

# SWINGING GONDOLA - FUN RIDE!

MEEN 363 - SP09 (LSA)

The figure shows a SWINGING GONDOLA (SG) at an amusement park. The gondola's mass and centroidal radius of gyration equal  $M$  and  $r_k$ . The total mass of the passengers does not affect significantly the magnitudes given above or CG location of SG. Two bars of length  $L$  and mass  $M_b$ ,  $I_b = (1/12) M_b L^2$ , attach the SG CG to a shaft connected to a drive motor. The shaft is supported on frictionless bearings. The motor applies torque  $T_m$  turning the bars and lifting the gondola to a release angle  $\theta_o$ . At this time, the motor is disengaged and the gondola starts to swing, giving a good thrill to its riders. Determine:

- System EOM after the motor is disengaged. Express the angular acceleration  $\ddot{\theta}$  in terms of the system parameters &  $\theta$ . [10]
- Given the (large) initial angle  $\theta_o$ , derive an analytical expression (symbolic) for the SG angular velocity  $\dot{\theta}$  in terms of the system parameters, and angles  $\theta$  and  $\theta_o$ . [10]
- Given,  $M=4000$  kg,  $r_k=4$  m,  $L=10$  m,  $M_b=0.05 M$ , and  $\theta_o = 125^\circ$ , calculate, when gondola crosses  $\theta=0^\circ$ , its angular speed [Hz], and SG CG tangential speed and radial acceleration. What is the type of "thrill" or sensation the riders feel when the SG passes through  $\theta=0^\circ$ ? [2.5 x 4]

## Data:

$$\theta_{o\_} := 125 \quad \text{degree} \quad \text{release angle}$$

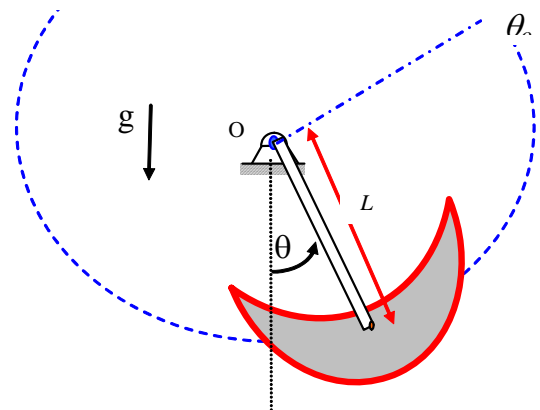
$$M := 4000 \cdot \text{kg} \quad L := 10 \cdot \text{m} \quad r_k := 4 \cdot \text{m}$$

$$M_b := 0.05 \cdot M$$

From Tables: (mass moment of inertia)

$$I_{b\_cg} = M_b \cdot \frac{L^2}{12} \quad I_{SG} := M \cdot r_k^2$$

bar Gondola inertia about its cg



## Moments of inertia about pivot O:

$$I_{b\_o} := M_b \cdot \frac{L^2}{12} + M_b \cdot \left(\frac{L}{2}\right)^2 \quad I_{SG\_o} := I_{SG} + M \cdot L^2 \quad \frac{I_{b\_o}}{I_{SG\_o}} = 0.014$$

**Total Mass moment of inertia about pivot O:**  $I_o := 2 \cdot I_{b\_o} + I_{SG\_o}$

## Equation of motion:

From summation of moments about pivot:  $I_o \cdot \frac{d^2}{dt^2} \theta = -M \cdot g \cdot L \cdot \sin(\theta) - 2 \cdot M_b \cdot g \cdot \frac{L}{2} \cdot \sin(\theta)$  assume: no friction

$$I_o \cdot \frac{d^2}{dt^2} \theta + g \cdot L \cdot \sin(\theta) \cdot (M + M_b) = 0 \quad [2] \quad \frac{d^2}{dt^2} \theta = \frac{-g \cdot L \cdot (M + M_b)}{I_o} \cdot \sin(\theta)$$



For **small amplitude motions**, Eqn(2) reduces to:

$$\sin(\theta) \equiv \theta$$

$$I_o \cdot \frac{d^2}{dt^2} \theta + g \cdot L \cdot (M + M_b) \cdot \theta = 0 \quad [3]$$

**natural frequency**, for small amplitude motions:

$$\omega_n := \left[ \frac{g \cdot L \cdot (M + M_b)}{I_o} \right]^{0.5}$$

**Natural period**

$$T_n := \frac{2 \cdot \pi}{\omega_n}$$

$$T_n = 6.764 \text{ s}$$

**Disregard bar masses**

$$\omega_{na} := \left[ \frac{\frac{g}{L} \cdot \frac{1}{1 + \left(\frac{r_k}{L}\right)^2}}{\right]^{0.5}$$

$$T_{na} := 2 \cdot \pi \cdot \left[ \frac{L}{g} \cdot \left[ 1 + \left(\frac{r_k}{L}\right)^2 \right] \right]^{0.5}$$

$$T_{na} = 6.834 \text{ s}$$

**Not much difference if bars are omitted from analysis**

### **SOLUTION for SMALL ANGLES $\theta$ - LINEARIZED EOM:**

The solution for the linear EOM (3) is  $\theta(t) = \theta_o \cdot \cos(\omega_n \cdot t) + \omega_o \cdot \sin(\omega_n \cdot t)$

