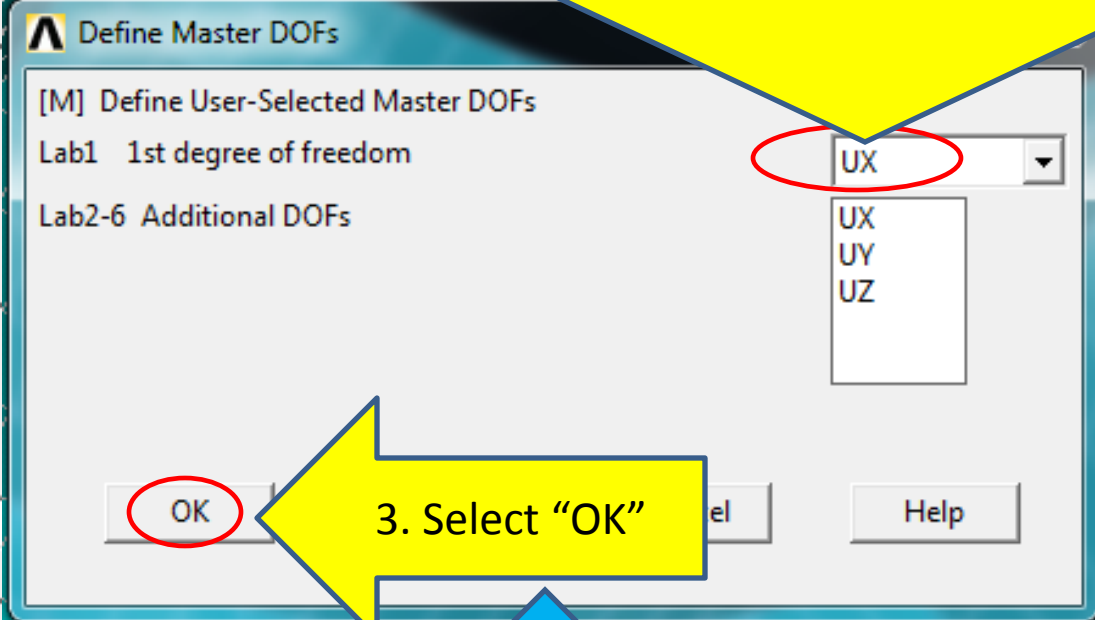
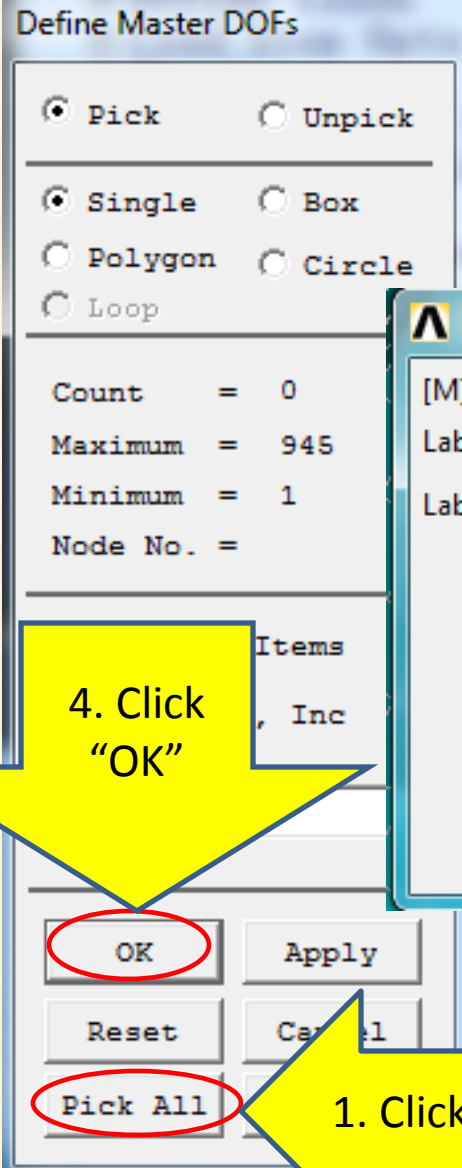
The dialog box contains the following options and buttons:

- Nodes
- Attached to
- Elements
- Keypoints
- Lines, all
- Lines, interior
- Areas, all
- Areas, interior
- Volumes, all
- Volumes, interior
- From Full
- Reselect
- Also Select
- Unselect
- Sele All
- Sele None
- OK
- Apply
- Plot
- Replot
- Cancel
- Help

Main interface



Define Master DOFs



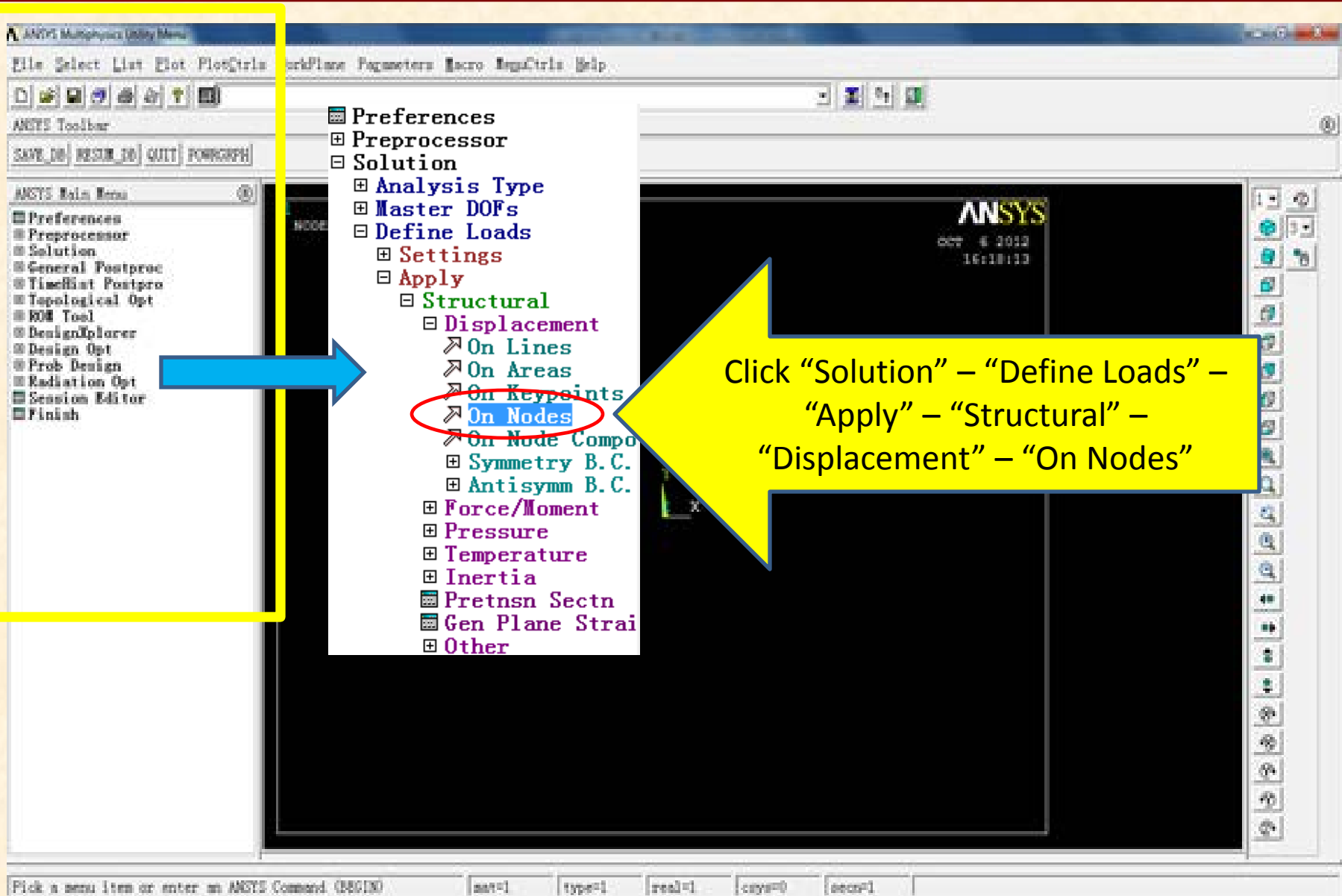
2. Select "UX"
NOTE: In ANSYS®, UX, UY and UZ stand for R (radial), ϑ (angular) and Z (axial) directions.

4. Click "OK"

3. Select "OK"

1. Click "Pick All"

Main interface



Apply U, ROT on Nodes

1
ELEMENTS
U
MAST

Apply U,ROT on Nodes

Pick Unpick

Single Box

Polygon Circle

Loop

Count = 0
Maximum = 945
Minimum = 1

ems

OK Apply

Reset Cancel

Pick All Help

1. Choose the nodes on the middle line of the lower surface of the pad model and the two lines close to the middle line

2. Click "OK"

stiffness_matrix_surface_1

ANSYS
OCT 7 2012
22:29:10

Apply U,ROT on Nodes

Apply U,ROT on Nodes

[D] Apply Displacements (U,ROT) on Nodes
Lab2 DOFs to be constrained

Apply as
If Constant value then:
VALUE Displacement value

AILDOF
UX
UY
UZ
VELX

Constant value

0

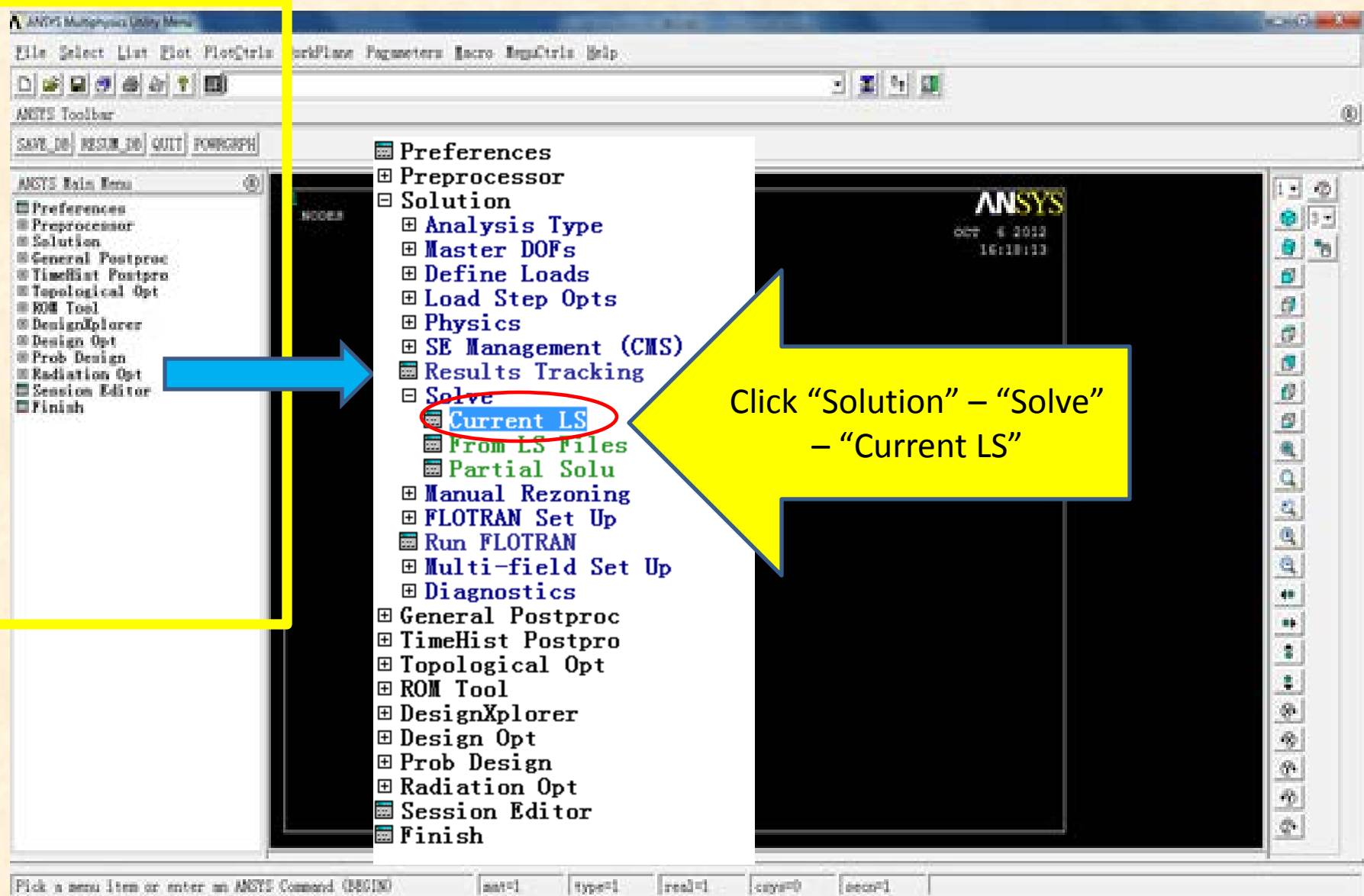
OK Cancel Help

1. Select "UX"

2. Type in "0"

3. Select "OK"

Main interface



Solve Current Load Step

Solve Current Load Step

[SOLVE] Begin Solution of Current Load Step

Review the summary information in the lister window (entitled "/STATUS Command"), then press OK to start the solution.

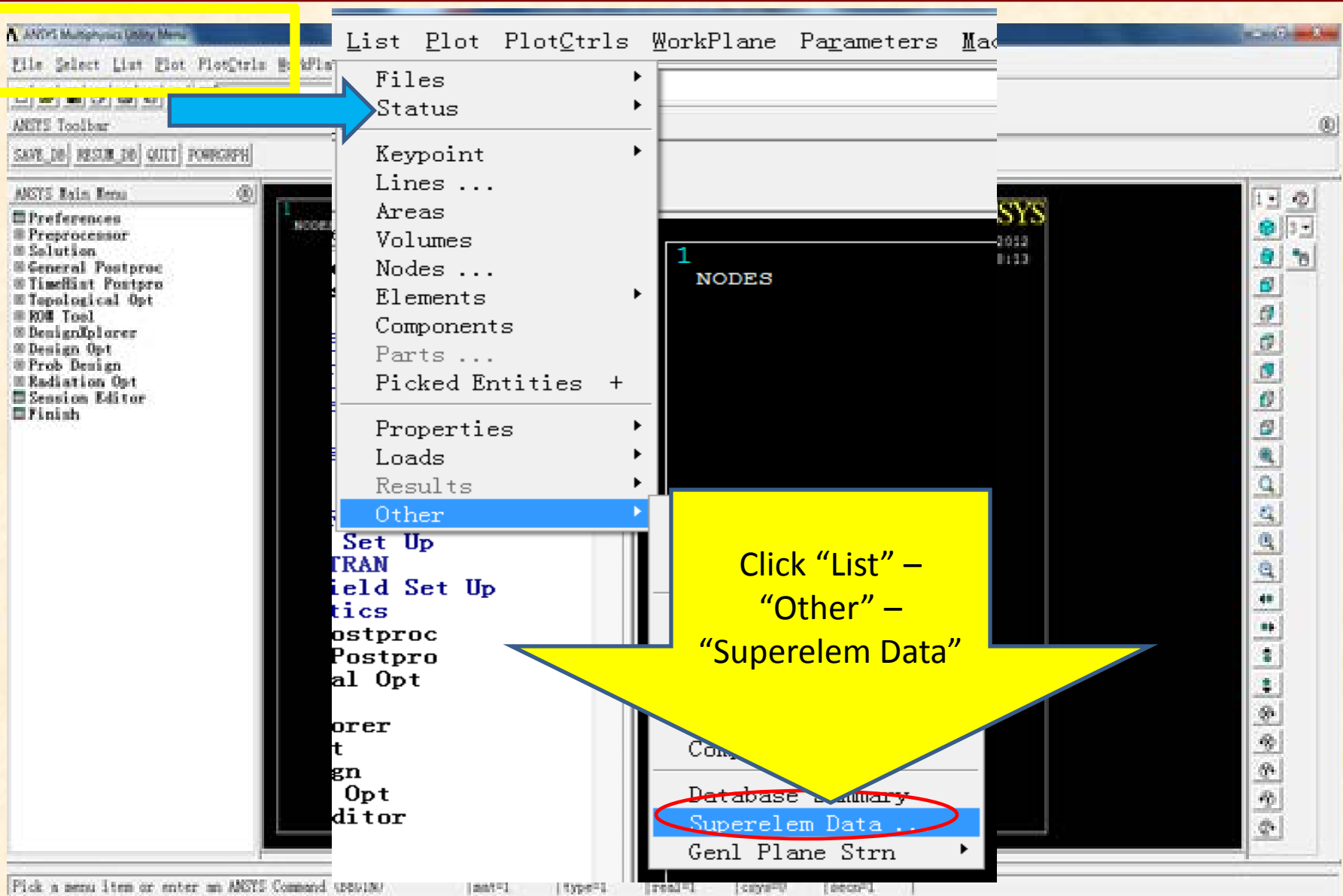
Select "OK"