where  $\theta_o := \theta_{o-} \cdot \frac{\pi}{180}$   $\omega_o = \frac{d}{dt} \theta$  are initial conditions

since  $\omega_o = 0$  starting from rest. Then

$$\theta(t) = \theta_o \cdot \cos(\omega_n \cdot t) \quad \text{and} \quad \omega(t) := -\theta_o \cdot \omega_n \cdot \sin(\omega_n \cdot t)$$

from these 2 equations, square both sides  $\frac{\theta^2}{\theta_o^2} = \cos(a)^2$   $\frac{\omega^2}{\theta_o^2 \cdot \omega_n^2} = \sin(a)^2$   $a = \omega_n \cdot t$

and add both equations to give

$$\frac{\omega^2}{\omega_n^2} = \theta_o^2 - \theta^2$$

Thus, the **linearized solution is**

$$\omega_{Lin}(\theta) := \omega_n \cdot (\theta_o^2 - \theta^2)^{0.5}$$

## Find angular speed of gondola for large release angles

From Conservation of Energy,  $T+V = V_0$ ; since gondola is released from initial angle  $\theta_0$  with null speed

**Kinetic energy**  $T = \frac{1}{2} \cdot I_0 \cdot \left( \frac{d}{dt} \theta \right)^2$

**Potential Energy**  $V = (M + M_b) \cdot g \cdot L \cdot (1 - \cos(\theta))$

$V_0 = (M + M_b) \cdot g \cdot L \cdot (1 - \cos(\theta_0))$

obtain:

$$\frac{1}{2} \cdot \left( \frac{d}{dt} \theta \right)^2 = \frac{(M + M_b) \cdot g \cdot L \cdot (\cos(\theta) - \cos(\theta_0))}{I_0} = \omega_n^2 \cdot (\cos(\theta) - \cos(\theta_0))$$

where

$$\omega_n := \left[ \frac{g \cdot L \cdot (M + M_b)}{I_0} \right]^{0.5}$$

Then, the gondola angular speed  $d\theta/dt = \omega$ :

$$\frac{d}{dt} \theta = \omega(\theta) := \omega_n \cdot [2 \cdot (\cos(\theta) - \cos(\theta_0))]^{0.5}$$

$$f(\theta) := \frac{\omega(\theta)}{2 \cdot \pi} \text{ in Hz} \quad f_n := \frac{\omega_n}{2 \cdot \pi}$$

**Tangential speed and normal acceleration at  $\theta=0$**   $f(0) = 0.262 \text{ Hz}$

$$f_n = 0.148 \text{ Hz}$$

$$V_t(\theta) := L \cdot \omega(\theta)$$

$$V_t(0) = 16.479 \frac{\text{m}}{\text{s}}$$

$$\frac{\omega(0)}{\omega_n} = 1.774$$

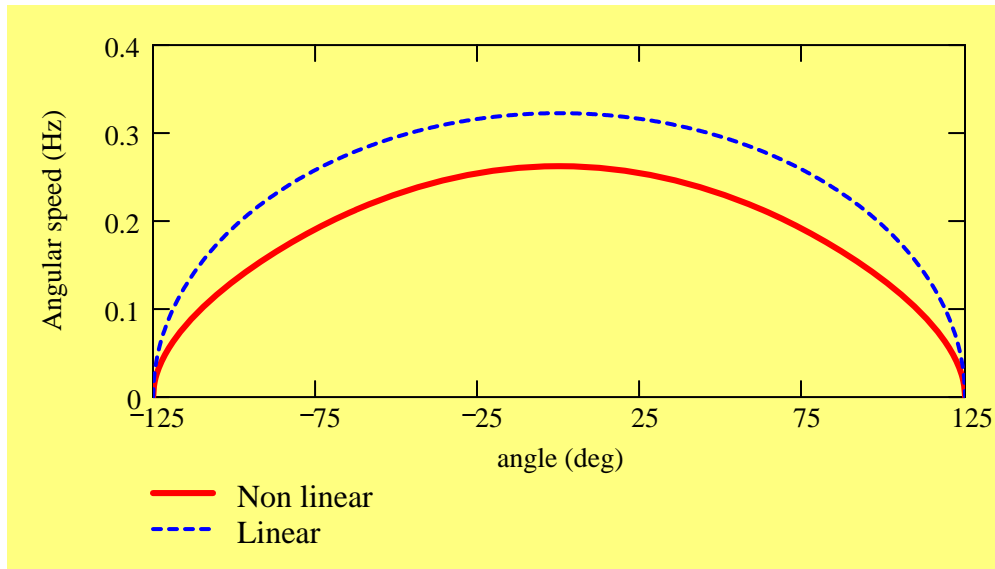
$$a_n(\theta) := -L \cdot \omega(\theta)^2$$

$$a_n(0) = -27.156 \frac{\text{m}}{\text{s}^2}$$

$$\frac{a_n(0)}{g} = -2.769$$

rider feels something in his/her tummy!  
(a weighless sensation)

Graphs: variation of gondola angular speed and radial acceleration for entire ride:



$\theta_{o\_} = 125$  degrees

Angular speed of gondola derived from NONLINEAR solution is smaller than using LINEAR solution.

The differences increase as the release angle increases

Greatest difference when  $\theta_o = 180$  deg