OK Cancel Help

Rotate Nodes into CS

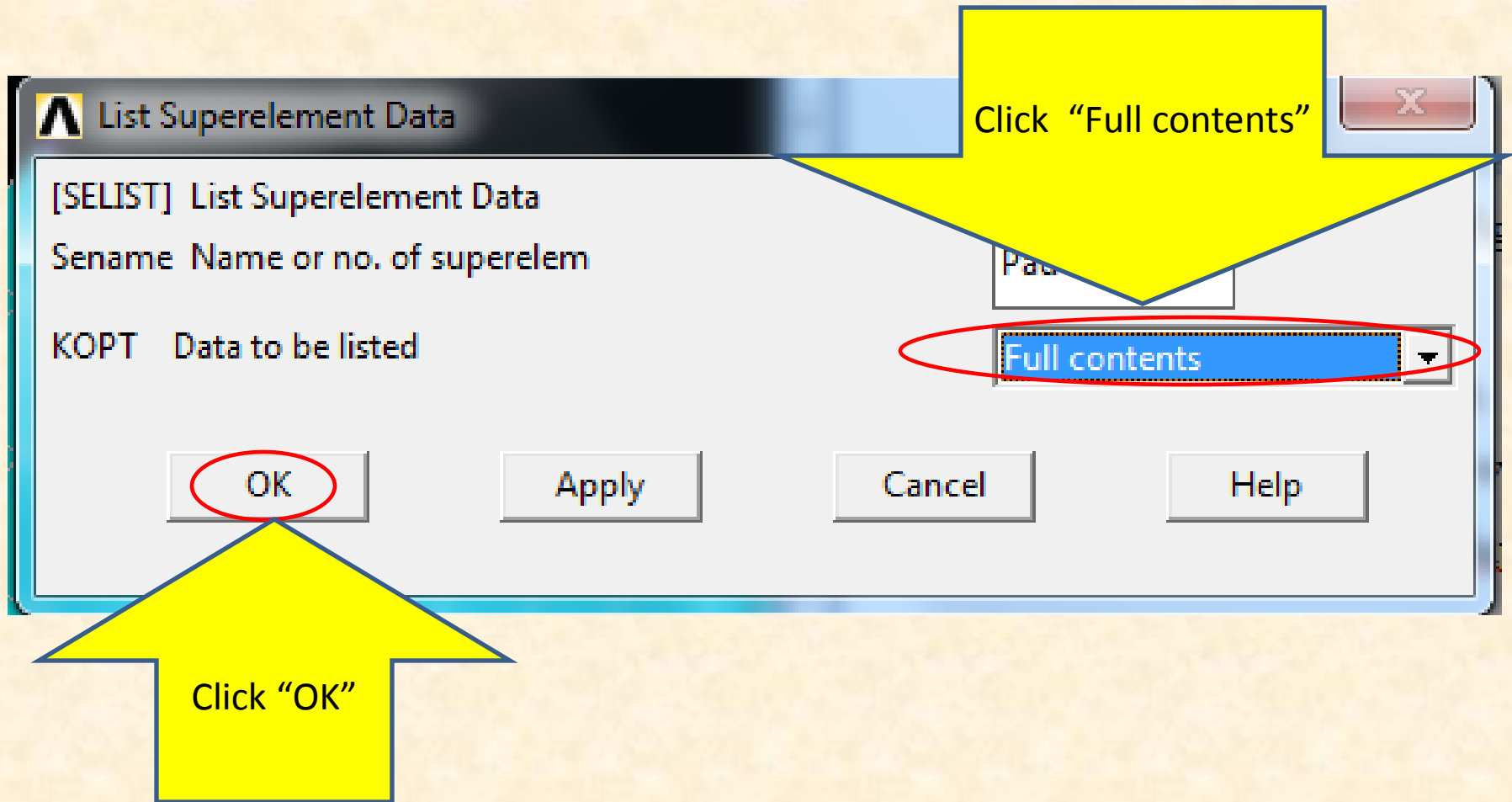
After solving the problem, the dialog box showing “Solution is done” will appear. Click “Close” to close the window.

Main interface



Click "List" –
 "Other" –
 "Superelem Data"

List Superelement Data



List of stiffness matrix

```

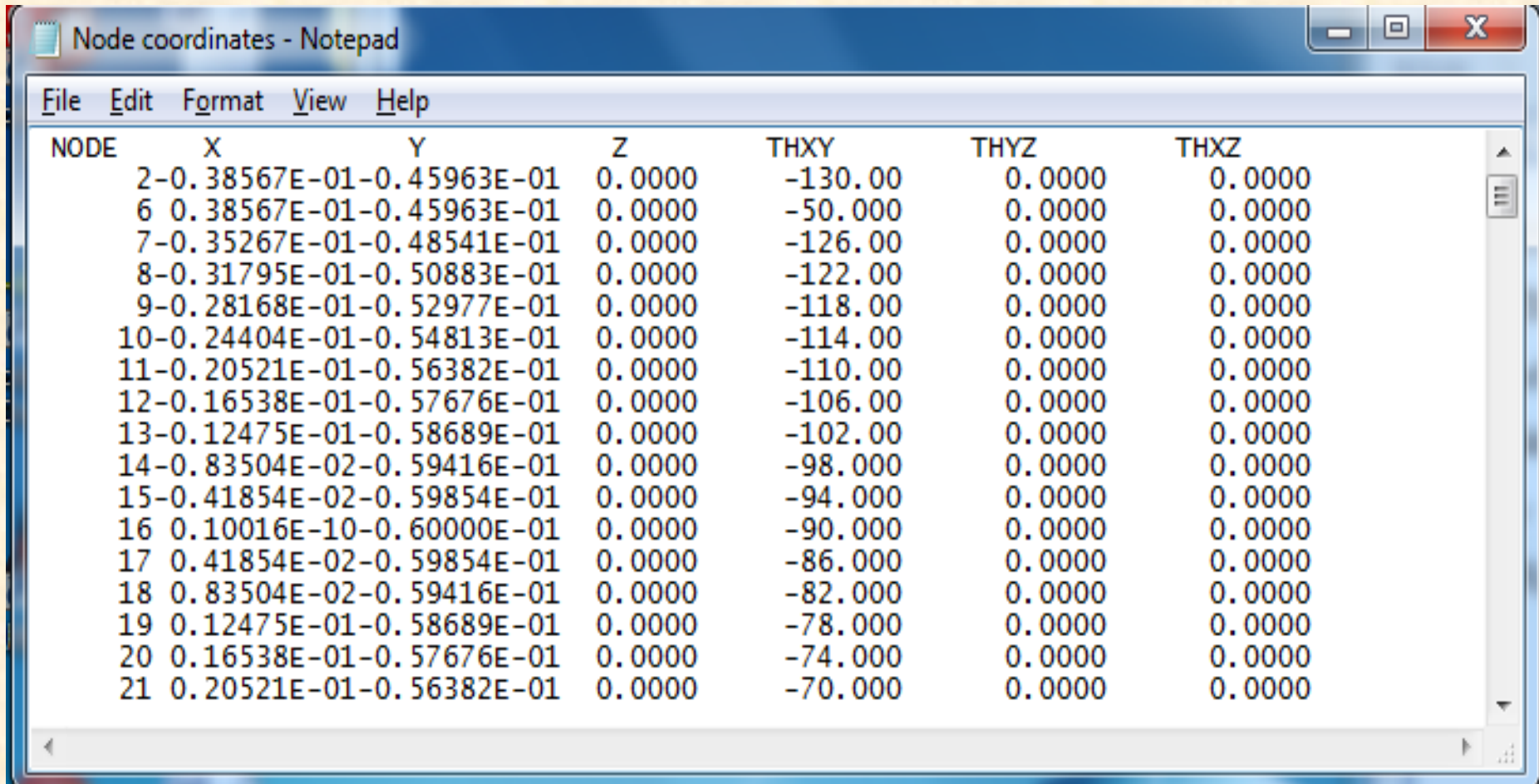
SELIST Command
File
PRINT CONTENTS OF SUPERELEMENT Pad-1995
PRINT OPTION =      3

HEADER =

      8          189          1          256          3
    945           0        945          189          0
      1           0           0           0           2
  2000000        2617          1        2835        81566
      0          9557        166          103        Pad-
  1995          9534          0
      922          928        1504        2452        2475
  2667          5502        9342           0        1120
  1312  109009479          0           0           3
      0           0           0           0           0
      419           0           0           0           0
      0           0           0           0           0

HEADER SUMMARY:
NUMBER OF ROWS =   189
NUMBER OF MATRICES =    1
NUMBER OF EDGE PLOT LINES =   256
NUMBER OF DEGREES OF FREEDOM PER NODE =    3
NUMBER OF DEGREE OF FREEDOM INDICES =   945
NUMBER OF NODES =   189
  
```

TXT file for nodal coordinates



Node coordinates - Notepad

NODE	X	Y	Z	THXY	THYZ	THXZ
2	-0.38567E-01	-0.45963E-01	0.0000	-130.00	0.0000	0.0000
6	0.38567E-01	-0.45963E-01	0.0000	-50.000	0.0000	0.0000
7	-0.35267E-01	-0.48541E-01	0.0000	-126.00	0.0000	0.0000
8	-0.31795E-01	-0.50883E-01	0.0000	-122.00	0.0000	0.0000
9	-0.28168E-01	-0.52977E-01	0.0000	-118.00	0.0000	0.0000
10	-0.24404E-01	-0.54813E-01	0.0000	-114.00	0.0000	0.0000
11	-0.20521E-01	-0.56382E-01	0.0000	-110.00	0.0000	0.0000
12	-0.16538E-01	-0.57676E-01	0.0000	-106.00	0.0000	0.0000
13	-0.12475E-01	-0.58689E-01	0.0000	-102.00	0.0000	0.0000
14	-0.83504E-02	-0.59416E-01	0.0000	-98.000	0.0000	0.0000
15	-0.41854E-02	-0.59854E-01	0.0000	-94.000	0.0000	0.0000
16	0.10016E-10	-0.60000E-01	0.0000	-90.000	0.0000	0.0000
17	0.41854E-02	-0.59854E-01	0.0000	-86.000	0.0000	0.0000
18	0.83504E-02	-0.59416E-01	0.0000	-82.000	0.0000	0.0000
19	0.12475E-01	-0.58689E-01	0.0000	-78.000	0.0000	0.0000
20	0.16538E-01	-0.57676E-01	0.0000	-74.000	0.0000	0.0000
21	0.20521E-01	-0.56382E-01	0.0000	-70.000	0.0000	0.0000

TXT file for stiffness matrix

```
SELIST - Notepad
File Edit Format View Help
ROW 1 MATRIX 1
0.18249045E+09 620692.45 -83860811. -20684658. -3303790.0
-2274234.1 -992794.60 -340246.36 45859.415 224886.43
276729.76 260701.46 228434.23 205086.55 202072.35
222247.26 265811.79 340076.73 454052.62 625187.54
894890.83 -614135.57 764060.93 1212209.7 947803.74
740108.15 586014.25 450996.83 340886.54 247067.57
169133.76 105863.93 53427.623 9029.1427 -31358.573
-74649.360 -126249.96 -193838.02 -287077.05 -417120.42
-600739.76 -878618.39 776251.40 372302.37 161815.86
10570.850 -140291.41 -350608.33 -757746.80 32802109.
-14786992. -920554.34 -1417117.3 241668.09 400431.61
1034970.0 -72318507. -15221866. -2471425.6 -1656048.9
324266.81 835013.90 1673999.9 -21485641. -5671629.7
-2729830.5 -403333.45 266422.23 860775.63 1321137.2
-3134638.7 -3348016.9 -1125850.8 -130385.39 525729.47
752946.55 1099094.4 -1915411.3 -1114810.7 -387548.66
321397.10 573330.69 703745.40 882853.01 -656490.38
```

CONCLUSIONS

Accurate characterization of mechanical components presently requires of physical models of ever increasing complexity that include all relevant geometrical aspects, material properties and fluid flow conditions, as per installation and operation, even envisioning operation well beyond their intended original design. This paper extends a computational thin film fluid flow model for tilting pad journal bearings to include both pivot flexibility and pad flexibility on the prediction of the static and dynamic forced performance of typical bearings. Presently, a FE structural commercial model builds the stiffness matrix for pad displacements, further reduced to show only the pad surface deformations due to an applied pressure field. The deformation field is integrated into the evaluation of film thickness for solution of the Reynolds equation delivering the hydrodynamic pressure field. A small amplitude perturbation analysis produces equations for the zeroth and first-order pressure fields from which the load capacity and a multitude of complex stiffnesses coefficients for each pad are determined. A pad assembly with frequency reduction method delivers the 4x4 complex dynamic stiffnesses for lateral or radial displacements of the shaft center. Curve fits of the force/displacement frequency functions delivers the bearing stiffness, damping and force coefficients.

In a parallel program, Gaines [1] performed measurements for three-pad LBP TPJBs. Pad thickness varied to change the pad flexibility. Gaines' test data show that pad flexibility affects little the journal eccentricity and maximum pad temperature, but has an impact on the bearing dynamic force coefficients. Current model predictions for the TPJBs in Ref. [1] correlate favorably with test data, except that the maximum temperature rise is underestimated by up to 40%. The reason may be due to that the oil pushed through the supply holes is not routed efficiently into the bearing pads. Pad flexibility reduces the predicted journal eccentricity and maximum temperature rise.

The real part of the TPJB dynamic stiffnesses (Z) at super-synchronous frequency ($\omega > \Omega$) shows a hardening effect, yet it is nearly constant for frequencies lesser than the shaft angular speed, $\omega < \Omega$. Hence, the virtual mass coefficients identified with a [K-C-

M] model are negative. However, the predicted virtual mass coefficients only affect the dynamic stiffnesses at a high frequency ($\omega > \Omega$). The imaginary part of the complex stiffnesses (Z) increases linearly with frequency. However, at 6 krpm (100 Hz), $\text{Im}(Z)$ show a decrease at a high frequency ($\omega > 180$ Hz).

In general, both the experimental and predicted dynamic force coefficients reduce due to pad flexibility. Pad flexibility causes up to a 20% reduction in predicted bearing stiffness. However, Gaines [1] reports an increase in bearing direct stiffnesses for the TPJB with the more flexible pad. Predicted bearing stiffnesses also show a slightly increase at a low load ($W/(LD) < 689$ kPa) due to pad flexibility but decrease at a high load (> 689 kPa).

Pad flexibility shows a more significant effect on the predicted dynamic force coefficients than the experimental results evidence, in particular at a high rotor speed ($\Omega = 12$ krpm), thus indicating the pad structural FE model delivers a smaller stiffness than that of the actual test pad. Measurements for the stiffness of the pads are needed to further improve the FE model. Predictions including pad flexibility deliver damping coefficients up to 46% lower than those obtained with a rigid pads model. Reducing the pad thickness from 11.5 mm to 8.5 mm causes also a reduction of up to 50% in predicted direct damping.

The paper also includes a parametric study to quantify the influence of pad thickness on the rotordynamic force coefficients of a sample TPJB with three pads of increasing preload=0, 0.25 (baseline) and 0.5. The bearing pads varies from rigid to flexible ($k_{pad} = \infty, 3.15$ and 7.33). Graphs for journal eccentricity, drag friction coefficient, and dimensionless force coefficients are presented versus Sommerfeld number ($0.2 < S < 2.2$). The operating journal eccentricity and the dynamic force coefficients of a TPJB reduce due to pad flexibility, in particular for operation at a large Sommerfeld number ($S > 0.8$). Pad flexibility shows a more pronounced effect for the TPJB with null pad preload.

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APPENDIX A. HIRS TURBULENT FLOW MODEL

XLTPJB© uses a laminar flow model to predict the static and dynamic performance of TPJBs. However, turbulent flow may occur in bearings operating at high rotational speeds, especially for bearings of large size, or with a lubricant of low viscosity [A.1]. The turbulent flow model based on Hirs [A.2] is introduced into the predictive model.

According to Ref. [A.3], the fully developed turbulent bulk-flow equations in a thin film are the

Continuity equation:

$$\frac{\partial}{\partial x}(hV_x) + \frac{\partial}{\partial z}(hV_z) + \frac{\partial h}{\partial t} = 0 \quad (\text{A.1})$$

Circumferential momentum equation:

$$-h \frac{\partial P}{\partial x} = \frac{\mu}{h} \left(\kappa_x V_x - \kappa_J \frac{U}{2} \right) + \rho \frac{\partial (hV_x)}{\partial t} \quad (\text{A.2})$$

Axial momentum equation:

$$-h \frac{\partial P}{\partial z} = \frac{\mu}{h} \kappa_z V_z + \rho \frac{\partial (hV_z)}{\partial t} \quad (\text{A.3})$$

where $V_x = \frac{1}{h} \int_0^h \bar{u} dy$ and $V_z = \frac{1}{h} \int_0^h \bar{w} dy$ are bulk flow velocities.

$\kappa_J = f_J R_J$, $\kappa_B = f_B R_B$ and $\kappa_x = \kappa_z = \frac{1}{2}(\kappa_J + \kappa_B)$ are bulk flow turbulence shear parameters. For laminar flows, $\kappa_x = \kappa_z = \kappa_J = \kappa_B = 12$. Above, h and P denote the film thickness and pressure, respectively, and $U = \Omega R$ is the journal surface velocity. Note equations (A.2-3) include temporal fluid inertia terms only, with the advection terms omitted.

Substituting the momentum Eqs.(A.2, A.3) into the continuity Eq.(A.1) gives

$$\begin{aligned} \frac{\partial}{\partial x} \left\{ \frac{h^3}{\kappa_x \mu_{(T)}} \frac{\partial P}{\partial x} \right\} + \frac{\partial}{\partial z} \left\{ \frac{h^3}{\kappa_z \mu_{(T)}} \frac{\partial P}{\partial z} \right\} &= \frac{\partial h}{\partial t} + \frac{\Omega}{2} \frac{\partial}{\partial \theta} \left(\frac{\kappa_J}{\kappa_x} h \right) \\ - \left\{ \frac{\rho h^2}{\kappa_x \mu_{(T)}} \frac{\partial^2 h}{\partial x \partial t} + \frac{\rho h^2}{\kappa_z \mu_{(T)}} \frac{\partial^2 h}{\partial z \partial t} + \frac{\partial (hV_x)}{\partial t} \frac{\partial}{\partial x} \left(\frac{\rho h^2}{\kappa_x \mu_{(T)}} \right) + \frac{\partial (hV_z)}{\partial t} \frac{\partial}{\partial z} \left(\frac{\rho h^2}{\kappa_z \mu_{(T)}} \right) \right\} & \end{aligned} \quad (\text{A.4})$$

Taking the time derivative of the continuity Eq.(A.1) gives

$$\frac{\partial^2}{\partial t \partial x}(hV_x) + \frac{\partial^2}{\partial t \partial z}(hV_z) + \frac{\partial^2 h}{\partial t^2} = 0 \quad (\text{A.5})$$

Assume $\kappa_x = \kappa_z = \kappa$. Substituting Eq. (A.5) into Eq.(A.4) yields a Reynolds equation for turbulent flow with temporal fluid inertia effects.

$$\frac{1}{R_j^2} \frac{\partial}{\partial \theta} \left\{ \frac{h^3}{\kappa \mu_{(T)}} \frac{\partial P}{\partial \theta} \right\} + \frac{\partial}{\partial z} \left\{ \frac{h^3}{\kappa \mu_{(T)}} \frac{\partial P}{\partial z} \right\} = \frac{\partial h}{\partial t} + \frac{\Omega}{2} \frac{\partial h}{\partial \theta} + \frac{\rho h^2}{\kappa \mu_{(T)}} \frac{\partial^2 h}{\partial t^2} \quad (\text{A.6})$$

The turbulence shear parameters are

$$\kappa_x = \kappa_z = \kappa = (\kappa_J + \kappa_B)/2 \quad (\text{A.7})$$

According to Refs. [A.3, A.4], κ_J and κ_B are determined from

$$\kappa_i = \begin{cases} 12 & \text{Re}_i < 1,000 \\ \frac{12}{\text{Re}_i} \left[1 - \left(\frac{\text{Re}_i - 1000}{20000} \right)^2 (3 - 2 \cdot \text{Re}_i) \right] + f_i \cdot \left(\frac{\text{Re}_i - 1000}{20000} \right)^2 (3 - 2 \cdot \text{Re}_i) & 1,000 \leq \text{Re}_i < 3,000 \\ f_i \text{Re}_i & \text{Re}_i > 3,000 \end{cases} \quad (\text{A.8})$$

over the flow regimes from laminar, to transition, to fully developed turbulence. Above,

$$f_i = 0.001375 \left[1 + \left(\frac{5 \cdot 10^5}{\text{Re}_i} \right)^{\frac{1}{3}} \right]$$

is Moody's friction factor, and $i=J$ (journal) or B (bearing)

surfaces. The bulk-flow Reynolds numbers (Re_i) relative to the bearing (B) and journal (J) surfaces are:

$$\text{Re}_J = \frac{\rho h}{\mu} \left[(V_x - \Omega R)^2 + V_z^2 \right]^{\frac{1}{2}}, \quad \text{Re}_B = \frac{\rho h}{\mu} \left[V_x^2 + V_z^2 \right]^{\frac{1}{2}} \quad (\text{A.9})$$

The bulk flow thermal energy transport equation with turbulent flow is [A.3]

$$C_p \left[\frac{\partial(\rho HUT)}{\partial x} + \frac{\partial(\rho HWT)}{\partial z} \right] + Q_s = H \frac{\Omega R}{2} \frac{\partial P}{\partial x} + \frac{\mu}{H} \left[\kappa_x \left(U^2 + W^2 + \frac{\Omega R}{2} U \right) + \kappa_J \Omega R \left(\frac{\Omega R}{4} - U \right) \right] \quad (\text{A.10})$$

Note the drag torque is obtained from

$$T_{\text{Torque}} = \iint \left\{ \frac{h}{2} \frac{\partial P}{\partial x} + \frac{\kappa_J \mu}{4 h} \Omega R \right\} R dx dy \quad (\text{A.11})$$

Predictions from the implemented turbulent flow model are compared against the test data and predictions of Taniguchi *et al.* [A.5] collected for a large size four-pad TPJB used in a steam turbine. Ref. [A.5] adopts the turbulent flow model by Ng and Pan [A.6]. Table A.1 lists the geometry and lubricant properties of the test TPJB and Figure A.1 shows a schematic view of the four-pad bearing.

Table A.1 Geometry and lubricant properties of the test TPJB in Ref [A.5]

Rotor diameter, D	479 mm
Pad axial length, L	300 mm
Pad arc angle, θ_p	80°
Pivot offset	0.5
Pad clearance, C_B	612 μm
Pad preload	0
Pad mass, m_p	119 kg
Pad inertia, I_p	2.74 $\text{kg}\cdot\text{m}^2$
Lubricant type	ISO VG32
Supply lubricant temperature, T_{in}	~40°C
Supply lubricant pressure	1 bar
Lubricant density, ρ	856.2 kg/m^3
Viscosity at 40°C, μ_0	0.0274 Pa·s
Viscosity temperature coefficient, α	0.0296 1/°C
Specific heat capacity at 40 °C	1950.7 J/(kg·K)
Journal speed, Ω	3000 rpm
Surface speed, ΩR	75.24 m/s
Unit load, $W/(LD)$	1,252 kPa
Average shear Reynolds number ⁶ , Re	2,043

⁶ $\text{Re}=\rho\Omega RC_B/\mu_{average}$ with $\mu_{average}$ is the lubricant viscosity at the average temperature in the fluid film.

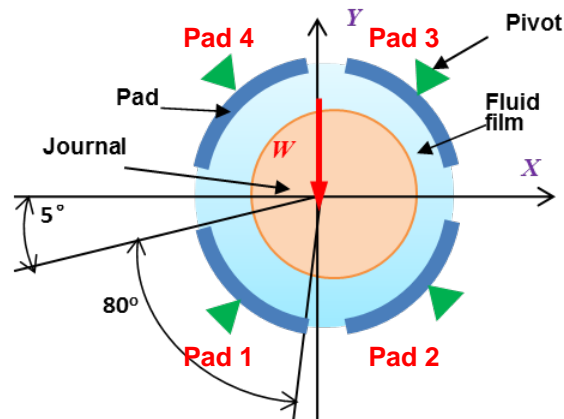


Figure A.1 Schematic view of large four-pad TPJB in Ref [A.5]

Note that the current model predictions do not account for either pivot or pad flexibility, as also in the model stated in Ref. [A.5]. The clearance is assumed constant, not affected by an increase in temperature, to keep consistency with Ref. [A.5]. The journal and bearing surfaces are taken as adiabatic. The inlet heat carry over coefficient (λ) is 0.8. The average shear Reynolds number ($Re=2,043$) is higher than 2,000, thus indicating a turbulent flow condition.

Figure A.2 shows the fluid film thickness at the mid-plane of the loaded two pads (#1 and #2). The current predictions and those from [A.5] agree with each other, albeit lower than the test results by 19% at the pad inlet, and correlating well for the magnitude of the minimum film thickness.

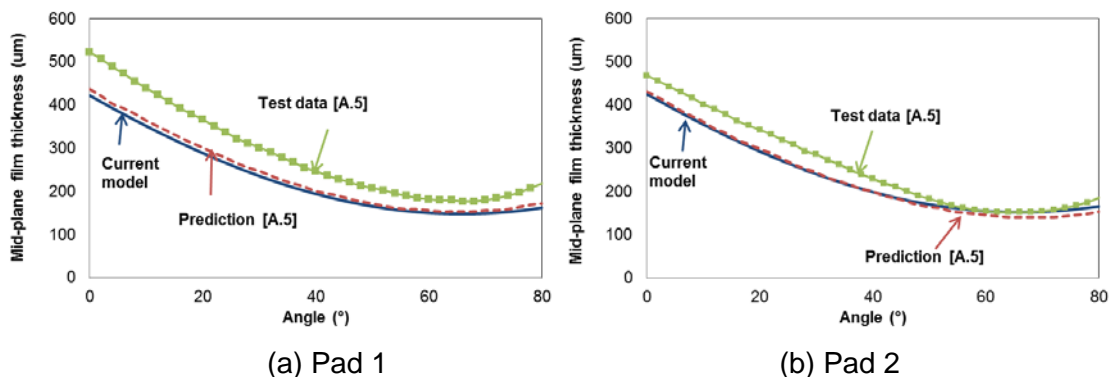


Figure A.2 Film thickness at mid-plane of pads #1 and #2. Current predictions and results in Ref. [A.5]. $W/(LD)=1,252$ kPa, $\Omega=3$ krpm, average $Re=2,043$.

Figure A.3 shows the pressure profile at the mid-plane of the loaded pads (#1 and #2). The current predictions give pressure profiles similar to those predicted in Ref. [A.5]. Note the models ignore the pressure rise at the pad inlet due to an entrance fluid inertia effect.

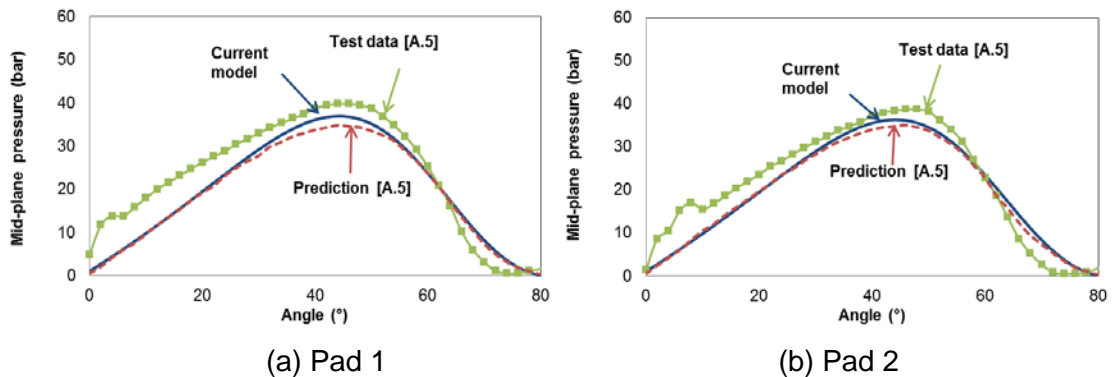


Figure A.3 Pressure profile at the mid-plane of pads #1 and #2. Current predictions and results in Ref. [A.5]. $W(LD)=1,252$ kPa, $\Omega=3$ krpm, average $Re=2,043$.

Figure A.4 depicts the temperature at the mid-plane of the loaded pads (#1 and #2). The current predictions and the data in Ref. [A.5] show good agreement at the pads' leading edges. However, at the pads' trailing edges, the current model delivers a temperature lower than the measured one and the predicted results in Ref. [A.5]. The current model takes a bulk temperature across the film whereas the measurement is obtained at 3 mm under a pad surface.

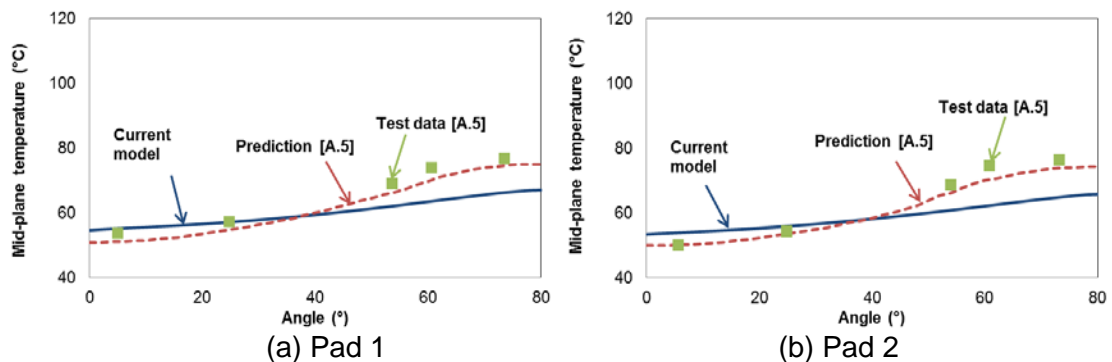


Figure A.4 Temperature at the mid-plane of pads #1 and #2. Current predictions and results in Ref. [A.5]. $W(LD)=1,252$ kPa, $\Omega=3$ krpm, average $Re=2,043$.

Figure A.5 depicts the drag power loss predicted by the current model and those in Ref. [A.5]. For rotor speeds from 1,200 rpm to 3,800 rpm, the predicted power loss agrees with the data in Ref. [A.5]. However, for shaft speeds above 3,800 rpm, the current model overestimates slightly the drag power loss. Note that during its operation, the thin film flow in the bearing is laminar ($Re < 1,000$) for rotor speeds below 2,900 rpm, in the transition flow regime for rotor speeds from 2,900 rpm to 3,700 rpm, and fully turbulent above 3,700 rpm.

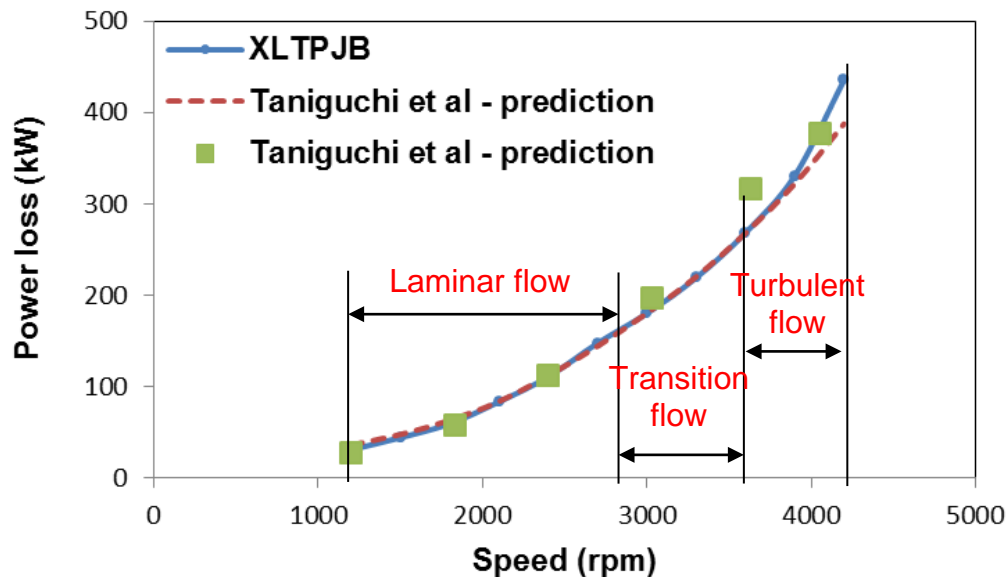


Figure A.5 Bearing friction loss at $W/(LD)=1,252$ kPa. Speed varying 1200 rpm to 4200 rpm. Current predictions and results in Ref. [A.5]. average Re from 1201 to 4774.

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APPENDIX B. A MODEL FOR BEARING LINER DEFORMATION

As the external load increases, the fluid film in a hydrodynamic bearing becomes thinner, thus producing an increase in fluid temperature that reduces the oil viscosity, thus further decreasing the oil film thickness [B.1]. Conventional hydrodynamic bearings commonly adopt Babbitt (white metal) as a liner material to protect the surfaces of rotors while ensuring low friction. Though the Babbitt liner provides good conformability and embedability, it loses strength with a raise in temperature, and has a high breakaway friction [B.1]. As a substitute for Babbitt, PTFE (polytetrafluoroethylene) shows better performance at high temperatures with a low breakaway friction [B.2]. However, the high flexibility of PTFE affects the static and dynamic forced performance of a fluid film bearing. Thus, a predictive model for bearings with a PTFE liner should account for its deformation.

The elastic radial deformation (δ) of a liner due to pressure is approximated as a plane strain problem [B.1-B.4],

$$\delta = \frac{s(P - P_a)(1 + \nu)(1 - 2\nu)}{E(1 - \nu)} \quad (\text{B.1})$$

where P is the hydrodynamic pressure, P_a is the ambient pressure, and s is the thickness of the liner. E and ν denote the elastic modulus and the Poisson's ratio for the liner material, respectively.

Ref. [B.1] also considers the thermal expansion (ΔR) of the liner into account, as

$$\Delta R = -\alpha s (T - T_a) \quad (\text{B.2})$$

where α is the thermal expansion coefficient of the liner, T is the linear surface temperature and T_a is the ambient temperature. The simple thermal expansion Eq. (B.2) is implemented into the current model as a change in bearing or pad clearance.

Kuznetsov *et al.* [B.1] analyze the effect of liner flexibility on the static performance of a plain journal bearing. Table B.1 lists the geometry and lubrication properties of the bearing in Ref. [B.1], and Figure B.1 shows a schematic view of the plain journal bearing.

Table B.1 Parameters of a plain journal bearing in Ref. [B.1]

Rotor diameter, D	198.5 mm
Bearing length, L	200 mm
Bearing cold clearance, C	237 μm
Bearing supply groove width	29.22°
Supply oil temperature	65°C
Oil viscosity (at 40°C)	0.033 Pa.s
Oil viscosity (at 100°C)	0.0056 Pa.s
Oil density	864.7 kg/m ³
Oil heat capacity	2008.5 J/(kg.K)
Oil thermal conductivity	0.13 W/(mK)
Unit load, $W/(LD)$	504-6,297 kPa
Journal speed	900 rpm

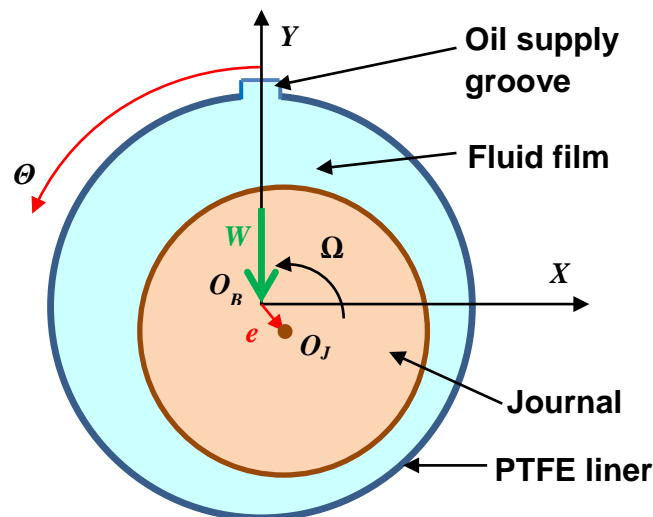
**Figure B.1 Schematic view of plain journal bearing with a liner, as per Ref. [B.1]**

Table B.2 lists the properties of the PTFE liner. Note that Ref. [B.1] does not inform on the ambient temperature. The current predictions assume the ambient temperature (T_a) to equal the the supply temperature 65°C.

Table B.2 PTFE liner properties in Ref. [B.1].

Elastic modulus, E	0.11 GPa
Poisson ratio, ν	0.46
thickness, s	2 mm
thermal expansion coefficient, α	$1.35 \times 10^{-4} \text{ K}^{-1}$

Ref. [B.1] also accounts for the thermal expansion of the shaft and bearing housing. Current predictive model offers a function to estimate the thermal expansion of the shaft as well. The equation to estimate the thermal expansion of the shaft is

$$\Delta R = -\alpha_{shaft} R (T_{shaft} - T_a) \quad (\text{B.3})$$

where α_{shaft} is the thermal expansion coefficient of the shaft, T_{shaft} is the temperature in shaft, and R is the shaft radius. In the current predictions, α_{shaft} is $1.11 \times 10^{-5} \text{ K}^{-1}$ for a steel journal.

Two cases are analyzed and predictions compared to the results in Ref. [B.1] (see Table B.3). Note that the deformation of the PTFE liner contains both the thermal expansion and mechanical deformation.

Table B.3 Cases to assess effect of liner flexibility on the performance of a plain journal bearing.

Case 1	Bearing with rigid liner
Case 2	Bearing with PTFE liner

Figure B.2 shows the fluid film thickness predicted by the current model for a 122 kN static load applied on the bearing ($W/(LD)=3,073 \text{ kPa}$). For both cases (rigid bearing and bearing with PTFE liner), the fluid film thickness predicted by the current model correlates well with that presented in Ref. [B.1]. The flexibility of PTFE liner tends to reduce the film thickness.

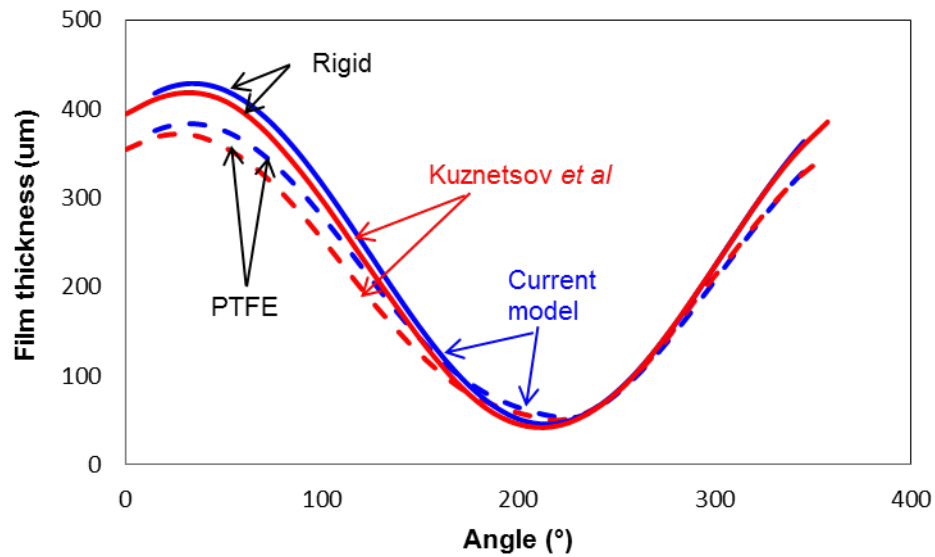


Figure B.2 Fluid film thickness at the bearing mid-plane. Static load $W/(LD)=3,073$ kPa, rotor speed $\Omega=900$ rpm. Current predictions and data from Ref. [B.1].

Figure B.3 shows the fluid film pressure at the bearing mid-plane; the current predictions show good correlation with those in Ref [B.1]. For a rigid bearing, the current model predicts a larger peak pressure. The peak fluid film pressure including the elasticity of the PTFE liner is 27% smaller than that assuming a rigid liner.

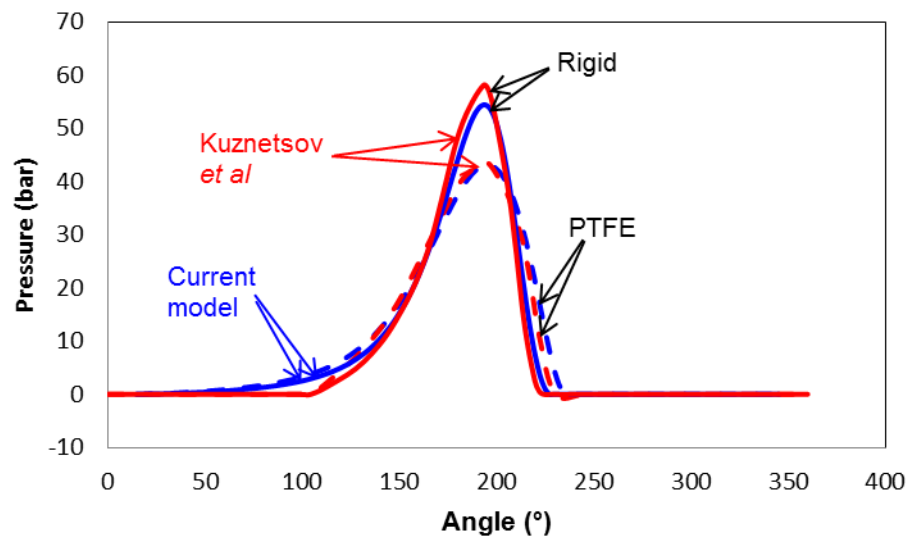


Figure B.3 Fluid film pressure at the bearing mid-plane. Static load $W/(LD)=3,073$ kPa, rotor speed $\Omega=900$ rpm. Current predictions and data from Ref. [B.1].

Figure B.4 depicts the maximum pressure, minimum fluid film thickness and eccentricity ratio versus unit load. The current predictions of maximum pressure and minimum fluid film thickness show good agreement with the results in Ref.[B.1]. The current model delivers a larger eccentricity ratio (=eccentricity/cold clearance) than that in Ref. [B.1]. The current predictions for the maximum temperature correlate well with the data in Ref. [B.1] at a low loads, $W/(LD)<2.5$ MPa; but are larger at high loads, $W/(LD)>4$ MPa. The PTFE liner tends to reduce the maximum film pressure and journal eccentricity ratio. The effect of PTFE liner flexibility (with thickness $s=2$ mm) on the minimum film thickness and maximum temperature is not significant.

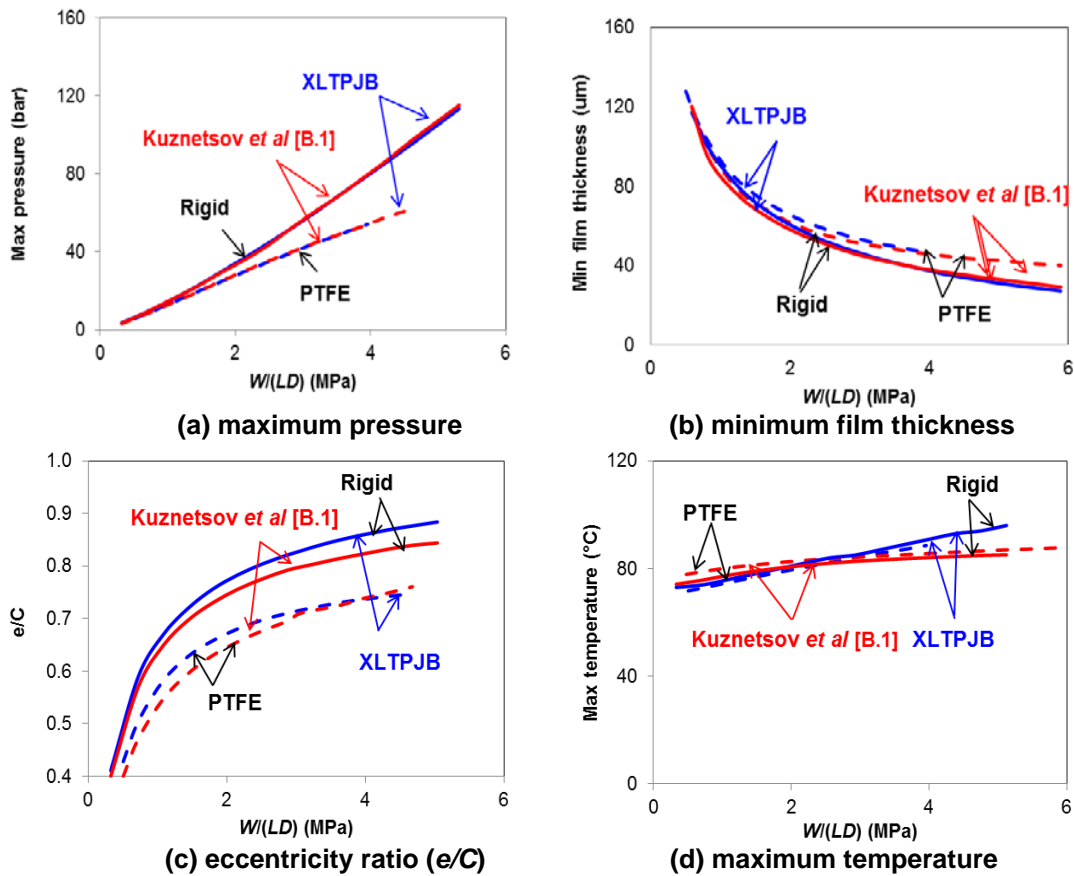


Figure B.4 (a) Maximum pressure, (b) minimum fluid film thickness and (c) eccentricity ratio and (d) maximum temperature at different load. $504 \text{ kPa} < W/(LD) < 6,297 \text{ kPa}$. $\Omega=900\text{rpm}$. Current predictions and data from Ref. [B.1].

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